1	The impact of thermal insulation on cooling energy consumption and optimal
2	insulation thickness for underground tunnel
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8	Abstract: The provision of mechanical cooling in deep mines comes with a significant energy cost
9	as a significant amount of heat transfers from surrounding rock to the airflow. Thermal insulation can
10	be applied to reduce such heat transfer, thereby cutting down the cooling load. In this study, the impact
11	of thermal insulation on reducing the heat flux through the rock was analytically investigated. The
12	optimal insulation thickness, life cycle saving and the payback period were also evaluated by using
13	the life cycle cost method as the economic benefit is heavily dependent on the insulation thickness.
14	Results show that heat flux between tunnel and airflow can be significantly reduced by the use of
15	thermal insulation, but the reduction varies with the tunnel and insulation conditions. The total cost
16	associated with using the thermal insulation firstly decreases and then increases when the insulation
17	thickness increases, implying an optimal insulation thickness. Nonetheless, both the optimal
18	insulation thickness and maximum life cycle saving can be increased by a rising rock temperature,
19	eventually leading to a reduced payback period.
20	Keywords: thermal insulation; cooling energy consumption; optimal insulation thickness; life cycle

21 cost analysis; underground tunnel

23	Nome	nclature
24	$C_0$	insulation cost ( $\frac{1}{m}$ )
25	$C_1$	electricity cost (¥/m)
26	$C_{\rm e}$	electricity price (¥/kWh)
27	$C_{t}$	total cost (¥/m)
28	Cp	specific heat capacity of air (kJ/(kgK))
29	СОР	coefficient of performance of the cooling system
30	$\Delta C_1$	difference between the electricity cost with and without insulation ( $\frac{1}{m}$ )
31	D	ratio of down payment to initial investment
32	d	market discount rate
33	Fo	Fourier number $(\alpha_1 t/r_1^2)$
34	h	convection coefficient (W/(m <sup>2</sup> K))
35	i	energy price rise rate
36	$J_0, J_1$	Bessel functions of the first kind of order 0 and 1
37	$k_1$	thermal conductivity of rock (W/(mK))
38	<i>k</i> <sub>2</sub>	thermal conductivity of insulation layer (W/(mK))
39	LCS	life cycle saving (¥/m)
40	$M_{\rm s}$	ratio of the annual maintenance and operation cost to initial investment
41	т	mass flow rate of air (kg/s)
42	<i>m</i> '	loan interest rate
43	Ν	payback period (year)
44	$N_{ m L}$	term of loan (year)

45	$N_{\min}$	years over which mortgage payments contribute to the analysis
46	Ne	analysis period (year)
47	$P_1$	ratio of cycle cost to the first-year fuel cost
48	$P_2$	ratio of life cycle expenditures to the initial investment
49	PWF	present worth factor
50	Q	heat gains of air (kJ)
51	q	heat flux (W/m <sup>2</sup> )
52	q'	heat transfer rate per meter (W/m)
53	$R_{\rm v}$	ratio of resale value at the end of analysis period to initial investment
54	r	cylindrical coordinate
55	$r_1$	radius of tunnel with insulation layer (m)
56	$r_2$	radius of tunnel without insulation layer (m)
57	<i>r</i> <sub>3</sub>	radius of outer boundary (m)
58	Т	temperature (°C)
59	Ta	inlet air temperature (°C)
60	$T_0$	original rock temperature (°C)
61	$T_1$	insulation layer temperature (°C)
62	$T_2$	rock temperature (°C)
63	t	time (s)
64	$\Delta T_{\rm a}$	air temperature difference between inlet and outlet (°C)
65	$Y_0, Y_1$	Bessel functions of the 2 <sup>ed</sup> kind of order 0 and 1
	<i>c</i> 1	

66 Greek symbols

- 67  $\alpha$  thermal diffusivity (m<sup>2</sup>/s)
- 68  $\alpha_1$  thermal diffusivity of rock (m<sup>2</sup>/s)
- 69  $\alpha_2$  thermal diffusivity of insulation layer (m<sup>2</sup>/s)
- 70  $\beta$ ,  $\beta_n$  eigenvalues
- 71  $\delta$  insulation thickness (m)
- 72 subscript
- 73 1 insulation layer
- 74 2 rock
- 75 superscript
- 76 (1) unsteady state solution
- 77 (2) steady state solution

## 78 **1. Introduction**

79 The increasing mineral production rate is essential to meet the growing demand of resources due to the rapidly growing population and industrialization [1], which necessitates the extraction of the 80 81 resources from underground at greater depths. It was reported that the air temperature in the Anglo-82 gold mine in South Africa was higher than 55 °C with the mining depth exceeding 3800 m. In China, there were more than 74 coal mines with mining depth exceeding 900 m in 2015, and the original 83 84 rock temperature ranges between 35-45 °C. Recently, an increasing number of coal and metal mines 85 are facing the engineering challenge caused by hot and humid mining environment, which imposes a negative impact on miners' health, mining facilities and productivity [2]. It is imperative for deep 86 87 mines to provide a suitable working environment through cost effective cooling measures.

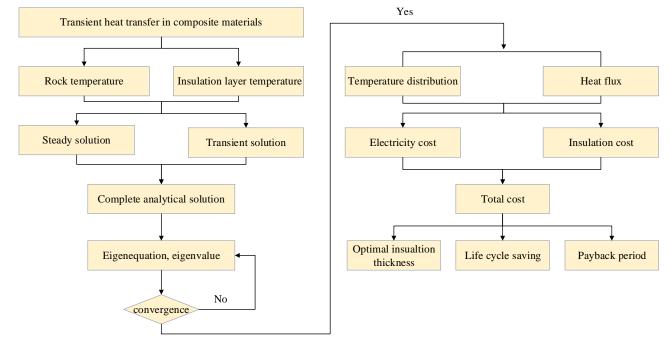
88 Various cooling systems with improved overall efficiency and operation schemes with reduced electricity cost have been adopted to achieve the required thermal environment in underground 89 90 spaces [3-5]. For instance, the operation scheme of the ice-cooling system was modified according to the peak-valley electricity tariff to reduce the electricity cost [4]. The ground air cooler, chilled 91 water close to 0 °C and ice refrigeration system were used to explore an energy efficient solution 92 for deep mines [6]. Inrushing mine water was used as the source of cooling energy to control the 93 94 environment, which helped to reduce the energy consumption of the cooling system [7]. The 95 electrical energy consumed by cooling system accounted for more than 20% of total energy 96 consumed by the coal mine. The variable speed drive technology was used to reduce the overall 97 electricity demand, with a total energy consumption reduction of 33% [8, 9]. The thermal-hydraulic 98 characteristics of an integrated mine cooling system was investigated by using a simulation model with a holistic view to assess the system's energy expenditure [10]. It was shown that the comfortable 99

100 and healthy environment conditions can be created without a significant increase in the capital cost. 101 Although above investigations indicate that these technologies can reduce the electricity cost, the cooling system's investment and electricity consumption are considerably high for mines. This is 102 103 because a significant amount of heat transfers from tunnel surrounding rock to the airflow in deep 104 underground space, which leads to a large cooling load. The cooling load can be decreased through 105 thermal insulation of the building envelopes, leading to a reduced investment and electricity cost of 106 the cooling system for residential buildings [11-13]. Intuitively, thermal insulation is applicable to reduce the overall energy consumption of cooling system in underground spaces [14]. Heat 107 108 transferred from the surrounding rock into underground space may account for more than 75% of 109 total cooling load in deep mines [15]. Therefore, thermal insulation presents the greatest potential 110 in minimizing the cooling requirement and reducing the electricity cost of cooling system [16]. Liu 111 et al. [17] studied the thermal insulation effects through various models based on one-dimensional 112 slab model with Dirichlet or Neumann boundary conditions. Many experimental investigations also indicated that the outlet airflow temperature was obviously decreased when thermal insulation was 113 114 applied [18-21].

However, greater insulation thickness can increase the materials' costs. An economic assessment should be carried out on the insulation design for a given underground tunnel to identify the economic balance between the decrease in the electricity consumption and the increase in the insulation cost, and such an approach has been broadly adopted in the industrial and residential buildings [22-25]. Ozel et al. [26] investigated the optimum insulation thickness of building walls by using life cycle cost method coupling with the consideration of the environmental impact analysis. Motaghian et al. [27] reported that the optimum insulation arrangements should be determined by

122	employing multi-objective optimization. Dombayci [28] evaluated the environmental impact of the
123	optimal thickness of external wall insulation for buildings and found out that the energy expenditure
124	was reduced by 46.6%. More reduction in heating and cooling loads (86.63%) was achievable when
125	the optimal insulation thickness ranging between 3-5 cm was used to insulate external walls [29].
126	Kecebas et al. [30] and Erturk et al. [31-33] investigated the optimal insulation thickness of pipes
127	by using life cycle cost method, in which the effect of air gap was considered. The optimal insulation
128	thickness in similar engineering areas was also investigated by Huang et al. [34] using the life cycle
129	cost analysis. Adamczyk et al. [35] proposed that the investment of thermal insulation was beneficial
130	for ecological environment because the energy demand was reduced. Such improvements on the
131	thermal performance economically benefit the application due to the reduced electricity cost. The
132	capital cost associated with the use of thermal insulation can be recovered within a certain period of
133	time. Therefore, the payback period is usually used as a direct measure of the economic benefit of
134	an invested technical improvement [36]. An investment payback period ranged from 1 year to 2.3
135	years was achieved when the rock wool and polystyrene were used as the insulation material based
136	on the life cycle cost analysis [37]. A shorter payback period was also possible when different
137	configurations were conducted [38, 39].

Researchers have made a tremendous progress in studying the impact of thermal insulation on building. However, the thermal insulation effect, including optimal insulation thickness, life cycle saving and the payback period in deep mines is rarely investigated. Especially, the effect of various rock temperatures on the optimal insulation thickness has not been reported, which indicates that a further investigation on the economic benefit analysis on the thermal insulation effect is needed. This paper is committed to investigating how the cooling load and the total cost are affected by thermal insulation and understanding how to identify the optimal insulation thickness for a given underground tunnel. Firstly, the cylindrical heat conduction model with Robin boundary condition was solved by using the separation-of-variables method. Secondly, based on above analytic results, the effects of insulation thickness, thermal conductivities of insulation layer and rock, tunnel radius and convection coefficient on the heat flux reduction were investigated. Finally, we investigated the optimal insulation thickness, life cycle saving and payback period under various rock temperatures by employing the life cycle cost analysis. The flowchart of this study is shown in Fig. 1.



152

151

Fig. 1. Flowchart of this study.

# 153 2. Mathematical description

- 154 *2.1. Heat transfer model*
- 155 2.1.1. Basic equations

156 When the air flows through the underground tunnel, heat is dissipated from tunnel rock to the

- 157 airflow driven by the temperature difference. To quantify this heat flux in a practical and reasonable
- 158 way, the following assumptions are made to facilitate the analytic model [40, 41]: (1) The thermal-

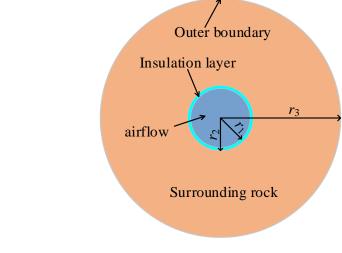
physical parameters of tunnel rock and insulation layer are constant, homogeneous and isotropic. (2) The tunnel rock is assumed dry, which means that the mass transfer between rock and airflow is not present. (3) The contact resistance between rock and insulation layer is negligible. (4) The crosssection shape of underground tunnel is circular. Therefore, the schematic diagram of the engineering problem can be illustrated by Fig. 2, which simplifies the heat exchange between the tunnel rock and insulation layer into one-dimensional transient heat conduction process in the cylindrical coordinate.

166 The governing equation is described by Eq. (1) [42]:

167 
$$\frac{\partial T_i}{\partial t} = \alpha_i \left( \frac{\partial^2 T_i}{\partial r^2} + \frac{1}{r} \frac{\partial T_i}{\partial r} \right)$$
(1)

168 where T is temperature, t is time, r is cylindrical coordinate,  $\alpha$  is thermal diffusivity. For the

169 subscript, i=1 represents the insulation layer, and i=2 represents the rock.



170 171

Fig. 2. Schematic of analytic model.

172 The initial and boundary conditions are given by Eq. (2) and Eq. (3):

173 
$$\begin{cases} T_1(r,0) = T_0 \\ T_2(r,0) = T_0 \end{cases}$$
(2)

174
$$\begin{cases} r = r_3 : T_2(r_3, t) = T_0 \\ r = r_1 : k_1 \frac{\partial T_1}{\partial r} = h(T_1 - T_a) \end{cases}$$
(3)

where  $T_0$  is the initial temperature for both rock and insulation layer;  $r_1$  is the radius of tunnel with insulation layer;  $r_3$  is the radius of outer boundary;  $T_a$  is air temperature; k is thermal conductivity. Because the contact thermal resistance is negligible, the contact boundary condition is described by Eq. (4):

179 
$$r = r_2: \quad \begin{cases} T_1(r_2, t) = T_2(r_2, t) \\ k_1 \frac{\partial T_1(r_2, t)}{\partial r} = k_2 \frac{\partial T_2(r_2, t)}{\partial r} \end{cases}$$
(4)

180 where  $r_2$  is radius of tunnel without insulation layer.

181 2.1.2. Analytic solution

182 The Eqs. (1)-(4) are solved by using the variable-separation-method [43]. The detailed process is

183 shown in the Appendix. The analytic solution is described by Eq. (5) and Eq. (6):

184 
$$T_{1}^{(1)}(r,t) = \sum_{n=1}^{\infty} A_{1n} \left[ J_{0}(\frac{\beta_{n}}{\sqrt{\alpha_{1}}}r) + \frac{B_{1n}}{A_{1n}}Y_{0}(\frac{\beta_{n}}{\sqrt{\alpha_{1}}}r) \right] \cdot \exp(-\beta_{n}^{2}t)$$
(5)

185 
$$T_{2}^{(1)}(r,t) = \sum_{n=1}^{\infty} A_{1n} \left[ \frac{A_{2n}}{A_{1n}} J_{0}(\frac{\beta_{n}}{\sqrt{\alpha_{2}}}r) + \frac{B_{2n}}{A_{1n}} Y_{0}(\frac{\beta_{n}}{\sqrt{\alpha_{2}}}r) \right] \cdot \exp(-\beta_{n}^{2}t)$$
(6)

186 The derivation of Eq. (5) and Eq. (6) can be obtained according to derivatives of *Bessel* function:

187 
$$\frac{\partial T_1^{(1)}(r,t)}{\partial r} = -\sum_{n=1}^{\infty} A_{1n} \cdot \exp(-\beta_n^2 t) \frac{\beta_n}{\sqrt{\alpha_1}} \left[ J_1(\frac{\beta_n}{\sqrt{\alpha_1}} r) + B_{1n}' Y_1(\frac{\beta_n}{\sqrt{\alpha_1}} r) \right]$$
(7)

188 
$$\frac{\partial T_2^{(1)}(r,t)}{\partial r} = -\sum_{n=1}^{\infty} A_{1n} \cdot \exp(-\beta_n^2 t) \frac{\beta_n}{\sqrt{\alpha_2}} \left[ A_{2n} J_1(\frac{\beta_n}{\sqrt{\alpha_2}} r) + B_{2n} Y_1(\frac{\beta_n}{\sqrt{\alpha_1}} r) \right]$$
(8)

189 where,  $B_{1n} = B_{1n} / A_{1n}$ ,  $B_{2n} = B_{2n} / A_{1n}$ ,  $A_{2n} = A_{2n} / A_{1n}$ .

190 The boundary condition at  $r=r_1$  is transformed to Eq. (9):

191 
$$-k_1 \frac{\beta_n}{\sqrt{\alpha_1}} J_1(\frac{\beta_n}{\sqrt{\alpha_1}} r) - h J_0(\frac{\beta_n}{\sqrt{\alpha_1}} r_1) = B_{1n} \left[ h Y_0(\frac{\beta_n}{\sqrt{\alpha_1}} r_1) + k_1 \frac{\beta_n}{\sqrt{\alpha_1}} Y_1(\frac{\beta_n}{\sqrt{\alpha_1}} r) \right]$$
(9)

192 In addition, the boundary conditions at  $r=r_2$  and  $r=r_3$  are changed to:

193
$$\begin{cases} A_{2n}^{'}J_{0}\left(\frac{\beta_{n}}{\sqrt{\alpha_{2}}}r_{3}\right)+B_{2n}^{'}Y_{0}\left(\frac{\beta_{n}}{\sqrt{\alpha_{2}}}r_{3}\right)=0\\ A_{2n}^{'}J_{0}\left(\frac{\beta_{n}}{\sqrt{\alpha_{2}}}r_{2}\right)+B_{2n}^{'}Y_{0}\left(\frac{\beta_{n}}{\sqrt{\alpha_{2}}}r_{2}\right)=J_{0}\left(\frac{\beta_{n}}{\sqrt{\alpha_{1}}}r_{2}\right)+B_{1n}^{'}Y_{0}\left(\frac{\beta_{n}}{\sqrt{\alpha_{1}}}r_{2}\right) \end{cases}$$
(10)

194 Then, the coefficients of  $A'_{2n}$ ,  $B'_{2n}$  are computed and described by Eq. (11):

195
$$\begin{bmatrix} A_{2n} \\ B_{2n} \end{bmatrix} = \begin{bmatrix} J_0 \left(\frac{\beta_n}{\sqrt{\alpha_2}} r_3\right) & Y_0 \left(\frac{\beta_n}{\sqrt{\alpha_2}} r_3\right) \\ J_0 \left(\frac{\beta_n}{\sqrt{\alpha_2}} r_2\right) & Y_0 \left(\frac{\beta_n}{\sqrt{\alpha_2}} r_2\right) \end{bmatrix}^{-1} \begin{bmatrix} 0 \\ J_0 \left(\frac{\beta_n}{\sqrt{\alpha_1}} r_2\right) + B_{1n} Y_0 \left(\frac{\beta_n}{\sqrt{\alpha_1}} r_2\right) \end{bmatrix}$$
(11)

196 Consequently, the eigenequation is obtained as:

197 
$$\frac{k_1}{\sqrt{\alpha_1}} \cdot \begin{bmatrix} 1 & B_{1n} \end{bmatrix} \cdot \begin{bmatrix} J_1(\frac{\beta_n}{\sqrt{\alpha_1}}r_2) \\ Y_1(\frac{\beta_n}{\sqrt{\alpha_1}}r_2) \end{bmatrix} - \frac{k_2}{\sqrt{\alpha_2}} \cdot \begin{bmatrix} A_{2n} & B_{2n} \end{bmatrix} \cdot \begin{bmatrix} J_1(\frac{\beta_n}{\sqrt{\alpha_2}}r_2) \\ Y_1(\frac{\beta_n}{\sqrt{\alpha_2}}r_2) \end{bmatrix} = 0$$
(12)

The eigenequation is a transcendental equation, which can only be solved through numerical measure. The eigenvalues are calculated by employing MATLAB in this study, and the solution's convergence is also investigated. The analytic solution does not converge when the first 15 eigenvalues are used. The analytic solution converges when the first 30 eigenvalues are employed, which is consistent with the solution of using the first 500 eigenvalues. Therefore, the first 30 eigenvalues are used for the following computation.

The coefficients of  $B'_{1n}, A'_{2n}, B'_{2n}$  are calculated by using the eigenvalues, and  $A_{1n}$  is determined by the initial condition presented in Eq. (A10). Substituting the general solution into the initial condition, 206 Eq. (13) is obtained as:

207  

$$\begin{cases}
T_{1}^{(1)}(r,0) = \sum_{n=1}^{\infty} A_{1n} \left[ J_{0}(\frac{\beta_{n}}{\sqrt{\alpha_{1}}}r) + B_{1n}'Y_{0}(\frac{\beta_{n}}{\sqrt{\alpha_{1}}}r) \right] = \sum_{n=1}^{\infty} A_{1n}\varphi_{1n}(r) \\
T_{2}^{(1)}(r,0) = \sum_{n=1}^{\infty} A_{1n} \left[ A_{2n}'J_{0}(\frac{\beta_{n}}{\sqrt{\alpha_{2}}}r) + B_{2n}'Y_{0}(\frac{\beta_{n}}{\sqrt{\alpha_{2}}}r) \right] = \sum_{n=1}^{\infty} A_{1n}\varphi_{2n}(r)
\end{cases}$$
(13)

208 where, 
$$\varphi_{1n}(r) = J_0(\frac{\beta_n}{\sqrt{\alpha_1}}r) + B_{1n}Y_0(\frac{\beta_n}{\sqrt{\alpha_1}}r), \quad \varphi_{2n}(r) = A_{2n}J_0(\frac{\beta_n}{\sqrt{\alpha_2}}r) + B_{2n}Y_0(\frac{\beta_n}{\sqrt{\alpha_2}}r)$$

209 The orthogonal expression for  $\varphi_{1n}(r)$  and  $\varphi_{2n}(r)$  is described by Eq. (14):

210 
$$\sum_{i=1}^{2} \frac{k_{i}}{\alpha_{i}} \int_{r_{i}}^{r_{i+1}} r \varphi_{1n}(r) \varphi_{2n}(r) dr = \begin{cases} 0 & m \neq n \\ N_{n} & m = n \end{cases}$$
(14)

211 where, 
$$N_n = \sum_{i=1}^2 \frac{k_i}{\alpha_i} \int_{r_i}^{r_{i+1}} r \varphi_{in}^2(r) dr$$

212 From Eq. (14), Eq. (15) can be obtained:

213 
$$\frac{k_1}{\alpha_1} \int_{r_1}^{r_2} T_1^{(1)}(r,0) \cdot r \cdot \varphi_{1n}(r) dr + \frac{k_2}{\alpha_2} \int_{r_2}^{r_3} T_2^{(1)}(r,0) \cdot r \cdot \varphi_{2n}(r) dr = A_{1n} \left[ \frac{k_1}{\alpha_1} \int_{r_1}^{r_2} r \cdot \varphi_{1n}^2(r) dr + \frac{k_2}{\alpha_2} \int_{r_2}^{r_3} r \cdot \varphi_{2n}^2(r) dr \right]$$
(15)

214 The coefficient  $A_{1n}$  is shown as:

215 
$$A_{1n} = \frac{1}{N_n} \left[ \frac{k_1}{\alpha_1} \int_{r_1}^{r_2} T_1^{(1)}(r,0) \cdot r \cdot \varphi_{1n}(r) dr + \frac{k_2}{\alpha_2} \int_{r_2}^{r_3} T_2^{(1)}(r,0) \cdot r \cdot \varphi_{2n}(r) dr \right]$$
(16)

### 216 2.2. Life cycle cost method

The economic benefit of thermal insulation is investigated by adopting the life cycle cost (LCC) analysis, which is a widely used method for evaluating all relevant costs over the life cycle time [22, 30, 39, 44]. The main procedures of LCC analysis are described as follows: (1) The initial cost is calculated. (2) The operation cost is evaluated when the influences of inflation and interest rate are considered. (3) The minimum total cost is then determined to provide the optimal insulation thickness.

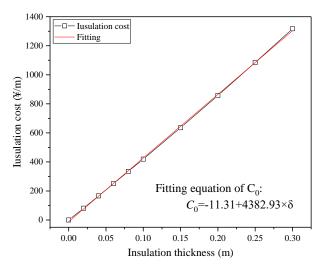
222 (4) The payback period and the life cycle saving are obtained.

The lightweight geopolymer concrete (LGC) with lower greenhouse gas emission is used as the insulation material in this study. The LGC with thermal conductivity of 0.1 W/(mK) and dry density of 450 kg/m<sup>3</sup> was produced by using fly ash and alkali solution at room temperature [45]. The price of insulation layer is about 325  $\pm$ /m<sup>3</sup> ( $\pm$ : RMB). The insulation cost per tunnel length under various insulation thicknesses is calculated.

Fig. 3 indicates that the insulation cost is approximately linearly increased when the insulation thickness increases from 0 m to 0.3 m. For instance, the insulation cost increases from 81.25/m to 1317/m when the insulation thickness increases from 0.02 m to 0.3 m. The empirical correlation between the insulation cost and the insulation thickness can be described by Eq. (17):

233 
$$C_0 = -11.31 + 4382.93 \times \delta$$
 (17)

where  $C_0$  is the insulation cost per tunnel length.  $\delta$  is the insulation thickness.



235 236

Fig. 3. The insulation cost versus thickness of insulation layer

237 2.2.2. Annual electricity cost

The annual electricity cost  $(C_1)$  is dependent on the total heat gain of the airflow. Due to the unsteady process of the convective heat transfer, the annual heat gain (cooling load) is calculated by the integration method:

241

256

$$Q = m \cdot \int_0^{365 \times 24 \times 3600} c_p \cdot \Delta T_a dt \tag{18}$$

where Q is the heat gain of air,  $c_p$  is the specific heat capacity of air, m is the mass flow rate,  $\Delta T_a$  is the temperature difference between inlet and outlet air.

The economic evaluation can be carried out through the  $P_1$ - $P_2$  economical method, which is shown as follows [34]:

246 
$$P_{1} = PWF(Ne, i, d) = \sum_{j=1}^{N_{e}} \frac{(1+i)^{j-1}}{(1+d)^{j}} = \begin{cases} \frac{1}{d-i} [1 - (\frac{1+i}{1+d})^{N_{e}}], i \neq d\\ \frac{Ne}{1+i}, i = d \end{cases}$$
(19)

247 
$$P_{2} = D + (1 - D) \frac{PWF(N_{\min}, 0, d)}{PWF(N_{L}, 0, m')} + M_{s} \times PWF(Ne, i, d) - \frac{R_{v}}{(1 + d)^{Ne}}$$
(20)

where  $P_1$  is the ratio of life cycle cost to the first-year electricity cost.  $P_2$  is the ratio of life cycle expenditures to the initial investment. *PWF* is the present worth factor. *Ne* is the analysis period, year.  $N_L$  is the term of loan, year. *d* is the market discount rate. *i* is the energy price rise rate. *m'* is the loan interest rate.  $N_{min}$  is the number of years over which mortgage payments contribute to the analysis (usually the minimum of *Ne* and *N<sub>L</sub>*). *M<sub>s</sub>* is the ratio of the annual maintenance and operation cost to initial investment. *R<sub>v</sub>* is the ratio of resale value at the end of analysis period to initial investment. *D* is the ratio of down payment to initial investment.

255 The annual electricity cost  $C_1$  is calculated by Eq. (21):

$$C_1 = \frac{Q}{3600 \times 250 \times \text{COP}} C_e \tag{21}$$

where COP is the coefficient of performance,  $C_e$  is the electricity price. All the parameters and their values are listed in Table 1.

### Table 1 Parameter for the life cycle cost analysis.

Parameters	Values
Interest rate $(i)$	8%
Lifetime (Ne)	10 Years
Market discount rate ( <i>d</i> )	6%
Coefficient of Performance (COP)	2.8
Electricity price $(C_e)$	1 ¥/kWh

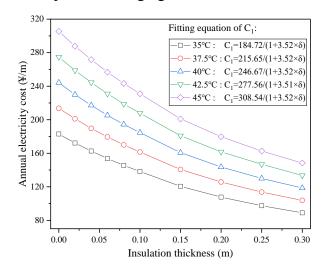
The heat gain of airflow from the tunnel with a length of 250 m, wind speed of 1 m/s, original surrounding rock temperature ranging from 35 °C to 45 °C, inlet air temperature of 20 °C and tunnel radius of 2 m is calculated after the tunnel is ventilated for 12 months. The heat gain of airflow and the annual electricity cost of the cooling system with different insulation thicknesses are listed in Table 2. It is observed that the heat gain of airflow and the annual electricity cost per meter decrease from 1844 MJ/m to 897 MJ/m and from 183 ¥/m to 89 ¥/m respectively when the insulation thickness increases from 0 to 0.3 m.

267

Table 2 Annual electricity cost versus insulation thickness.

Insulation thickness (m)	0	0.04	0.08	0.1	0.15	0.2	0.25	0.3
Heat gains (MJ/m)	1844	1640	1467	1394	1216	1086	983	897
Annual electricity cost (¥/m)	183.0	162.7	145.6	138.3	120.6	107.7	97.5	89.0

Fig. 4 shows the influence of rock temperature on the annual electricity cost under various insulation thickness. The annual electricity cost is observed to be decreased with the increased insulation thickness. It is known that the heat transfer can be significantly reduced by a larger insulation thickness, thereby giving a reduced cooling load, which then leads to a lower annual electricity cost. The empirical correlations between the annual electricity cost and the insulation thickness under various rock temperatures ranging from 35 °C to 45 °C are also presented in Fig. 4.





### Fig. 4. Annual electricity cost under various original rock temperature.

## 276 2.2.3. Total cost, life cycle saving and payback period

The total cost (C<sub>t</sub>) consists of the electricity cost over the life time and the insulation cost, which can be described by Eq. (22):

279

282

# $C_{t} = P_{1}C_{1} + P_{2}C_{0} \tag{22}$

280 Life cycle saving (LCS) is the difference between the saved energy cost over the life cycle time

and the insulation cost, which can be calculated by Eq. (23):

$$LCS = P_1 \Delta C_1 - P_2 C_0 \tag{23}$$

283 where  $\triangle C_1$  is the electricity cost difference with and without the insulation layer.

By setting Eq. (23) to zero, the payback period N can be calculated:

285 
$$N = \begin{cases} \frac{\ln[1 - ((P_2 \times C_0)/(\Delta C_1)) \times (d-i)]}{\ln((1+i)/(1+d))}, i \neq d\\ ((P_2 \times C)/(\Delta C_1)) \times (1+i), i = d \end{cases}$$
(24)

## 286 **3. Results and discussion**

### 287 3.1. Verification of analytic solution

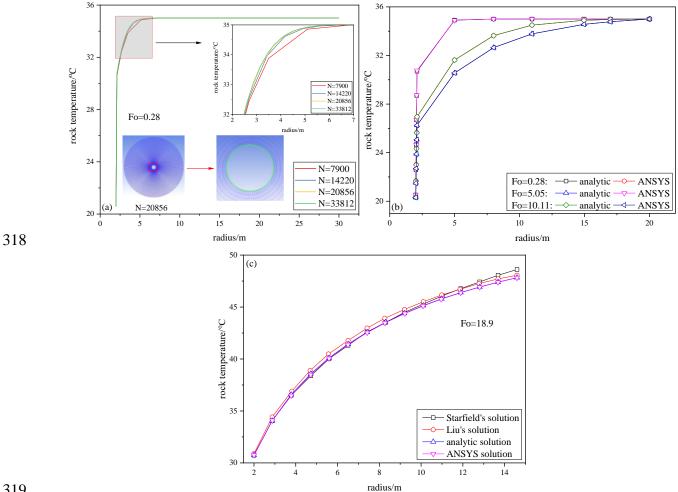
A numerical model is simulated by ANSYS Fluent 18.0 to verify the analytic solution. The quadrilateral mesh with high quality is created through ICEM CFD (Integrated Computer Engineering and Manufacturing code for Computational Fluid Dynamics). To decrease the total mesh number and shorten the computing time, uniform grid is employed in the insulation layer and nonuniform grid is used in surrounding rock region, respectively.

293 The key thermal properties and boundary conditions are listed in Table 3. The thermal properties 294 of tunnel and ventilation parameters are obtained from a coal mine in Xuzhou, China. The properties of an insulation material are provided by a literature [45]. The thermal penetration length is defined 295 296 as the depth at which the dimensionless temperature  $(T - T_a)/(T_0 - T_a)$  inside the surrounding rock 297 reduces to 0.99 based on the boundary layer theory [46]. To eliminate the boundary effect, the outer 298 boundary radius should be larger than the thermal penetration length [18, 40]. The analytic analysis 299 indicated that the outer boundary radius in Table 3 is larger than the thermal penetration depth, which is not discussed in this paper. The grid independence investigation is conducted on this numerical 300 301 model. Four grids with mesh numbers of 7900, 14220, 20856 and 33812 are generated, respectively. 302 The rock temperature distribution under different meshing are compared when the tunnel is ventilated 303 for about ten days (Fo=0.28), as shown in Fig. 5 (a). It is seen that the calculated rock temperature 304 does not vary when the grid number increased to 20856. To minimise the computational time without 305 compromising the calculation accuracy, the mesh with the grid number of 20856 is employed for the following investigation. 306

Parameters	Value
Original rock temperature, $T_0$	35 °C
Thermal conductivity of rock (TCR), $k_1$	2.4 W/(mK)
Thermal diffusivity of rock, $\alpha_1$	$1.3 \times 10^{-6}  m^2/s$
Tunnel radius, $r_1$	2 m
Convection coefficient, h	20 W/(m <sup>2</sup> K)
Inlet air temperature, $T_{\rm a}$	20 °C
Thermal conductivity of insulation layer (TCIL), $k_2$	0.1 W/(mK)
Thermal diffusivity of insulation layer, $\alpha_2$	$0.4 \times 10^{-6} \text{ m}^2/\text{s}$
Outer boundary radius, $r_3$	30 m
Insulation thickness, $\delta$ ( $\delta = r_2 - r_1$ )	0.1 m

Table 3 Key thermal properties and boundary conditions.

308 The temperature distribution calculated from the analytic solution when Fo=0.28, Fo=5.05 and 309 Fo=10.11 are compared with that given by Fluent, as shown in Fig. 5 (b). The tunnel surface temperature difference between analytic and Fluent is 0.18 °C, 0.23 °C and 0.03 °C when Fo=0.28, 310 311 5.05 and 10.11 respectively, which shows that the analytic results are in good agreement with the 312 numerical results. Starfield et al. [47] and Liu et al. [17] computed rock temperature when the tunnel 313 rock is dry. The rock temperature is also solved through the analytic solution in this study and the 314 comparison of rock temperature distribution is presented in Fig. 5 (c). The tunnel surface temperature 315 of Starfield's, Liu's, analytic and ANSYS solution is 30.73 °C, 30.89 °C, 30.75 °C and 30.75 °C 316 respectively when Fo=18.9. The maximum relative error is lower than 0.46%. Therefore, the validity 317 of the analytic solution in this study can be proved.



## 319

320 Fig. 5. Model validation: (a) grid independence verification, analytic solution versus (b) ANSYS and (c)

321

### literatures' results.

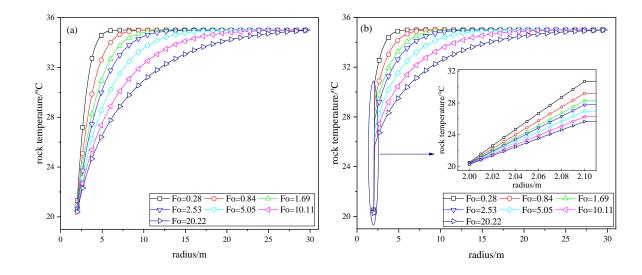
#### 322 *3.2.Thermal insulation effect*

#### 323 *3.2.1. Rock temperature*

324 The rock temperature distributions in the cross section of the tunnel without and with insulation layer are investigated, as shown in Fig. 6. As can be seen from Fig. 6 (a), the rock temperature 325 326 considerably decreases near the tunnel wall due to the cooling effect caused by the airflow with lower temperature. Such cooling effect weakens as it moves further away from the tunnel wall. The rock 327 328 temperature stays almost at a constant level when the distance is above a certain threshold, which is the length of the outer boundary that needs to be determined in the analytic model. The rock 329

330 temperature is expected to decrease rapidly at the initial stage due to the transient heat transfer process. 331 After a period of time, the reducing rate of the rock temperature decreases gradually. For instance, the tunnel surface temperature decreases from 35 °C to 20.90 °C when Fo increases from 0 to 0.84, 332 333 with a temperature difference of 14.10 °C. In contrast, the rock temperature difference only reduces by 0.83 °C when Fo increases from 0.84 to 10.11. This is because that the temperature difference 334 between airflow and tunnel is reduced by the increased Fo number. Similar rock temperature 335 336 distribution is achieved by Wang et al. [40, 48] through experimental and analytic investigations. 337 Fig. 6 (b) presents the rock temperature distribution of tunnel with insulation layer. It is seen that 338 the temperature gradient in insulation layer is significantly larger than that in surrounding rock. This 339 is because that the TCIL is smaller than that of surrounding rock. For instance, the temperature 340 gradient in insulation layer is 101.79 °C/m, in comparison with that of 8.47 °C/m in surrounding rock 341 at Fo=0.28. The temperature gradient in the insulation layer also decreases when the Fo number 342 increases, which is caused by the reduced temperature difference between airflow and tunnel. As we can also see from Fig. 6, the temperature gradient near the tunnel wall is large, which suggests that 343 344 the closer the distance to the tunnel wall, the higher the heat flux is [49, 50]. Therefore, an excellent

thermal insulation performance can be achieved when the insulation layer is installed on the tunnelwall [42].





348

Fig. 6. Rock temperature of tunnel (a) without and (b) with insulation layer.

349 To verify the thermal insulation effect, the surplus temperature is monitored during the unsteady heat transfer process. The surplus temperature is defined by the temperature difference between tunnel 350 351 wall and airflow, which can be used to compute the heat convection between tunnel and airflow [42]. 352 The variation of surplus temperature with Fo is listed in Table 4, which shows that the surplus temperature of original tunnel is always higher than that of insulated tunnel. Although the surplus 353 354 temperature's ratio decreases gradually when the Fo increases, the heat convection of original tunnel 355 is still about 1.43 times higher than that of insulated tunnel when Fo=20.22. It indicates that thermal insulation can reduce about 30% of total heat convection when the tunnel is ventilated for two years. 356 Therefore, the thermal insulation can considerably reduce heat convection between tunnel and airflow, 357 358 which contributes to the decreased cooling load.

Table 4 Surplus temperature versus Fo number.

Fo	0.28	0.84	1.69	2.53	5.05	10.11	20.22
Original tunnel (°C)	1.30	0.90	0.73	0.65	0.55	0.47	0.40
Insulated tunnel (°C)	0.52	0.45	0.40	0.38	0.34	0.30	0.28

### 360 *3.2.2. Heat flux reduction*

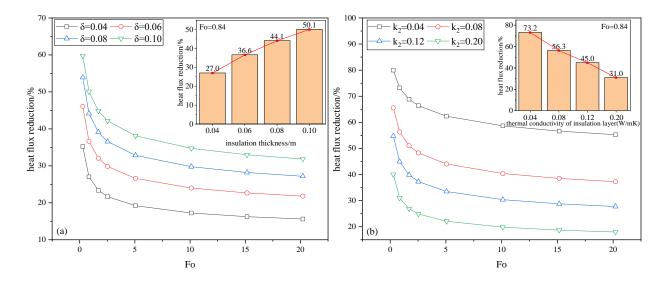
361 Because the thermal insulation effect is dependent on the parameters related to insulation layer, 362 tunnel and ventilation conditions, it is necessary to investigate the impact of insulation thickness, TCIL, TCR, tunnel radius and convection coefficient on thermal insulation effect. A series of 363 calculations of heat flux of the tunnel with/without insulation layer have been conducted under 364 various parameters. Therefore, the effects of these various parameters on heat flux reduction are 365 366 presented. Because the moisture leads to an increased TCIL and a coupled heat and mass transfer, the thermal insulation effect can be considerably reduced by moisture transfer [45]. Therefore, it may not 367 368 be feasible for wet tunnel to control underground environment by using thermal insulation, which is 369 not discussed in this study.

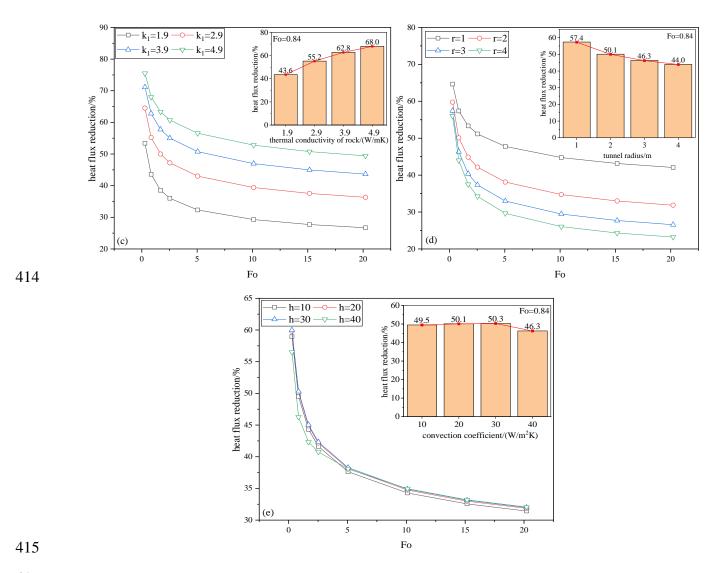
Fig. 7 shows the changes of heat flux reduction achieved by various parameters, such as insulation 370 thickness, TCIL, TCR, tunnel radius and convection coefficients. The impact of insulation thickness 371 on heat flux reduction is shown in Fig. 7 (a). Taking the profile under Fo=0.84 as an example, it is 372 found that the heat flux reduction increases when a thicker insulation is used, which indicates that 373 374 larger insulation thickness can weaken the heat transfer from tunnel rock to airflow. For instance, the heat flux reduction increases from 27.0% to 50.1% when the insulation thickness increases from 0.04 375 m to 0.1 m. However, larger insulation thickness implicates an increased insulation cost. A balance 376 377 needs to be achieved among them to achieve the best cost effectiveness, which is discussed in more details in section 3.3. Fig. 7 (b) shows that the heat flux reduction increases when the TCIL decreases. 378 379 Zhou et al. [51] also found that the heat flux reduction non-linearly decreased with the increased 380 TCIL, which is similar with the variation presented in Fig. 7 (b). This is because a lower TCIL can

381	result in a larger thermal resistance [42]. As the heat flux reduction is strongly dependent on the TCIL,
382	preference should be given to the insulation layer with a lower thermal conductivity. Zhang et al. [52]
383	also indicates that a smaller TCIL can lead to a lower heat flux and a better thermal insulation effect.
384	Fig. 7 (c) illustrates the impact of TCR on the heat flux reduction with insulation thickness of 0.1
385	m and TCIR of 0.1 W/(mK). It is seen that an increase in the TCR contributes to an enhancement in
386	heat flux reduction, suggesting that the thermal insulation effect is significant when the TCR is high.
387	For instance, the heat flux reduction increases from 43.6% to 68.0% when the TCR is increased from
388	1.9 W/(mK) to 4.9 W/(mK). Zhu [19] indicated that an increase in TCR (thermal diffusivity) could
389	lead to a quicker decrease in rock temperature, which is consistent with a larger heat flux reduction.
390	It is known that a higher TCR is equivalent to a smaller thermal resistance of surrounding rock. When
391	the thermal resistance of the surrounding rock is coupled with that of the insulation layer, the smaller
392	the thermal resistance of the surrounding rock, the larger the increasing rate of total thermal resistance
393	is. Therefore, one can expect that a higher TCR coupled with the insulation layer can lead to a
394	significant thermal insulation effect. The influence of the tunnel radius on the heat flux reduction with
395	an insulation thickness of 0.1 m and TCIR of 0.1 W/(mK) is shown in Fig. 7 (d), from which we can
396	find that the heat flux reduction decreases when the tunnel radius increases. This indicates that the
397	thermal insulation effect is significant when the tunnel radius is smaller. The reason is that the total
398	thermal resistance is increased by the decreased tunnel radius. Compared with the tunnel with a small
399	radius, a thicker insulation layer (or smaller TCIL) can achieve the same thermal insulation effect for
400	those tunnels with a large radius.

401 Fig. 7 (e) shows the impact of convection coefficient on the heat flux reduction. It can be observed
402 that the heat flux reduction is not strongly dependent on the convection coefficient. For instance, the

403 heat flux reduction only increases from 46.3% to 50.3% when the convection coefficient increases from 10 W/( $m^2$ K) to 40 W/( $m^2$ K). The heat transfer can be enhanced by increased convection 404 coefficient. However, the decreased tunnel surface temperature also leads to a reduced temperature 405 406 difference between tunnel and airflow. This indicates that the increase in the convection coefficient 407 is offset by the decrease in the temperature difference. Therefore, the thermal insulation effect does not significantly depend on the convection coefficient. Zhang [18] experimentally proved that the 408 409 dimensionless rock temperature was not strongly dependent on the Bi number (convection coefficient), which is consistent with the results showed in Fig. 7 (e). As we can see from the profiles 410 411 shown in Fig. 7, all these effects of insulation thickness, TCIL, TCR, tunnel radius and convection 412 coefficient on the heat flux reduction are non-linear.





416 Fig. 7. Heat flux reduction versus (a) insulation thickness, (b) TCIL, (c) TCR, (d) tunnel radius, (e) convection

417

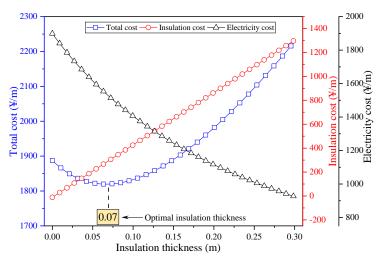
coefficient.

418 *3.3.Economic benefit analysis* 

419 3.3.1. Total cost

The insulation cost, electricity cost and the total cost under different insulation thicknesses are calculated based on the life cycle cost method. The effect of insulation thickness on the total cost is presented in Fig. 8. It is seen from Fig. 8 that the total cost firstly decreases and then increases when the insulation thickness increases. For instance, the total cost decreases from 1887  $\frac{1}{2}$ /m to 1820  $\frac{1}{2}$ /m when the insulation thickness increases from 0 m to 0.07 m. The total cost begins to increase when the insulation thickness is larger than 0.07 m. It is known that the optimum insulation thickness should be determined against the lowest total cost [53]. Therefore, the optimal insulation thickness is 0.07 m when the rock temperature is 35 °C, with the lowest total cost of 1820  $\frac{1}{2}$ /m.

It is also found from Fig. 8 that the electricity cost decreases slowly with the increased insulation thickness. The reason for this phenomenon may attributed to the fact that the thermal conductivity of insulation layer is not small enough to reduce more heat dissipation from tunnel to the airflow. One can expect that the insulation layer with lower thermal conductivity can lead to a remarkable reduction in the electricity cost. In addition, the thermal insulation effect for the tunnel with higher rock temperature may be more significant, which can be proved through the results shown in Fig. 4.

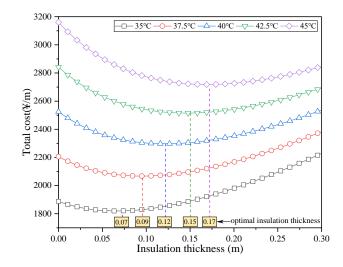


434

435

Fig. 8. Total cost versus the insulation thickness (rock temperature=35 °C).

In order to investigate the influence of rock temperature on the optimal insulation thickness, the total costs under various rock temperatures are also evaluated and shown in Fig. 9. It is found that an increase in the rock temperature can lead to an increase in the total cost. Under the same insulation thickness, the larger the rock temperature, the higher the total cost is. In addition, the optimal insulation thickness is observed to be increased when the rock temperature increases. For instance, the optimal insulation thickness increases from 0.07 m to 0.17m when the rock temperature increases 442 from 35 °C to 45 °C. This may be attributed to the fact that a larger insulation thickness is needed to



443 reduce the cooling load for the case with a higher rock temperature.



445

Fig. 9. Total cost versus the insulation thickness (various rock temperatures)

## 446 *3.3.2. life cycle saving and payback period*

447 The effect of the rock temperature on the life cycle saving is investigated and presented in Fig. 10. It is seen from Fig. 10 (a) that the life cycle saving firstly increases and then decreases when the 448 449 insulation thickness increases, which indicates that there is a maximum life cycle saving for various cases with different rock temperatures. As shown in Fig. 10 (b), the maximum life cycle savings per 450 451 tunnel length are 61.2 ¥/m, 128.2 ¥/m, 213.2 ¥/m, 310.1 ¥/m and 421.8 ¥/m when the rock temperature is 35 °C, 37.5 °C, 40 °C, 42.5 °C and 45 °C, respectively. In addition, it is also observed from Fig. 10 452 453 that the optimal insulation thicknesses corresponding to the maximum life cycle saving under various 454 rock temperatures are consistent with the optimal insulation thickness illustrated in Fig. 9.

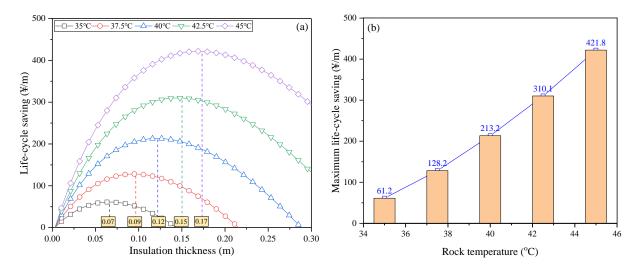




Fig. 10. Life cycle saving versus (a) insulation thickness, (b) rock temperature.

The optimal insulation thickness, maximum life cycle saving and payback period are listed in Table
5. It is found that the optimal insulation thickness and maximum life cycle saving increase from 0.07
m and 61.2 ¥/m to 0.17 m and 421.8 ¥/m when the rock temperature increases from 35 °C to 45 °C.
Interestingly, a shorter payback period is achieved by the increased rock temperature. For instance,
the payback period reduces from 9.96 years to 4.91 years when the rock temperature increases from
35 °C to 45 °C.

463 Table 5 Optimal insulation thickness, maximum life cycle saving and payback period under various rock

464

### temperatures.

Rock temperature (°C)	Optimal insulation thickness	Maximum life cycle saving (¥/m)	Payback period (year)	
	(m)			
35°C	0.07	61.2	9.96	
37.5°C	0.09	128.2	8.3	
40°C	0.12	213.2	6.85	

42.5°C	0.15	310.1	5.72
45°C	0.17	421.8	4.91

### 465 **4.** Conclusions

The transient heat conduction model of insulated underground tunnel is solved by using the separation-of-variable method in this study. The rock temperature and heat flux reduction achieved by applying thermal insulation are investigated. In addition, the optimal thickness, life cycle saving and payback period are discussed based on the life cycle cost method. The main findings are as follows:

(1) The temperature gradient in the insulation layer is considerably larger than that in surrounding
rock. The surplus temperature of original tunnel is 2.5-1.43 times higher than that of tunnel with
insulation layer when Fo increases from 0.28 to 20.22.

474 (2) An increase in the insulation thickness and TCR results in an increase in the heat flux reduction.
475 A decrease in TCIL and tunnel radius have a positive impact on the heat flux reduction. The heat flux
476 reduction is not strongly dependent on the convection coefficient.

477 (3) The total cost firstly decreases and then increases when the insulation thickness increases. The 478 optimal insulation thickness and the maximum life cycle saving increase from 0.07 m to 0.17 m and 479 from 61.2/m to 421.8/m, the payback period reduces from 9.96 year to 4.91 years respectively 480 when the rock temperature increases from 35 °C to 45 °C.

This study presents a preliminary investigation on how the thermal insulation can be used to save the cooling energy and overall system cost in the mining environment in an effective and economic way. More factors, including CO<sub>2</sub> emission and environmental conditions, need to be accounted for in further investigations to gain an in-depth understanding on how to effectively minimise the overall 485 energy consumption in the mining industry, thereby contributing to the carbon neutral target.

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# 490 Appendix:

491 The Eqs. (1)-(4) can be solved by using the variable-separation-method. According to the 492 superposition principle, the solution for above equations can be expressed by Eq. (A1):

493 
$$\begin{cases} T_1(r,t) = T_1^{(1)} + T_1^{(2)} \\ T_2(r,t) = T_2^{(1)} + T_2^{(2)} \end{cases}$$
(A1)

494 where the superscript '(1)' represents the unsteady state solution, and the superscript '(2)' represents 495 the steady-state solution.

### 496 For the steady-state solution, the basic equations are:

497 
$$\frac{\partial^2 T_2^{(2)}}{\partial r^2} + \frac{1}{r} \frac{\partial T_2^{(2)}}{\partial r} = 0$$
(A2)

498 
$$\frac{\partial^2 T_1^{(2)}}{\partial r^2} + \frac{1}{r} \frac{\partial T_1^{(2)}}{\partial r} = 0$$
(A3)

499 
$$T_2^{(2)}(r_3,t) = T_0$$
 (A4)

500 
$$r = r_1: \quad k_1 \frac{\partial T_1^{(2)}}{\partial r} = h \left( T_1^{(2)} - T_a \right)$$
(A5a)

501  

$$r = r_{2}: \begin{cases} T_{1}^{(2)}(r_{2},t) = T_{2}^{(2)}(r_{2},t) \\ k_{1} \frac{\partial T_{1}^{(2)}(r_{2},t)}{\partial r} = k_{2} \frac{\partial T_{2}^{(2)}(r_{2},t)}{\partial r} \end{cases}$$
(A5b)

502 The general solutions for Eq. (A2) and Eq. (A3) are given by Eq. (A6) and Eq. (A7):

503 
$$T_1^{(2)}(r) = C_1 \ln r + D_1$$
 (A6)

504 
$$T_2^{(2)}(r) = C_2 \ln r + D_2$$
 (A7)

Eq. (A6) and Eq. (A7) are substituted into Eq. (A4) and Eq. (A5), the coefficients  $C_1$ ,  $C_2$ ,  $D_1$ ,  $D_2$ 

506 can be obtained by using Eq. (A8):

507

$$\begin{cases} C_{1} = (T_{0} - T_{a}) \left[ \ln r_{2} + \frac{k_{1}}{hr_{1}} - \ln r_{1} - \frac{k_{1}}{k_{2}} (\ln r_{2} - \ln r_{3}) \right]^{-1} \\ C_{2} = \frac{k_{1}}{k_{2}} C_{1} \\ D_{1} = \frac{k_{1}}{hr_{1}} C_{1} - (C_{1} \ln r_{1} - T_{a}) \\ D_{2} = T_{0} - \frac{k_{1}}{k_{2}} C_{1} \ln r_{3} \end{cases}$$
(A8)

508 For the unsteady state solution, the governing equation and the contact boundary condition are also 509 described by Eq. (1) and Eq. (4). The initial condition and boundary condition are given by Eq. (A9) 510 and Eq. (A10):

511 
$$\begin{cases} T_2^{(1)}(r_3,t) = 0\\ k_1 \frac{\partial T_1^{(1)}}{\partial r} \Big|_{r=r_1} = h T_1^{(1)} \Big|_{r=r_1} \end{cases}$$
(A9)

512 
$$\begin{cases} T_1^{(1)}(r,0) = T_0 - T_1^{(2)}(r) \\ T_2^{(1)}(r,0) = T_0 - T_2^{(2)}(r) \end{cases}$$
(A10)

513 The unsteady state equation can be solved by using the separation-of-variables method [40, 43], 514 and the following form of temperature is adopted for both insulation layer and surrounding rock 515 regions (subscripts and superscripts are omitted):

516 
$$T(r,t) = R(r) \cdot \Gamma(t)$$
(A11)

517 Substituting Eq. (A11) into the governing equation, two ordinary differential equations can be 518 obtained in the form of Eq. (A12) and Eq. (A13):

519 
$$\Gamma'(t) + \alpha \beta^2 \Gamma(t) = 0$$
 (A12)

520 
$$r^2 R''(r) + r R'(r) + \beta^2 r^2 R(r) = 0$$
 (A13)

where  $\beta$  is the positive eigenvalue. 521

522 The fundamental solution of Eq. (A12) is given by Eq. (A14):

523 
$$\Gamma(t) = \exp(-\alpha\beta^2)$$
(A14)

The general solution of Eq. (A13) can be expressed by Eq. (A15): 524

525 
$$R(r) = A_1 J_0(\beta r) + B_1 Y_0(\beta r)$$
(A15)

526 Therefore, the general form of the unsteady state solution is derived and described by Eq. (A16):

527 
$$T_i^{(1)}(r,t) = \sum_{n=1}^{\infty} \left[ A_{in} J_0(\beta_{in} r) + B_{in} Y_0(\beta_{in} r) \right] \cdot \exp(-\alpha_i \beta_{in}^2 t)$$
(A16)

528 The coefficients of  $A_{in}$ ,  $B_{in}$ ,  $\beta_{in}$  can be solved from the initial condition and boundary conditions

529 and the contact boundary condition.

530 From the contact boundary condition, Eq. (A17) can be obtained:

531 
$$\alpha_1 \beta_{1n}^2 = \alpha_2 \beta_{2n}^2 = \beta_n^2$$
 (A17)

532 Substituting Eq. (A17) into Eq. (A16), it is obtained as:

533 
$$T_i^{(1)}(r,t) = \sum_{n=1}^{\infty} \left[ A_{in} J_0(\frac{\beta_n}{\sqrt{\alpha_i}}r) + B_{in} Y_0(\frac{\beta_n}{\sqrt{\alpha_i}}r) \right] \cdot \exp(-\beta_n^2 t)$$
(A18)

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