# 1 Experimental application of FRF-based model updating approach to estimate soil mass and

- 2 stiffness mobilised under pile impact tests
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## 29 Abstract

30 The dynamic response of structures in contact with soil is receiving increasing interest and there is a 31 growing need for more accurate models capable of simulating the behaviour of these systems. This is particularly important in the field of offshore wind turbines, where accurate estimates of system 32 33 frequency are needed to avoid resonance, and in the structural health monitoring fields, where 34 accurate reference damage models are used. Previous work has shown that there is significant 35 uncertainty in how to specify mobilised soil stiffness for dynamic soil-pile interaction modelling. 36 Moreover, the contribution of soil mass in dynamic motion is often ignored. This paper applies a 37 finite-element iterative model updating approach previously developed by the authors to two experimental piles to ascertain the mobilised soil stiffness and mass profiles from impact test data. 38 39 The method works by obtaining a frequency response function (FRF) from an impact test performed 40 on a test pile, developing a numerical model of this system, applying initial estimates of soil mass and 41 stiffness, and updating these properties to match the experimental FRF with that generated in the 42 numerical model. A range of elements are investigated including multiple runs of the approach to test 43 repeatability, the influence of different starting estimates for stiffness, the effect of variability in 44 experimental test data, and the influence of the pile length over which masses are distributed. 45 Moreover, potential sources of error are discussed. The method provides reasonably consistent 46 estimates of the soil stiffness and mass acting in the lateral dynamic motion of a given pile tested in 47 this paper. The approach may be useful in the continued improvement of Soil-Structure Interaction 48 (SSI) modelling for dynamic applications.

49 Keywords: Soil Stiffness; Model-Updating; Dynamics; Mass; Winkler; Soil-Structure Interaction

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## 51 **1. Introduction**

52 There is increasing interest in the dynamic response of structures incorporating soil-structure 53 interaction, particularly in the fields of Earthquake [1,2] and Offshore Engineering [3–6] among 54 others. For offshore wind turbines, accurate knowledge of the soil-structure interaction behaviour is 55 paramount to the safe operation of these structures due to the potential for resonance from waves and the spinning rotor, which can exacerbate fatigue. In recent times, the field of vibration-based 56 57 Structural Health Monitoring (SHM), which traditionally focussed on detecting damage in super-58 structural components such as bridge beams [7-9], has begun to focus on damage detection of 59 foundations [10-14]. These recent developments have led to an urgency relating to the need for 60 accurate models capable of encapsulating the behaviour of soil-structure interaction systems.

61 The development of numerical models for structural simulations has been the recourse for design 62 engineers for many years, since it is not possible to experimentally trial every load-case a structure may incur. It is unusual for a developed numerical model of a given structural system to perfectly 63 64 model the behaviour at the first trial, therefore the field of Finite-Element (FE) model updating has 65 focussed on utilising information from the actual structural response to modify the parameters of the 66 numerical model in order to minimise the differences in behaviour between the model and the real 67 system. This is particularly important in the field of structural damage detection where reference 68 numerical models of assets such as bridges are required to benchmark normal operating behaviour. In 69 dynamic modelling fields, model updating approaches have received much attention in recent years 70 [15–23]. Imregun et al. [15] developed a Frequency Response Function (FRF)-based model updating 71 approach and investigated its performance against several barriers for implementation including noisy 72 experimental data and the uniqueness of the updated model when applied to the case of a beam. 73 Experimental noise posed an issue to the accuracy of the method. Nalitolela et al. [16] demonstrated a 74 FRF-based approach using experimental and simulated data, which was based on the addition of 75 artificial stiffness to the structure. A sensitivity procedure was used to update the model parameters. 76 Esfandiari et al. [20] developed a model updating approach to identify the presence of damage by

updating the stiffness and mass of the structure using a FRF-based method applied to a truss model. A similar study by Hwang and Kim [18] focussed on estimating damage severity and location using FRFs for a cantilever beam and a helicopter rotor blade model. Wu et al. [23] presented a FRF-based approach to estimate the mass and stiffness of soil contributing to the lateral dynamic motion of simulated foundation piles, and demonstrated the method using numerically simulated data for typical pile geometries and soil spring stiffness.

This paper is an advancement on work presented by Prendergast and Gavin [6] and Wu et al. [23]. 83 84 Prendergast and Gavin [6] investigated the variation in modelled dynamic response of soil-pile 85 systems through the implementation of different formulations of soil spring stiffness. The various 86 formulations, termed coefficients of subgrade reaction (in static case), require the specification of pile 87 structural and geometric parameters such as Young's modulus (E), second moment of area (I), pile diameter (D) and soil properties including small-strain stiffness ( $E_0$ ) and Poisson's ratio ( $v_s$ ). These 88 89 expressions, originally derived for static applications under specified operational strain, led to 90 significantly varied dynamic responses in the study conducted in [6], both in predicted acceleration 91 magnitude and frequency. This study highlighted the significant uncertainty that persists in the 92 selection of an appropriate subgrade reaction model to transform identical soil and pile properties, as 93 significantly different responses were predicted. The present study applies the FRF-based model 94 updating approach developed by Wu et al. [23] to the experimental case study data of two piles in [6], with a view to estimating the soil mass and stiffness mobilised in the dynamic motion. The FRF of a 95 96 given pile is derived using the input force time-history and the output acceleration-time history from 97 experimental testing, and this is used as the target in the updating method. A numerical beam-Winkler 98 model is developed with an initial soil stiffness profile, estimated using a variety of subgrade reaction 99 formulations and available geotechnical data [6]. This stiffness is applied in the numerical model and 100 the soil mass is initially guessed. The method then updates the stiffness and mass at the soil-structure 101 interface in the beam-Winkler model until the experimental FRF and the numerical FRF generated in 102 the model match within a defined tolerance. The approach aims to reduce the uncertainty in the 103 selection of a soil stiffness profile by enabling a simple model updating approach using a single FRF 104 from the target structure.

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### 106 **2. Numerical modelling of piles**

In this section, the methods adopted to formulate numerical FE models of piles to model theirdynamic responses are described.

#### 110 2.1 Mathematical formulation

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Numerical beam-Winkler models are developed to simulate the behaviour of real test piles, described in Section 4. A FE model from which to obtain the dynamic response of a pile to a lateral impact is modelled in this paper using Euler-Bernoulli beam elements [24] to model the pile, and Winkler spring elements [25,26] to model the soil. Soil mass is incorporated by adding lumped masses to the nodes connecting Winkler spring elements to the pile elements. The global dynamic response is governed by Eq. (1).

117 
$$\mathbf{M}_{\mathbf{G}}\{\ddot{\mathbf{x}}(t)\} + \mathbf{C}_{\mathbf{G}}\{\dot{\mathbf{x}}(t)\} + \mathbf{K}_{\mathbf{G}}\{\mathbf{x}(t)\} = \{\mathbf{P}(t)\}$$
(1a)

where  $\mathbf{M}_{G}$ ,  $\mathbf{C}_{G}$  and  $\mathbf{K}_{G}$  are the ( $N \times N$ ) global mass, damping and stiffness matrices for the pile-soil system; *N* is the total number of degrees of freedom (DOF) and

120 
$$\mathbf{x}(t) = \{x_1(t) \quad x_2(t) \quad \dots \quad x_N(t)\}^T$$
(1b)

$$\dot{\mathbf{x}}(t) = \{\dot{x}_1(t) \quad \dot{x}_2(t) \quad \dots \quad \dot{x}_N(t)\}^T$$
(1c)

$$\ddot{\mathbf{x}}(t) = \{ \ddot{x}_1(t) \quad \ddot{x}_2(t) \quad \dots \quad \ddot{x}_N(t) \}^T$$
(1d)

123 
$$\mathbf{P}(t) = \{P_1(t) \quad P_2(t) \quad \cdots \quad P_N(t)\}^T$$
(1e)

where  $\mathbf{x}(t)$ ,  $\dot{\mathbf{x}}(t)$  and  $\ddot{\mathbf{x}}(t)$  are the displacement, velocity and acceleration of each DOF in the model, 124 125 for each time step. Damping is modelled using Cauchy damping, employing a two-term Rayleigh formulation [27]. The damping ratio used is measured from the experimental signals, see Section 4. 126 The dynamic response is obtained by solving Eq. (1) using the Wilson- $\theta$  integration scheme [28,29]. 127 The natural frequencies and mode shapes of the soil-pile system may be calculated by solving the 128 129 Eigenproblem [27] of the system matrix  $\mathbf{D}_{SYS} = \mathbf{M}_{G}^{-1}\mathbf{K}_{G}$ . Further details on the numerical modelling 130 employed are available in Wu et al. [23]. In this paper, the mass and stiffness matrices for the pile 131 model are derived using the material and geometrical properties of the test piles, described in Section 4. The force vector  $\mathbf{P}(t)$  is populated using the force time-history from a modal hammer impact, 132 133 described in Section 4.

## 134 2.2 Soil stiffness using subgrade reaction approach

135 The present paper is an evolution of work presented by Prendergast and Gavin [6] which assessed the 136 performance of five particular formulations of subgrade reaction in modelling the small-strain

dynamic response of laterally vibrating piles. These models were developed by Biot [30], see Eq. (2),

138 Vesic [31,32], see Eq. (3), Meyerhof and Baike [33,34], see Eq. (4), Klopple and Glock [33–35], see 139 Eq. (5) and Selvadurai [34,35], see Eq. (6). The research in [6] concluded that for the given field conditions and pile parameters considered, the Vesic model (Eq. 3) provided the closest 140 141 approximation to the frequency response of two experimental piles, with deviations of 16.6% and 142 3.9% respectively. However, the analysis highlighted the significant disparity in predicted response 143 depending on which formulation was implemented, and moreover the analysis assumed no soil mass 144 contributed to the dynamic behaviour of the pile-soil system. In this paper, these subgrade reaction 145 models are used to specify the initial stiffness guess in the model-updating approach.

146 
$$k_{s} = \frac{0.95E_{0}}{D(1-v_{s}^{2})} \left[ \frac{E_{0}D^{4}}{(1-v_{s}^{2})EI} \right]^{0.108}$$
(2)

147 
$$k_{s} = \frac{0.65E_{0}}{D(1-v_{s}^{2})} \left[\frac{E_{0}D^{4}}{EI}\right]^{1/12}$$
(3)

148 
$$k_s = \frac{E_0}{D(1 - v_s^2)}$$
(4)

149 
$$k_{s} = \frac{2E_{0}}{D(1+v_{s})}$$
(5)

150 
$$k_s = \frac{0.65}{D} \frac{E_0}{(1 - v_s^2)}$$
(6)

where  $E_0$  is the small-strain Young's modulus of soil (N/m<sup>2</sup>), *D* is the pile diameter (m),  $v_s$  is the Poisson ratio, *E* is the Young's modulus of the pile material (N/m<sup>2</sup>) and *I* is the cross-sectional moment of inertia (m<sup>4</sup>). The  $E_0$  profile for a given site can be estimated using shear wave velocity measurements [36,37], or from correlations to other geotechnical site investigation tests such as Cone Penetration Test (CPT) data [3,38–40]. The method for converting the moduli of subgrade reaction to individual spring moduli is detailed in Prendergast et al. [13].

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## 158 **3** Soil mass and stiffness iterative updating method

A graphical representation of the model updating approach developed by Wu et al. [23] to estimate the soil mass and stiffness acting along a pile is shown in Fig. 1 and an overview of the procedure is summarised herein. An experimental FRF is obtained from an impact test on the pile for which the soil stiffness and mass are sought, using Eq.(7) [6,41,42].

164 
$$H_{a}(\overline{\omega}) = \frac{\ddot{X}(\overline{\omega})}{P(\overline{\omega})}$$
(7)

where  $P(\overline{\omega})$  is the Fourier transform of the input force time-history p(t) from a modal hammer and 165  $\ddot{X}(\overline{\omega})$  is the Fourier transform of the output acceleration time-history  $\ddot{x}(t)$  from an accelerometer. 166 The amplitude of the complex-valued FRF in Eq. (7) is denoted by  $F_a(\overline{\omega}) = |H_a(\overline{\omega})|$ . It is assumed 167 168 that the material and geometry of this pile are known to the user so that a reference beam-Winkler 169 numerical model of the system can be created using the approach in Section 2.1. Using site 170 investigation data such as shear wave measurements and employing a subgrade reaction model such 171 as in Eqs. (2)-(6), soil spring stiffnesses can be applied in the numerical model as the initial educated 172 guess as to the acting soil stiffness in the system. A stiffness weighting,  $w_k$  is initially assumed as 1 173 times this profile. An initial guess of soil mass is postulated from a uniform distribution of mass 174 weightings,  $w_m$  between 0 and 30, to be multiplied by the known pile mass,  $m_p$  and distributed among the sprung pile nodes in the reference numerical model. The information is used to assemble mass, 175 176  $M_G$  and stiffness,  $K_G$  matrices using the approach in Section 2.1. The numerical model also requires 177 an estimate of the damping of the real system and, as a Rayleigh formulation is adopted in the 178 modelling, the damping ratio of the first mode  $\xi_1$  is required. This can be estimated from the 179 experimental time-domain response using the logarithmic decrement technique [43] or through fitting 180 exponential decay functions [44]. This can also be estimated in the frequency domain using the half-181 power bandwidth method [43]. The damping matrix  $C_{G}$  is then formulated as a linear combination of 182  $M_G$  and  $K_G$ , using this specified damping ratio [27]. Once a numerical model employing an initial guess of the soil properties of the real system is developed, one can generate a first estimate numerical 183 184 FRF by applying the force time-history from the experimental test to a node in the numerical model 185 close to the point of application on the real system, and the acceleration response of the system may be calculated by solving Eq. (1). The output acceleration from the node closest to the accelerometer 186 187 on the real system is used in the FRF specification. After the first run of the numerical model, one 188 now has a FRF from the experimental test, and a FRF from the numerical model. A mass ratio is defined as  $r_m = F_{a,EXPT} / F_{a,NUM}$  where  $F_{a,EXPT}$  is the peak amplitude of the experimental acceleration 189 190 FRF and  $F_{a,NUM}$  is the peak amplitude of the calculated numerical FRF. A frequency ratio is defined as 191  $r_{\omega} = f_{NUM} / f_{EXPT}$  where  $f_{NUM}$  is the frequency associated with  $F_{a,NUM}$  and  $f_{EXPT}$  is the frequency associated with  $F_{a,EXPT}$ . The peak information (amplitude and frequency) from both FRFs can be used 192 to obtain  $r_m$ ,  $r_\omega$  and subsequently to calculate  $r_k = r_m \times (r_\omega)^2$ . These values are stored for use later in 193

the linear projection. Two convergence criteria are defined;  $\mathcal{E}_{\omega}$  is the frequency convergence tolerance and  $\mathcal{E}_m$  is mass convergence tolerance. For all experimental trials in this paper, the convergence criteria are set to 1%.

197 For the second run of the iterative method, the soil mass estimate is either increased or decreased depending on the magnitude of  $r_m$  from the initial run. If  $r_m^{(0)} < 1$ , the mass should increase as this was 198 underestimated in the numerical model in the first run. If  $r_m^{(0)} > 1$ , the mass weighting should 199 decrease. The mass weighting is increased or decreased by an arbitrary value of 10 for the second 200 201 guess, with a minimum mass of zero applied (no negative mass). The value '10' is not important, as 202 the actual mass weighting is calculated in later iterations using the two starting estimates from (0) iteration<sup>(0)</sup> and iteration<sup>(1)</sup>. For the stiffness weighting, the second guess is chosen from a uniform 203 distribution of values between 0.7 and 1.3, to be multiplied by the initial soil stiffness profile. Once 204 205 again, the actual value is unimportant, as two starting estimates are required in the iterative approach to allow the system minimise the difference in the FRF peak information and converge on mobilised 206 207 weightings to be applied to the stiffness and mass estimates. Once the second run stiffness and mass 208 weightings are specified (and stored), the system checks if the results of the initial first run are within 209 the defined tolerance, i.e. less than 1% difference in FRF peak amplitudes and frequencies between 210 experimental and numerical FRFs. If they are not, the second guess weightings are applied to the 211 profiles in the numerical model. New  $M_G$ ,  $K_G$  and  $C_G$  matrices are assembled, the force time-history 212 is applied, the output acceleration is calculated, and a new FRF is generated. There now exists two 213 estimates of the FRF of the system, iteration<sup>(0)</sup> and iteration<sup>(1)</sup>. Both of these estimates are used to 214 initiate the linear projection method to calculate further weightings for stiffness and mass towards convergence. These further weightings are updated using the mass ratio,  $r_m$ , and the frequency ratio, 215  $r_{\omega}$ , from the current and previous iterations, and the stiffness ratio defined by  $r_k = r_m \times (r_{\omega})^2$ . The 216 linear projection aims to minimise the difference in FRF peak value and frequency between the 217 218 generated numerical FRF and the target experimental FRF. Once the calculated weightings lead to the 219 generation of a numerical FRF that converges on the experimental FRF, the method terminates and 220 outputs the converged soil stiffness profile and added soil mass.

Due to the tendency for error propagation in automated optimisation processes, some inadmissibility checks and boundary conditions are implemented in the procedure. It is possible for the linear projection method to postulate a negative weighting for stiffness or mass. If this happens, the linear projection method automatically re-calculates the new weighting using the  $j^{th}$  and  $(j-2)^{th}$ ,  $j^{th}$  and  $(j-2)^{th}$ ,  $j^{th}$  and  $(j-2)^{th}$ ,  $j^{th}$  and  $(j-2)^{th}$  iterations until admissible weightings are produced. Should the  $(j-i)^{th}$  iteration reach the first iteration of the method without an admissible weighting being obtained, the new weighting is calculated by multiplying the value of the  $j^{th}$  iteration by a random value between 0.9 and 1.1 (i.e. the current weighting is varied by  $\pm 10\%$ ), then the method continues as normal. Additionally, if convergence is not achieved within (an arbitrary) 15 iterations, the system resets and re-initialises all of the parameters.



Fig. 1. Flow chart of iterative algorithm

#### 233 4 Experimental pile tests

Data from a field test conducted in Prendergast and Gavin [6] is used to test the iterative updating approach developed in Wu et al. [23]. A summary of the field test and information relating to the new analysis is described herein. Lateral vibration tests were conducted on two 0.34 m diameter openended steel piles driven into dense, over-consolidated sand at a quarry in Blessington, southwest of Dublin, Ireland. Prior to testing, both piles were excavated by different amounts to give L/D ratios of 13 and 9 for Pile 1 and 2, respectively, see Fig. 6(a).

- The test quarry has been characterised in detail [45] and used to investigate the performance of a number of model, prototype and full scale foundation concepts over the last number of years [46–49]. A full description of the geotechnical properties of the site can be obtained in [37,45,46,50]. The small-strain stiffness properties of the site, measured using Multi-Channel Analysis of Surface Waves (MASW), see [36], are required for the approach in this paper. The shear wave velocity profile, Fig. 2(a) is used to derive the small-strain Young's modulus profile, Fig. 2(b) by first calculating the
- small-strain shear modulus (G<sub>0</sub>) using  $G_0 = \rho v_s^2$  and  $E_0 = 2G_0(1+v)$ , where  $\rho$  is the soil density
- 247 (kg/m<sup>3</sup>) and v is the small-strain Poisson ratio, taken as 0.1.





Fig. 2. Small-strain soil stiffness data. (a) shear wave velocity measurements, (b) derived  $E_0$  profile.



## Fig. 3. Photo of impact testing on Pile 2

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Each pile was fitted with three accelerometers distributed along the exposed portion of the pile shaft, 253 254 see Fig. 6(a), and these accelerometers were programmed to scan at 1000Hz. Note, only the top accelerometer is used in the procedure while the remaining two accelerometers are used to ensure 255 256 consistency in the data. The test procedure (for a given pile) involved impacting the pile laterally with 257 a PCB Piezotronics 086D50 model sledgehammer-type modal hammer [51](tip mass = 5.5 kg) and 258 measuring the resulting acceleration signal from the accelerometers, see Fig. 3. A number of hammer 259 impacts were undertaken on each pile to investigate repeatability. Each acceleration signal was low-260 pass filtered with a cut-off at 60 Hz to reduce the contribution of higher modes and noise, and a FRF 261 is then generated, which is used as the target data in the numerical analysis to estimate the stiffness 262 and mass contribution of the soil.

The damping ratio is estimated for each impact test by fitting an exponential curve to the peaks of the filtered acceleration signal in the time-domain, see [6,44], and validated using a logarithmic decrement technique [43].

FRFs of velocity and displacement are derived from the acceleration FRF using Eqs. (8) and (9). These FRFs are used to test the convergence of the iterative approach in the sense that if the converged soil mass and stiffness estimates provide a match in  $F_a$ ,  $F_v$  and  $F_d$ , this acts as an additional check to mitigate false positives. Note,  $F_v$  and  $F_d$  are not used directly in the iterative updating approach (see Fig. 1), but only used as a check in the converged model. Note also that these are derived from  $F_a$  because the pile velocity and displacement are not measured in the experiment.

272 
$$F_{\nu}(\overline{\omega}) = \left| \frac{H_a(\overline{\omega})}{i\overline{\omega}} \right| = \frac{F_a(\overline{\omega})}{\overline{\omega}}$$
(8)

273 
$$F_d(\overline{\omega}) = \left| \frac{H_a(\overline{\omega})}{(i\overline{\omega})^2} \right| = \frac{F_a(\overline{\omega})}{\overline{\omega}^2}$$
(9)

where  $\overline{\omega}$  is the variable of excitation. The FRFs for five impact tests conducted on Pile 1 and 2 respectively are shown in Fig. 4 and the data is presented in Table 1. Damping data specified is from the curve fitting approach. Fig. 4(a) shows the frequency content of the force time-histories for the five impacts applied to Pile 1. Fig. 4(b) shows the acceleration FRFs for these five impacts on Pile 1. Fig. 4(c) shows the frequency content of the force time-histories for the five impacts applied to Pile 2. Fig. 4(d) shows the acceleration FRFs for these five impacts on Pile 2. Fig. 4(d) shows the acceleration FRFs for these five impacts on Pile 2.

0	Q	1
4	0	I

#### Table 1 Experimental data

Test	Frequency (Hz)	Damping ratio (%) – curve
		fitting method
P1 T1	20.26	1.77
P1 T2	20.02	1.72
P1 T3	20.02	1.85
P1 T4	20.02	1.77
P1 T5	20.02	1.93
P2 T1	12.21	1.07
P2 T2	12.21	1.24
P2 T3	12.21	1.30
P2 T4	12.21	1.33
P2 T5	12.21	1.30

282 \*P1 = Pile 1, P2 = Pile 2, T = Test No.



283

Fig. 4. Pile impact test data. (a) Frequency content of force time-history for five impact tests T1-T5 on
Pile 1, (b) FRF from each impact test on Pile 1, (c) Frequency content of force time-history for five
impact tests T1-T5 on Pile 2, (d) FRF from each impact test on Pile 2.

Using Eqs. (8) and (9),  $F_a$  can be converted to  $F_v$  and  $F_d$ . Fig. 5 shows the derived  $F_v$  and  $F_d$  from the first impact test conducted on both Pile 1 and Pile 2. These are used as a means to check the converged mass and stiffness weightings at the end of applying the method.



Fig. 5. Frequency Response Functions for the first impact test conducted on Pile 1 and 2. (a) Pile 1  $F_a$ , (b) Pile 1  $F_v$ , (c) Pile 1  $F_d$ , (d) Pile 2  $F_a$ , (e) Pile 2  $F_v$ , (f) Pile 2  $F_d$ .

293

#### 294 **5** Analysis

#### 295 5.1 Numerical modelling of field data

296 Two field piles were experimentally tested, as described in Section 4. Two reference numerical 297 models were developed, shown in Fig. 6(b) and (c) for Pile 1 and 2 respectively, using the procedure 298 described in Section 2.1. Pile 1 contains 72 Euler-Bernoulli beam elements, each of length 0.1m, and 299 46 Winkler spring elements to model the soil. Since Pile 1 was initially excavated from an embedment 300 of 7m to 4.5m, there still exists soil within the pile (as it is an open-ended tube). The level of internal soil (plug) was approximately 2m below the original ground level. This was incorporated in the 301 numerical model as an extra mass, assuming a (packed) density for the internal soil at 2000 kg/m<sup>3</sup>. 302 303 External soil (added) masses are initially set to zero except for the top quarter of the springs, in line 304 with the procedure in [23], due to the fact that an embedded pile impacted laterally at the head will 305 have little modal displacement at depth (Section 5.6 investigates apportioning masses over increasing portions of the piles). The external impact force is applied at a distance of 1m below the pile head, 306 307 close to the point of application on the real system. Pile 2 is modelled similarly to Pile 1, except that 32 Winkler springs are used to model the lesser embedded depth. The soil plug is taken the same as 308

- 309 for Pile 1, as an added mass to a depth of 2 m below the original embedded length (i.e. a soil plug 5m
- 310 long from the pile tip). The impulse force is applied to a node in the model at a distance of 2m below
- 311 the pile head, in accordance to the real situation.



313

Fig. 6. Model schematic (dimensions in mm), (a) experimental pile geometry, (b) numerical schematic
 for Pile 1, (c) numerical schematic for Pile 2

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#### 317 5.2 Example of applying the iterative updating method

An example of running the model is demonstrated in this section and the Pile 1 model with an initial starting soil stiffness estimate using the Biot approach (Eq. 2) is shown. The results are presented in Fig. 7 for the first run of the model (with the random starting estimates for mass weighting), and the final converged values of  $F_a$ , since it is the acceleration FRF that is solely used in the procedure, see Section 3. To show that the method accurately calculates the operating parameters,  $F_v$  and  $F_d$  are also shown as calculated in the model overlain on the derived FRFs from the experimental data. Fig. 7(a) shows the experimental  $F_a$  and the first estimate of the numerical  $F_a$ . Fig. 7(b) and (c) show the same information for  $F_v$  and  $F_d$  respectively. Fig. 7(d) shows the experimental  $F_a$  and the converged numerical  $F_a$ . Fig. 7(e) and (f) show the same information for  $F_v$  and  $F_d$ . A plot of the initial estimate and final converged acceleration signal, used to develop the numerical  $F_a$  is shown in Fig. 8. Fig. 8(a) shows the predicted acceleration for the first iteration overlain on the experimental signal and corresponds to the FRFs shown in Fig. 7(a). Fig. 8(b) shows the final converged numerical acceleration overlain on the experimental signal and corresponds to the FRFs shown in Fig. 7(d). This figure demonstrates how the approach matches the real-measured response in the time-domain.

- The method takes 21 iterations to converge (1 global loop of 15 iterations followed by resetting and 6 further iterations). The values of the parameters of interest (mass and stiffness weightings, ratios and tolerances) for all 21 iterations are reported in Table 2. The method stops when all three tolerances (mass, frequency and inferred stiffness, see Fig. 1) are less than 0.01 (1%). The method estimates that the Biot profile applied to the numerical model should be multiplied by 0.95 and soil mass equating to
- site blot prome applied to the numerical model should be maniplied by 0.95 and son mass equating to
- 6 times the pile mass should be distributed to the top quarter of the pile springs in order to match the
- 338 experimental FRF.



340Fig. 7. Example of running the method for Biot starting profile – Pile 1. (a)  $F_a$  experimental and341numerical iteration 1, (b)  $F_v$  experimental and numerical iteration 1, (c)  $F_d$  experimental and numerical342iteration 1, (d)  $F_a$  experimental and converged numerical, (e)  $F_v$  experimental and converged343numerical, (f)  $F_d$  experimental and converged numerical.



Fig. 8 Experimental and predicted accelerations – Pile 1. (a) Iteration 1 of the method, (b) Final
 iteration (21) of the method

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## Table 2 Parameters during iterative process

Global loop	Iteration	Wm	Wk	r <sub>m</sub>	rω	$\mathbf{r}_{\mathbf{k}}$	Tol <sub>m</sub>	$\mathrm{Tol}_\omega$	$\mathrm{Tol}_k$
0	1	24.442	1.000	6.077	0.621	2.342	5.077	0.379	1.342
0	2	14.442	1.079	2.711	0.802	1.744	1.711	0.198	0.744
0	3	9.359	1.178	1.366	0.965	1.273	0.366	0.035	0.273
0	4	7.974	1.235	1.087	1.025	1.142	0.087	0.025	0.142
0	5	7.542	1.297	0.988	1.055	1.100	0.012	0.055	0.100
0	6	7.592	1.444	0.919	1.089	1.089	0.081	0.089	0.089
0	7	7.533	2.658	0.651	1.273	1.055	0.349	0.273	0.055
0	8	7.611	4.588	0.564	1.400	1.107	0.436	0.400	0.107
0	9	7.220	0.647	1.633	0.834	1.135	0.633	0.166	0.135
0	10	7.451	19.388	0.504	1.641	1.358	0.496	0.641	0.358
0	11	9.023	36.318	0.499	1.718	1.472	0.501	0.718	0.472
0	12	10.250	97.502	0.495	1.821	1.641	0.505	0.821	0.641
0	13	11.206	699.121	0.496	1.970	1.923	0.504	0.970	0.923
0	14	12.402	2236.178	0.498	2.029	2.049	0.502	1.029	1.049

0	15	13.486	10241.427	0.500	2.083	2.169	0.500	1.083	1.169
1	1	2.926	1.000	0.650	1.131	0.831	0.350	0.131	0.169
1	2	12.926	0.867	2.780	0.763	1.620	1.780	0.237	0.620
1	3	4.571	0.971	0.800	1.065	0.908	0.200	0.065	0.092
1	4	5.413	0.958	0.910	1.029	0.964	0.090	0.029	0.036
1	5	6.108	0.949	1.016	1.001	1.018	0.016	0.001	0.018
1	6	6.001	0.952	0.998	1.006	1.009	0.002	0.006	0.009

## 349 5.3 Converged results for different starting stiffness profiles

350 In this section, the results of applying each of the five subgrade models (Eqs. 2-6) as the initial 351 starting estimate are trialled for Pile 1 and Pile 2. Each model is run one time, and the results of the 352 converged mass and stiffness weightings for each stiffness profile and both piles are shown in Table 353 3. It is important to note that the converged stiffness weighting should be different for each model, as 354 this is multiplied by the initial profile (Biot, Vesic, etc.) to obtain the converged soil stiffness profile. 355 The mass weighting should be relatively consistent between runs, since this is multiplied by the 356 constant that is the pile mass (for a given pile). In Table 3, it can be seen that for Pile 1, a relatively consistent estimate of the mass weighting is obtained from each model. The converged mass 357 weighting for Pile 2 is a little more variable, though still reasonably consistent. 358

# Table 3 Converged stiffness and mass weightings for one run of updating method for each subgrade reaction model – Pile 1 & 2

	Pl	ILE 1	PILE 2		
Model	Converged <i>w</i> <sub>k</sub>	Converged <i>w</i> <sub>m</sub>	Converged <i>w<sub>k</sub></i>	Converged <i>w</i> <sub>m</sub>	
Biot	0.874	5.633	1.779	18.176	
Vesic	1.266	5.838	2.399	17.790	
Meyerhof & Baike	0.703	5.984	1.369	18.796	
Klopple & Glock	0.377	5.858	0.847	20.915	
Selvadurai	0.994	5.691	2.101	18.813	

<sup>361</sup> 

362 As mentioned above, it is expected the converged stiffness weightings be different for each model, as 363 this is multiplied by the specified soil stiffness profile to obtain the converged stiffness profile. This is best demonstrated as in Fig. 9, which shows the starting and converged stiffness profiles with depth for each of the five subgrade reaction models for Pile 1. The stiffness is shown in terms of spring stiffness units (N/m). Fig. 9(a) shows the initial spring stiffness profiles (the markers show the individual springs) as derived from the site data in Fig. 2(b) using each subgrade model (Eqs. 2-6). Fig. 9(b) shows the results of multiplying each of these profiles by the associated converged stiffness weighting for Pile 1 in Table 3. This plot demonstrates visually how the profiles converge toward one enother to establish the actine soil stiffness for Pile 1

another to establish the *acting* soil stiffness for Pile 1.



371

Fig. 9. Converged stiffness profiles after one run of each model - Pile 1. (a) Original stiffness profiles
 from each subgrade reaction formulation, (b) Converged weighted stiffness profile after one run of
 each model.

## 375 5.4 Multiple runs for a given stiffness profile

The previous section presents the results of running each model once until convergence is achieved. However, since each run begins with effectively random starting estimates (between 0 and 30 for the mass weighting for the first run, and between 0.7 and 1.3 for the stiffness weighting for the second run), it is of interest to assess repeatability between multiple runs of a given model. Pile 1 with an initial stiffness profile defined by the Biot model (Eq. 2) is run five times until converged mass and

stiffness weightings are obtained. Fig. 10 shows the path of each weighting toward convergence for 381 382 each run, Fig. 10(a) for the mass weightings and Fig. 10(b) for the stiffness weightings. Each run (R1-383 R5) takes a different number of iterations to converge. R1 takes 20 iteration to converge and ends with  $w_m = 5.63$  and  $w_k = 0.87$ . R2 takes only 4 iterations to converge and ends with  $w_m = 5.99$  and  $w_k$ 384 385 = 0.94. R3 converges after 6 iterations with  $w_m = 5.89$  and  $w_k = 0.93$ . R4 takes 5 iterations and converges with  $w_m = 6.02$  and  $w_k = 0.94$ . Finally, R5 converges after 36 iterations with  $w_m = 6.14$  and 386  $w_k = 0.98$ . Note also that the system resets if convergence is not achieved in 15 iterations, where all 387 388 the parameters are reinitialised and the procedure starts over, see Fig. 1. The converged mass and 389 stiffness weightings do vary a little between runs however in the context of obtaining stiffness 390 information for geotechnical applications, they are reasonably consistent. Some of the reasons for the 391 difference in the converged values is discussed in Section 5.7.



392



The results for the same analysis on Pile 2 is summarised in Table 4. The mass and stiffness weightings are reasonably consistent between runs for this pile with the Biot model.

Table 4 Results of 5 runs of Biot model – Pile 2

Analysis run No.	w <sub>m</sub> converged	w <sub>k</sub> converged	Iterations
1	18.18	1.78	4
2	18.78	1.85	6
3	18.97	1.86	267
4	18.00	1.78	7
5	19.09	1.87	517

400

### 401 5.5 Consistency between different experimental impact tests

402 Until now, only one set of experimental data from each pile, namely  $F_a$  from 1 impact test (P1 T1 and 403 P2 T1 Table 1) has been considered. In this section, the ability for the method to calculate consistent 404 mobilised stiffness and mass weightings from a number of impact tests conducted on both Piles 1 and 405 2 is evaluated. The target FRFs for five impact tests are shown in Fig. 4. The method is run one time 406 for each of the starting soil stiffness models (Eqs. 2-6), for each of the five impact tests conducted on 407 both piles (Table 1), resulting in a total of 50 runs. Table 5 shows the values of the converged stiffness 408 and mass weightings from each run for Pile 1 and Table 6 shows the results for Pile 2.

#### 409

#### Table 5 Pile 1 Analysis of five impact tests

Impact Test No.	VE	SIC	SELVA	DURAI	BI	ОТ	MEYE	RHOF	KLO	PPLE
	$W_m$	$W_k$								
1	5.838	1.266	5.691	0.994	5.633	0.874	5.984	0.703	5.858	0.377
2	4.570	1.007	5.112	0.939	5.233	0.847	4.767	0.569	4.731	0.312
3	3.861	0.963	3.661	0.750	3.774	0.680	3.661	0.484	3.684	0.270
4	4.179	1.000	4.058	0.775	4.125	0.704	4.110	0.518	4.006	0.276
5	3.316	0.874	3.152	0.684	3.391	0.638	3.237	0.451	3.265	0.252

# 410

411

Table 6 Pile 2 Analysis of five impact tests

Impact Test No.	VES	SIC	SELVAI	DURAI	BIC	)T	MEYEI	RHOF	KLOP	PLE
	Wm	$W_k$	$W_m$	$W_k$	$W_m$	$W_k$	$W_m$	$W_k$	$W_m$	$W_k$
1	17.790	2.399	18.813	2.101	18.176	1.779	18.796	1.369	20.915	0.847
2	16.763	2.375	16.203	1.881	18.413	1.829	17.606	1.323	17.862	0.738
3	17.959	2.486	17.918	2.045	17.771	1.790	17.991	1.334	17.649	0.710
4	15.771	2.366	15.625	1.905	15.371	1.670	14.932	1.199	16.107	0.708
5	15.220	2.030	15.376	1.686	14.249	1.379	15.757	1.104	15.633	0.622

413 Observing Table 5 and 6, the data from different impact tests lead to slightly different estimates of 414 converged mass weightings in each case, for both piles. It is noteworthy that for a given impact test, 415 the converged mass weightings for each of the soil stiffness models are relatively consistent for a 416 given pile. There are two potential reasons for this, (i) the mass weighting is very sensitive to the 417 quality of  $F_a$  and any variations in this strongly affect the converged mass weighting, or (ii) depending 418 on the magnitude of the impact applied in each case, different amounts of mass may have been 419 mobilised in the soil surrounding the pile. Converged stiffness weightings for a given soil profile also 420 vary somewhat between impact tests. Further potential reasons for these differences are discussed in Section 5.7. 421

## 422 5.6 Influence of changing the active length over which masses are apportioned

423 All previous analyses consider the added soil masses apportioned to the top quarter of the springs in each model, as an approximate estimate for the mobilised mass of soil contributing to the first mode 424 425 of vibration of each system. In reality, there will be some depth over which the soil mass will be 426 effectively mobilised, due to the nature of the pile head bending when impacted. The active length, or 427 effective depth of a pile, is the length beyond which further increases in pile length do not have any 428 additional influence on pile head displacements, or rotations (or frequency) [52]. Quantifying the 429 active length is an area of much uncertainty and previous studies have suggested several formulations for this parameter, which vary depending on the constraints applied to the pile head, the pile rigidity, 430 and the nature of applied loading [53-58]. In this section, the influence of distributing masses over 431 432 different lengths of a pile on the converged stiffness and mass weightings is studied. Active lengths equating to 25%, 50%, 75% and 100% of the embedded pile length are considered. Pile 1 impact test 433 434 1 (P1 T1, Table 1) is used as the test case and a Biot soil profile is adopted as the initial soil stiffness estimate. Each model is run five times for a given mass length distribution, and the results are 435 436 presented in Table 7 as the average  $\pm$  standard deviation of converged mass and stiffness weightings, 437 for each mass distribution case.

438

Table 7 Influence of mass length distribution on converged weightings

Masses distributed over	Average $w_m \pm$ Standard	Average $w_k \pm$ Standard
length, L (L <sub>p</sub> = pile length)	deviation	deviation
L=0.25L <sub>p</sub>	5.83±0.27	0.92±0.05
L=0.5L <sub>p</sub>	11.54±0.39	0.92±0.04
L=0.75L <sub>p</sub>	18.10±0.70	0.97±0.04
L=L <sub>p</sub>	23.90±0.44	0.97±0.02

440 Increasing the length over which masses are apportioned has limited influence on the converged 441 stiffness weighting, with these values remaining sufficiently consistent for each case, considering the 442 nominal errors present due to the natural variability in the algorithm convergence process. However, 443 the converged mass weighting increases proportionally to the increase in mass distribution length, 444 changing from  $w_m$ =5.83 for masses distributed over 25% of the pile embedment to  $w_m$ =23.90 for 445 masses distributed along the entire embedded depth. In the procedure to add point masses to the pile, 446 the mass weighting is multiplied by a fixed 'added mass', which is the pile mass, and this is then 447 divided equally among the 'active spring nodes', namely the nodes with non-zero added masses. So, for the first case, a weighting of 5.83 is multiplied by the pile mass and divided among 12 springs (a 448 449 quarter of the 46 springs), giving  $\approx 0.5$  times the pile mass added to each spring. For the last case, a 450 pile weighting of 23.90 is multiplied by the pile mass and divided among all 46 springs, again giving 451  $\approx 0.5$  times the pile mass added to each spring. Therefore, when one normalises the converged weighting to the number of springs with non-zero added masses, the added point mass at each spring 452 453 is approximately the same.

454 This finding highlights that no matter how many springs are specified to attach masses, the added point mass at each spring will be approximately the same. This result may seem counterintuitive as 455 456 the global mass added increases with the number of active springs, and suggests that the approach is 457 therefore very sensitive to the specified active length by the user. However, this result may be understood by observing the influence of added point masses on the  $F_a$  peak height for the first mode 458 459 of the pile. Herein, the model for Pile 1 with a Biot stiffness profile subjected to an impact test is 460 shown for the case where fixed point masses are added sequentially to the springs starting from ground level. The first run contains no added soil mass, the second run has one added mass, etc., until 461 462 all the springs contain the same added point mass. With the increasing number of added masses, the 463 FRF  $F_a$  peak height ( $F_{a,max}$ ) decreases logarithmically, see Fig. 11. It is noteworthy that the peak 464 heights,  $F_{a,max}$  for the cases with masses added to 12 springs (L/L<sub>p</sub>=0.25) and masses added to 46 465 springs (L/L<sub>p</sub>=1) do not vary significantly, which explains why the result appears insensitive to the length over which masses are added. Note, to further investigate this influence would require 466 observing higher modes of vibration, which would be influenced strongly by a given mass 467 468 distribution. However, this is beyond the scope of the present study. It is recommended that potential 469 users of the method specify an active length using the most applicable approach available.



471 Fig. 11 Influence of increasing the number of added masses along the pile on the  $F_a$  peak value

### 473 5.7 Sources of error in the method

The iterative model updating approach presented in [23] was developed and validated using 474 475 numerically simulated data of piles. Application of the approach to real experimental data has 476 unearthed some issues. Variability and noise in experimental data inevitably affects the quality of 477 results. One of the key issues may relate to the time-length of the signals available for the 478 experimental analysis. The impact tests conducted on both piles contained 3 seconds of acceleration 479 data. The impact of this is investigated in Figs. 12 and 13. Fig. 12(a) shows how the FRF  $F_a$  peak amplitude varies for different mass and stiffness weightings applied to the numerical model of Pile 1 480 481 with a Biot soil stiffness profile. The surface plot in Fig. 12(a) is generated using time signals of 482 length T=200s, the same as the analyses conducted throughout this paper. Also shown as a horizontal 483 plane in grey is the peak amplitude of the experimental  $F_a$  as measured in the first impact test on Pile 1. An immediately obvious trait is that the numerical  $F_a$  peak amplitude is affected by changes in both 484 485 mass and stiffness weighting, which deviates significantly from the theory of how single-degree-of-486 freedom (SDOF) models should behave, see [23]. The curve along which both the experimental and 487 numerical planes intersect provides the solution combinations  $\{w_m, w_k\}$ , which lead to the same  $F_a$ 488 peak amplitude in the numerical model as in the experimental data. It is important to note that the 489 other criterion of matching the frequency is required in the iterative procedure, but not shown in these 490 plots. This explains why the procedure always converges on broadly similar values for a given situation, and not a large range, as would be the case if the  $F_a$  peak alone were sought. Fig. 12(b) 491 shows the same information as Fig. 12(a) but this time for the FRF  $F_d$  peak amplitude. The 492 493 experimental data (horizontal grey plane) is the  $F_d$  peak amplitude derived from the experimental  $F_a$ 

- 494 using Eq. (9). Once again there is an intersection curve of  $\{w_m, w_k\}$  combinations that enables the 495 numerical model have the same  $F_d$  as the experiment. The influence of time on signal quality is
- 496 investigated in Figs. 12(c) and (d), where an acceleration time series of length T=3s is used for each
- 497 run. The difference between the surface plots in (a) and (b) to those in (c) and (d) is best demonstrated
- 498 in the contour plots shown in Fig. 12(e) and (f). The result of using a time series of length T=200s for
- the analyses is shown by the smoothness of the solid contour lines in parts (e) and (f). Reducing the
- 500 time series to T=3s (in line with the experimental data) leads to a more jagged contour plot, denoted
- 501 by the dashed lines in (e) and (f). This roughness in the peak  $F_a$  amplitude infers that for convergence
- 502 to be achieved between the 'rough' experimental  $F_a$  and the 'smooth' numerical  $F_a$  some errors are
- 503 introduced. For Pile 1 with a Biot stiffness profile, this is quite minor, however Fig. 13 shows the
- same information for Pile 2, which is significantly affected by signal length issues.





507 Fig. 12. Influence of signal length on FRF peak height for different mass and stiffness weightings for Pile 1 – Biot model. (a) variation of peak amplitude of  $F_a$  with  $w_m$  and  $w_k$  compared to experimental  $F_a$ 508 509 Impact Test 1 - T = 200s, (b) variation of peak amplitude of  $F_d$  with  $w_m$  and  $w_k$  compared to 510 experimental  $F_d$  (derived) Impact Test 1 – T=200s, (c) variation of peak amplitude of  $F_a$  with  $w_m$  and 511  $w_k$  compared to experimental  $F_a$  Impact Test 1 - T = 3s, (b) variation of peak amplitude of  $F_d$  with  $w_m$ 512 and  $w_k$  compared to experimental  $F_d$  (derived) Impact Test 1 - T = 3s, (e) contour plot of peak 513 amplitude of  $F_a$  with  $w_m$  and  $w_k$  for both T=200s and T=3s runs, (d) contour plot of peak amplitude of 514  $F_d$  with  $w_m$  and  $w_k$  for both T=200s and T=3s runs.

Fig. 13 shows the results for Pile 2 with a Biot soil stiffness profile. Fig. 13(a) shows a surface plot of the  $F_a$  peak amplitude and how it varies with mass and stiffness weightings. Fig. 13(b) shows this information for the  $F_d$  peak amplitude. Also shown as a horizontal grey plane is the experimental  $F_a$ peak amplitude in (a) and derived  $F_d$  peak amplitude in (b) from the first impact test on Pile 2. The

520 smooth surface plots in (a) and (b) are derived from analysis of signals that are T=200 s long. Figs. 521 13(c) and (d) show the same information as (a) and (b) respectively, but are generated from time signals that are T=3s long. For this case there is a substantial decrease in the smoothness of each plot, 522 523 which highlights the potential errors that are introduced by the use of short time signals in the 524 experimental data analysis. The results from the four surface plots in Figs. 13(a)-(d) are shown as 525 contour plots in (e) and (f), where the solid contours are generated from T=200s signals and the 526 jagged contours from T=3s. This highlights that use of the short experimental signals is a potential 527 source of model error, which may be significant. This may account for some of the difference in calculated stiffness weightings between Pile 1 and 2. Note, all of the analyses in the previous sections 528 529 used T=200s for the numerical modelling while the experimental signals contained only 3 seconds of 530 data.





533 Fig. 13. Influence of signal length on FRF peak height for different mass and stiffness weightings for Pile 2 – Biot model. (a) variation of peak amplitude of  $F_a$  with  $w_m$  and  $w_k$  compared to experimental  $F_a$ 534 535 Impact Test 1 - T = 200s, (b) variation of peak amplitude of  $F_d$  with  $w_m$  and  $w_k$  compared to 536 experimental  $F_d$  (derived) Impact Test 1 – T=200s, (c) variation of peak amplitude of  $F_a$  with  $w_m$  and 537  $w_k$  compared to experimental  $F_a$  Impact Test 1 - T = 3s, (b) variation of peak amplitude of  $F_d$  with  $w_m$ 538 and  $w_k$  compared to experimental  $F_d$  (derived) Impact Test 1 - T = 3s, (e) contour plot of peak 539 amplitude of  $F_a$  with  $w_m$  and  $w_k$  for both T=200s and T=3s runs, (d) contour plot of peak amplitude of 540  $F_d$  with  $w_m$  and  $w_k$  for both T=200s and T=3s runs.

541 While the short time-length of the processed signals may be the largest source of error, an additional 542 source of error arises from the experimental impact testing. Each pile is an open-ended steel cylinder 543 and, when subjected to impacts from a modal hammer, this induces an in-plane excitation in the pile 544 annulus. This in-plane excitation manifests as a high-frequency pollution in the bending signal. Prior 545 to transforming the time-signal to a FRF, the signal is low-pass filtered to remove the contribution of 546 this noise [6,41]. This process will have some influence on the quality and nature of the FRF.

547 Further sources of error might arise due to the stepped nature of the available soil stiffness (E) data from the multi-channel analysis of surface waves. Any errors here may be exacerbated in the 548 549 procedure, which uses a single stiffness weighting for the entire profile depth. Moreover, since Pile 2 550 has less embedded depth than Pile 1, any errors in this profile will be exacerbated further. It should be 551 noted that the same E profile is used for both piles, as this is in effect an average profile for the test 552 site, so some errors can be expected as to the actual acting magnitudes at each depth. In terms of the 553 reference numerical models developed, there is some question over the mass density of the internal 554 plugged soil in each pile, which had to be estimated for the purposes of this paper. Additionally, the 555 numerical method involves simplifying the pile to a 1D beam-Winkler system, which may deviate in 556 behaviour from the real continuous pile system. Due to numerical constraints in iterative analyses of 557 this nature, it is infeasible to use a full 3D model as it would be computationally too expensive.

#### 558 6. Conclusion

In this paper, the application of a finite-element model updating approach to estimating the mobilised soil stiffness and mass in laterally impacted piles is studied. The reason behind the development of this method is due to the ongoing uncertainty surrounding the specification of soil-structure interaction stiffness in pile-soil interaction. Moreover, any contribution of soil mass is typically ignored. The method, which was previously derived and applied to simulated data, is demonstrated using experimental pile data in this paper.

565 Impact tests are performed on two piles with varying L/D ratios to derive frequency response 566 functions, which are used as the target in an algorithm to estimate the mobilised soil stiffness and 567 mass. Five subgrade reaction formulations are used to specify the initial starting stiffness. The 568 analysis updates the soil stiffness and mass in a numerical model of the pile to converge on the 569 experimental FRF. For the case where each of the five subgrade reaction models are used, the method 570 converges on broadly similar added mass weightings and the converged stiffness profiles are 571 relatively similar. This is better for Pile 1 than for Pile 2, which exhibits more variability (less 572 embedded depth leads to more errors potentially). For a given impact test, the effect of running the 573 model multiple times is studied to ascertain if significant variability exists between different runs. The 574 results do vary a little, due to the random nature of the starting estimates for mass in the first iteration 575 and stiffness in the second iteration, though the converged values are broadly similar for each trial. 576 More variability is evident when different impact tests are used as the target FRF for each case. In 577 general, for a given impact test, the converged mass weighting for each subgrade reaction model is 578 relatively similar for a given pile. However, the difference between the converged weightings for the 579 different impact tests warrants some discussion. Experimental errors in the FRF peak height is most 580 likely the reason for this variation, though there is potentially some influence from the amount of 581 mobilised mass surrounding the pile as a result of the intensity of a given impact from the modal 582 hammer. Additionally, the influence of the active depth over which masses are distributed is also 583 investigated and it is shown for the conditions tested that masses distributed over a length beyond 584 20% of the embedment have limited further influence on the first mode of vibration. However, the 585 effect on higher modes was not evaluated and would require further study. Finally, the sources of 586 error due to time-length of signals is studied with a view to shedding some light on the importance of 587 accurate experimental data. It is recommended that future studies use longer time signals for the experimental data than those available in the present study to mitigate against these signal-processing 588 589 related issues. Short time signals lead to poorly spaced frequency vectors in the FRF, which may 590 strongly influence the converged results.

Aside from some issues, the method was applied with relative success in this paper, and shows that a simple impact test may be useful to obtain better estimates of the mobilised soil-structure interaction stiffnesses and masses acting in the small-strain dynamic soil-pile behaviour. The research may be useful for the development of more accurate damage quantification models for SSI applications or in the growing offshore monopile fields.

596 Future work will investigate extension of the approach to use of multiple vibration modes to provide 597 further insight into the behaviour and, to potentially enable depth-dependant weightings be obtained. 598 The latter may be more useful for cases where large-strain deformations are experienced at pile heads 599 relative to at-depth, thereby enabling calculation of the mobilised strain-dependant stiffness at the pile 600 head. Furthermore, expansion of the approach to different types of foundation structures such as 601 shallow pads or suction caissons should form part of future work. It should be noted that the approach 602 in this paper uses an impact from a modal hammer to excite a structure; therefore there are some 603 limitations of this approach. Large-diameter monopiles may not be sufficiently excited by impact 604 from a modal hammer in order to obtain reliable FRFs. Moreover, highly damped systems suffer the same issues. Expansion of the approach to these types of systems may require investigation ofdifferent excitation sources to generate FRFs.

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