

# Promising emission behavior in Pr<sup>3+</sup> / In selenide-chalcogenide-glass small-core step index fiber (SIF)

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## Abstract

Selenide-chalcogenide glass, small-core, step-index fiber (SIF), core-doped with Pr<sup>3+</sup>: 9.51 x 10<sup>24</sup> ions m<sup>-3</sup> (500 ppmw) is fabricated for the first time with indium to help solubilize Pr<sup>3+</sup>. Core diameters of 20 or 40 μm are confirmed using scanning electron microscopy and near-field imaging; fibre numerical aperture is ~0.4. Optical loss is ≥4.9 dB m<sup>-1</sup> across the 3-9 μm mid-infrared (MIR) spectral range. On pumping at 1.55 μm or 2.013 μm, the SIFs give broad MIR emission across 3.5-6 μm assigned to <sup>3</sup>H<sub>6</sub>→<sup>3</sup>H<sub>5</sub> and <sup>3</sup>H<sub>5</sub>→<sup>3</sup>H<sub>4</sub>. The Pr<sup>3+</sup> emission-lifetime at 4.7 μm decreases from *bulk*-glass (10.1 ±0.3 ms), to *intermediately* processed fiber (8.10 ±0.5 ms) to SIF (7.1 ±0.5 ms) induced by the processing. On end-pumping SIFs at 2.013 μm, the output pump-power and emission intensity at 4.7 μm became sub-linear and super-linear, respectively, suggesting MIR excited-state saturation is occurring.

## KEYWORDS

mid-infrared, chalcogenide glass, rare earth ions, emission, emission-lifetime, excited-state saturation.

## 1. Introduction

The mid-infrared (MIR) spectral region spans the 3-50  $\mu\text{m}$  spectral range [1]. MIR sources include blackbodies (*e.g.* *Globar*<sup>®</sup>), but these display low brightness. In contrast, MIR quantum cascade lasers, an emerging technology, and MIR OPOs (*optical parametric oscillators*) and gas (HeNe, CO, CO<sub>2</sub>) lasers, which are mature technologies, exhibit high brightness. New MIR fiber narrow-line lasers are being developed along two strands: (i) nonlinear conversion, using stimulated Raman gain scattering [2] and (ii) direct-emission, rare earth ion (RE)-doped [3-6], - the topic of this study. MIR RE-doped fiber lasers have not yet been demonstrated at  $\geq 4 \mu\text{m}$  [7], yet potentially offer advantages of compactness, high quantum efficiency, high brightness, excellent beam quality, ability to be pulsed, and greater reliability over gas lasers. MIR RE-doped fiber-lasers have prospective applications in providing new wavelengths for cutting/welding of soft materials, including polymers and in medical fiber-laser-surgery of human-tissue, and as narrow-line MIR molecular sensors [8]. Importantly, MIR RE-doped fiber lasers are potential pumps for MIR fiber supercontinuum (SC) laser sources to achieve all-fiber solutions for portable, real-time, broadband MIR molecular sensing, for instance for early diagnosis of cancer [8-15].

Chalcogenide glasses are promising RE hosts for MIR fiber-lasing due to their low phonon energy, large refractive indices hence large RE absorption/emission cross sections, long fluorescent-decay lifetimes [6] and potential for low optical-loss fiber-fabrication [see 16 and refs. therein]. Selenide-chalcogenide glasses, selected here, retain longer-wavelength near-infrared (NIR) transparency for pumping with commonly available lasers.

RE solubility is poor in binary chalcogenide glasses *e.g.* As<sub>40</sub>Se<sub>60</sub> [17]. A Ga-solubilizer is commonly added to aid RE solubility in chalcogenide glasses [18] based on Ge-As/Sb-S/Se [18-23]. The RE solubility is considered to be enhanced by means of local [Ga-(S/Se)-RE] chemical-complexing [18]. We have made bulk selenide glasses doped with up to 6000 ppmw RE in the presence of Ga [24]. Also, with a Ga solubilizer, we have fabricated 500 ppmw Pr<sup>3+</sup>/Ga small-core selenide-glass SIF (step-index fiber) which for the first time exhibited the same emission-lifetime at 4.7  $\mu\text{m}$  as its parent bulk-glass - 7.8 ms [23]. This result implied that glass homogeneity was retained during the SIF glass-fabrication processing which was verified by a painstaking study to image and analyse the SIF small-core, and core/cladding interface, using high resolution transmission electron microscopy [23].

An In solubilizer was used to make a Pr<sup>3+</sup>/In bulk selenide-chalcogenide glasses and large-core SIF [25-27]. In is in the same Group of the Periodic Table, as Ga, but heavier. Thus, potentially an [In-(S/Se)-RE] chemical-complex could offer a local, lower phonon energy environment to improve RE radiative-efficiency compared to the [Ga-(S/Se)-RE] chemical-complex, notwithstanding that RE 4f inner level transitions tend to be shielded from the local chemical environment. Indeed, on pumping at 1.55  $\mu\text{m}$ , we found that the Pr<sup>3+</sup> emission across 3.5-6  $\mu\text{m}$ , and emission at 4.7  $\mu\text{m}$ , were of greater intensity, and longer lifetime, respectively, in the Pr<sup>3+</sup>/In bulk glass than in the Pr<sup>3+</sup>/Ga bulk glass [25]. Further work on producing high

purity  $\text{Pr}^{3+}$  doped multi-component chalcogenide glasses can be found in [28-30], including the use of indium iodide to incorporate indium into Ge-As-Se-In-I [29] and Ge-Sb-Se-In-I [30] glasses.

Here, we report on a study of the emission behavior and fabrication for the first time of  $\text{Pr}^{3+}/\text{In}$  small-core SIFs based on selenide-chalcogenide glasses. The concentration of In solubilizer in the  $\text{Pr}^{3+}$  core was fixed at 1 atomic % (at. %). The SIF core-diameters were 20  $\mu\text{m}$  and 40  $\mu\text{m}$ , with a NA (numerical aperture) of  $\sim 0.4$  (estimated from [23]) and V-parameter  $\sim 5.0$ , giving multi-moded behavior at 4-6  $\mu\text{m}$  wavelength, corresponding to the  ${}^3\text{H}_6 \rightarrow {}^3\text{H}_5$  and  ${}^3\text{H}_5 \rightarrow {}^3\text{H}_4$  emissions and potential lasing emissions.

On pumping the SIF at 1.55  $\mu\text{m}$  or 2.013  $\mu\text{m}$ , broad MIR emission at 3.5-6  $\mu\text{m}$  was observed. The  $\text{Pr}^{3+}$  emission-lifetime at 4.7  $\mu\text{m}$  decreased from *bulk*-glass (10.1  $\pm 0.3$  ms), to *intermediately* processed fiber (8.10  $\pm 0.5$  ms) to SIF (7.1  $\pm 0.5$  ms); the decrease may have been induced by the processing. On end-pumping SIFs at 2.013  $\mu\text{m}$ , the output pump-power collected from the opposite fiber-end to that pumped and emission intensity at 4.7  $\mu\text{m}$  collected from same fiber-end to that pumped, became sub-linear and super-linear, respectively, suggesting MIR excited-state saturation is occurring for the first time in a chalcogenide-glass fiber.

## 2. Experimental

### 2.1 Bulk glass preparation

#### 2.1.1 Cladding-glass boule (for extrusion to tube)

A Ge-As-In-Se-S cladding-glass boule (see Fig.1) was prepared. Ge (5n Materion), As (7n5 Furakawa Denshi, prior heat-treated at  $10^{-3}$  Pa), indium (6n5 Alfa Aesar), Se (5n Materion, prior heat-treated at  $10^{-3}$  Pa) and S (5N, prior boiled under  $10^{-3}$  Pa for 5 mins) were batched inside a glovebox (MBraun:  $< 0.1$  ppm  $\text{H}_2\text{O}$  and  $< 0.1$  ppm  $\text{O}_2$ ) and melted 850°C/8h in a silica glass ampoule (prior air-baked then vacuum-baked, each 1000°C/6h) before being quenched and annealed. 3 atomic% (at%) Se substitution by S manifested useful contrast in optical microscopic and SEM imaging; 1 at% Se replaced by S in chalcogenide glasses has been reported to reduce refractive index by 0.005 at 1.8  $\mu\text{m}$  [31]. Glasses were annealed at the DSC *onset*-T<sub>g</sub> [32].

#### 2.1.2 Core-glass rod for caning and *intermediate* fiber-drawing

$9.51 \times 10^{24}$  ions  $\text{m}^{-3}$  (500 ppmw)  $\text{Pr}^{3+}$ -doped Ge-As-In-Se host core-glass (Fig.1) was prepared. First, Ge-As-Se was prepared as in section 2.1.1 and the glass transferred to a silica-glass still with 500 ppmw Al-wire (5N, Alfa Aesar: O-getter). The still was sealed under vacuum ( $10^{-3}$  Pa) and distillation executed within a two-zone furnace (Instron). The distilled Ge-As-Se was re-melted 800°C/7h, quenched, annealed and transferred to a new silica-glass ampoule (air-baked then vacuum-baked, each at 1000°C/6h) and re-melted with both indium (6n5 Alfa Aesar) and 500 ppmw  $\text{Pr}^{3+}$  (3n Alfa Aesar) for 6h/850°C to form the doped core-glass rod. Again, glasses were annealed at the DSC *onset*-T<sub>g</sub> [32].

## 2.2 Extrusion, caning, fiber-drawings and SIF fabrication

The extrusion, caning, fiber-drawings and SIF fabrication are depicted in processes (a), (b), (i) and (ii), respectively, in Fig. 1 :

(a) A melt-derived Ge-As-In-Se-S cladding-glass boule (from section 2.1.1) had an outside-diameter (OD) =  $28.7 \pm 0.1$  mm and length (L) =  $17 \pm 0.1$  mm. This boule was extruded [33] to give a cladding-tube (non-RE-doped Ge-As-In-Se-S) of OD =  $10.5 \pm 0.2$  mm and ID (inner diameter) =  $1.95 \pm 0.05$  mm (see Fig.1(a)).

(b) A melt-derived  $9.51 \times 10^{24}$  ions  $\text{m}^{-3}$  (500 ppmw)  $\text{Pr}^{3+}$ -doped Ge-As-In-Se core-glass rod (from section 2.1.2) was directly caned to give unclad cane of OD =  $1.5 \pm 0.1$  mm (see Fig.1(b)).

(i) The melt-derived  $9.51 \times 10^{24}$  ions  $\text{m}^{-3}$  (500 ppmw)  $\text{Pr}^{3+}$ -Ge-As-In-Se core-glass rod (from section 2.1.2) was directly fiberised to give unclad *intermediate*-fiber, OD =  $230 \mu\text{m} \pm 20 \mu\text{m}$  (see Fig.1(i)).

(ii) The  $9.51 \times 10^{24}$  ions  $\text{m}^{-3}$  (500 ppmw)  $\text{Pr}^{3+}$ -Ge-As-In-Se unclad cane (see (b) above) was inserted in the Ge-As-In-Se-S cladding-tube (see (a) above) and then fiber-drawn as ‘rod(b)-in-tube(a)’ (under  $\text{N}_2$  gas (BOC)) to make small-core SIF with  $9.51 \times 10^{24}$  ions  $\text{m}^{-3}$  (500 ppmw)  $\text{Pr}^{3+}$  doped Ge-As-In-Se core and undoped Ge-As-In-Se-S cladding, of core-OD as- designed/fiber OD =  $20 \mu\text{m} / 130 \mu\text{m}$  and =  $40 \mu\text{m} / 270 \mu\text{m}$ , respectively, (see Fig.1(ii)).

From now on, the term: ‘ $\text{Pr}^{3+}$ /In small-core SIF’ will be used to denote the two types of small-core SIFs each with  $9.51 \times 10^{24}$  ions  $\text{m}^{-3}$  (500 ppmw)  $\text{Pr}^{3+}$  doped Ge-As-In-Se core and undoped Ge-As-In-Se-S cladding, of core-OD as-designed/fiber OD =  $20 \mu\text{m} / 130 \mu\text{m}$  or =  $40 \mu\text{m} / 270 \mu\text{m}$ , respectively. Additionally, the term ‘unstructured  $\text{Pr}^{3+}$ /In *intermediate* fiber’ will be used to denote the unstructured fiber composed of  $9.51 \times 10^{24}$  ions  $\text{m}^{-3}$  (500 ppmw)  $\text{Pr}^{3+}$ -Ge-As-In-Se and of OD =  $230 \mu\text{m} \pm 20 \mu\text{m}$  (see Fig.1(i)).

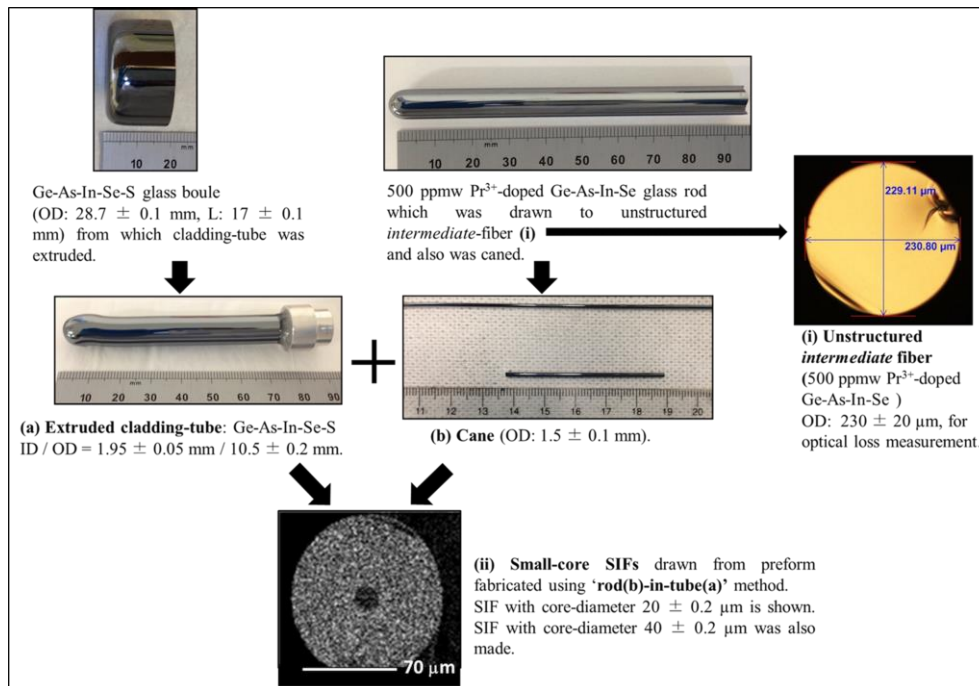


Fig. 1. Processing to produce ‘ $\text{Pr}^{3+}$ /In small-core SIFs’ (core/cladding = 500 ppmw  $\text{Pr}^{3+}$ -Ge-As-In-Se / Ge-As-In-Se-S, with as-designed  $20 \mu\text{m}$  and  $40 \mu\text{m}$  core diameters). OD is outside diameter.

## 2.3 Characterization of bulk glasses and fiber

### 2.3.1 Glass stability

Powder-XRD was done on samples to test for amorphicity in a Siemens Krystalloflex 810 X-ray diffractometer, with  $\text{CuK}\alpha$  radiation, in the range  $10 - 70^\circ 2\theta$ , in steps of  $0.05^\circ 2\theta$  per 40 s with each XRD pattern collected in  $\sim 13$  h.

### 2.3.1 Fiber optical loss

Optical loss of the ‘unstructured  $\text{Pr}^{3+}/\text{In}$  intermediate fiber’ (see Fig. 1(i)) was measured in the wavelength range  $1 - 9 \mu\text{m}$ , using the cut-back method (detailed in [16]) with an IFS 66/S, Bruker Ft-MIR spectrometer and InGaAs, InSb and MCT detectors; the optical path was not purged. Fig. 2 shows the selected ‘best’ fiber cleaves used in the fiber-loss calculation.

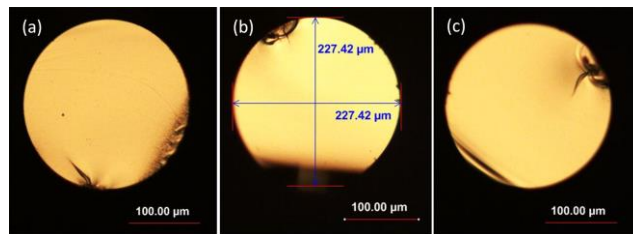


Fig. 2. (a) - (c) Optical micrographs of “best fiber-cleaves” used in fiber-loss calculation [16] of the ‘unstructured  $\text{Pr}^{3+}$ -doped intermediate fiber’ (see section 2.2, Fig. 1(i))

### 2.3.3 ‘ $\text{Pr}^{3+}/\text{In}$ small-core SIF’ cross sectional geometry

The ‘ $\text{Pr}^{3+}/\text{In}$  small-core SIFs’ (Fig. 1(ii)) were analysed as follows. SIFs were cleaved, and cross sections carbon-coated then imaged and analysed by means of (E(environmental)) SEM-BSE and (E)SEM-EDX (FEG XL30 ESEM) with an Oxford Instruments INCA x-sight Si(Li) detector with ATW2 window.

Near-field, NIR imaging of  $55\text{-}60$  mm long samples was *via* a tunable laser:  $1.465 - 1.575 \mu\text{m}$  (Agilent; 8164B).  $1.465 \mu\text{m}$  was selected as being off-centre from the  $1.45 \mu\text{m}$   $\text{Pr}^{3+}$  absorption [23], with absorption falling sufficiently by  $1.465 \mu\text{m}$  for detection of the guided light. The  $1.465 \mu\text{m}$  light was launched into the ‘ $\text{Pr}^{3+}/\text{In}$  small-core SIFs’ using a tapered silica-fiber, mounted on a XYZ translation stage, with focused spot size  $2.5 \mu\text{m}$  at the chalcogenide SIF launch-end. The output light was collected using a microscope  $\times 10$  objective and collimated onto a Siemens XQ1112 vidicon tube camera.

### 2.3.4 Emission intensity and lifetime

#### A. Emission intensity

The  $\text{Pr}^{3+}$ -doped fiber was held on a XYZ optical-stage and pumped at  $2.013 \mu\text{m}$  wavelength with a water-cooled, thulium fiber-laser, pumped at  $793 \text{ nm}$  (LISA Laser Products OHG, Germany). The emission from each  $\text{Pr}^{3+}$ -doped fiber was collected through a Spex MiniMate monochromator which had a diffraction-grating blazed at  $6 \mu\text{m}$ . In order to increase the signal-to-noise ratio, the fluorescence signal was modulated using a chopper ( $95 \text{ Hz}$ ) and a lock-in amplifier was used (Brookdeal Electronics LTD 9503SC). The detection system was based on a room temperature MCT (mercury-cadmium-telluride) detector (Vigo

System PVI-6), preamplifier for the detector (Judson PA-6) and data acquisition card (NI USB-6008 National Instruments).

### B. Emission lifetime

Fluorescent decay lifetimes at  $4.7\ \mu\text{m}$  were measured at RT using on/off modulation at 6.4 Hz of the pump laser and collected using different setups (Fig. 3), for instance the ‘Pr<sup>3+</sup>/In small-core SIFs’, 145 mm long were end-pumped at  $2.013\ \mu\text{m}$  / 12 mW with collection at the same fiber-end as pumped (see Fig. 3(c)).

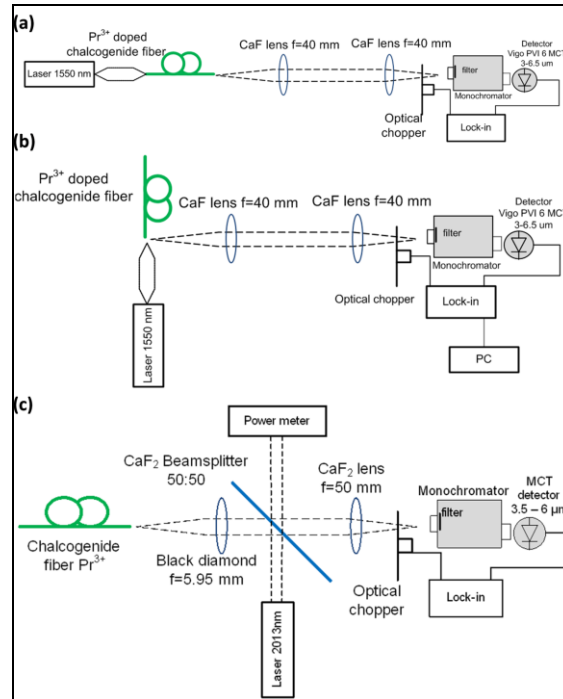


Fig. 3. Optical setups for measuring fiber photoluminescent (PL) intensity and lifetime (LT): collecting PL/LT from: (a) fiber-end opposite that pumped; (b) fiber-side and (c) from the same fiber-end as that pumped. Laser pump was  $1.55\ \mu\text{m}$  in (a) and (b), and  $2.013\ \mu\text{m}$  in (c).

For this, the pump light was passed through a collimator (Thorlabs), beam splitter (Thorlabs;  $\varnothing 2''$  CaF<sub>2</sub> plate beamsplitter, coating: 2 - 8  $\mu\text{m}$ , thickness = 3 mm) and chalcogenide-glass black-diamond lens (Thorlabs 390028-E; effective focal length ‘EFL’ = 5.95 mm); the diameter of the  $2.013\ \mu\text{m}$  wavelength excitation-beam after collimation was 2.2 mm and became  $8.2\ \mu\text{m}$  (~ circular) after focusing onto the end of the ‘Pr<sup>3+</sup>/In small-core SIFs’ (core-diameters nominally 20 and 40  $\mu\text{m}$ ). In addition, end-collection of the ‘Pr<sup>3+</sup>/In small-core SIFs’ was carried out at the opposite fiber end to that pumped when pumping at  $1.55\ \mu\text{m}$ . PL intensity and lifetimes were also measured in the ‘Pr<sup>3+</sup>/In intermediate fiber’ (see Fig.1(i)) using end-collection at the opposite fiber end to that pumped (see Fig. 3(a)) and side-collection (see Fig. 3(b)) and pumping at  $1.55\ \mu\text{m}$ . Measurements were in ambient air and corrected for system response using a Global<sup>®</sup> blackbody at 1430 K. Input pump-power was measured prior to launch into the Pr<sup>3+</sup>-doped fibers, and on exit from the fiber, using a power-meter (Thorlabs; Thermal Sensors S302C).

### 3. Results and discussion

#### 3.1 X-ray diffraction (XRD) crystallization study and optical loss spectra

XRD patterns (Fig.4) show amorphicity of: the Pr<sup>3+</sup>/ In core-glass rod (see section 2.1.2, Fig.1), the undoped cladding-tube (section 2.2(a), Fig.1(ii)) and the Pr<sup>3+</sup>/In unclad-cane (section 2.2(b), Fig.1(b)), although the XRD detection threshold is 2-5 vol.% of crystals [34]).

Fig.5 shows optical loss spectra spanning 1 - 9 μm of the ‘unstructured Pr<sup>3+</sup>/In *intermediate* fiber’ (see section 2(i), Fig.1(i)). Fig. 5 also presents, for comparison purposes, the optical loss of a Pr<sup>3+</sup>/In large-core *intermediate* SIF, reported previously in [25] (core / cladding = 498 ppmw Pr<sup>3+</sup>-Ge-As-In-Se / Ge-As-In-Se of fiber OD = 270 ± 30 μm and core-diameter = 80-85% of fiber OD, formed by extrusion then fiber-drawing).

Fig.6 presents the electronic energy level diagram of an isolated Pr<sup>3+</sup>-ion [35]. Comparing Figs.5 and 6, electronic absorption of Pr<sup>3+</sup> is evident at 1.5 μm and 2 μm, assigned to transitions from the ground-state <sup>3</sup>H<sub>4</sub> to two sets of degenerate excited-states (<sup>3</sup>F<sub>4</sub>, <sup>3</sup>F<sub>3</sub>) and (<sup>3</sup>F<sub>2</sub>, <sup>3</sup>H<sub>6</sub>), respectively. Absorption of Pr<sup>3+</sup> centered at 4.5 μm is assigned as ground-state <sup>3</sup>H<sub>4</sub>→<sup>3</sup>H<sub>5</sub>; this masks underlying broad extrinsic absorption due to the glass impurity: [H-Se]- [6]. From Fig.5, further unwanted impurity extrinsic absorption bands are: -[O-H], [H-Se]-, [H-O-H], ≡[Ge-O]-/≡[As-O]- and ≡[Si-O]- at 2.9, 3.5, 6.3, 7-7.9 and >9 μm, respectively. Loss spectra (Fig.5) indicate that the extra distillation steps followed here, compared to our earlier work [25], have reduced impurity-related absorption, especially: -[O-H], [H-O-H] and ≡[Ge-O]-/≡[As-O]- peaks at 2.9, 6.3 and 7-7.9 μm, respectively. The minimum baseline optical loss of the ‘unstructured Pr<sup>3+</sup>/In *intermediate* fiber’ (section 2.2 and Fig.1(i)) was 4.9 ± 0.3 dB/m at 6.84 μm, across 1-9 μm. These losses are higher than the best loss results to date for Pr<sup>3+</sup> doped unstructured doped glass fiber, published recently in [28]. The results of [28] give confidence that background losses can be reduced to below the 1 dB m<sup>-1</sup> limit that modeling suggests is desirable if efficient MIR fiber lasers are to be realized [36].

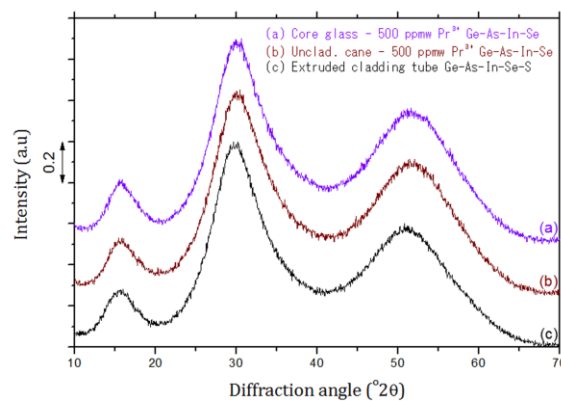


Fig. 4. XRD patterns of: (a) Pr<sup>3+</sup>/ In core-glass rod (see section 2.1.2 and Fig.1); (b) Pr<sup>3+</sup>/In unclad-cane (section 2.2(b) and Fig.1(b)) and (c) undoped cladding-tube (section 2.2(a) and Fig.1(ii)).

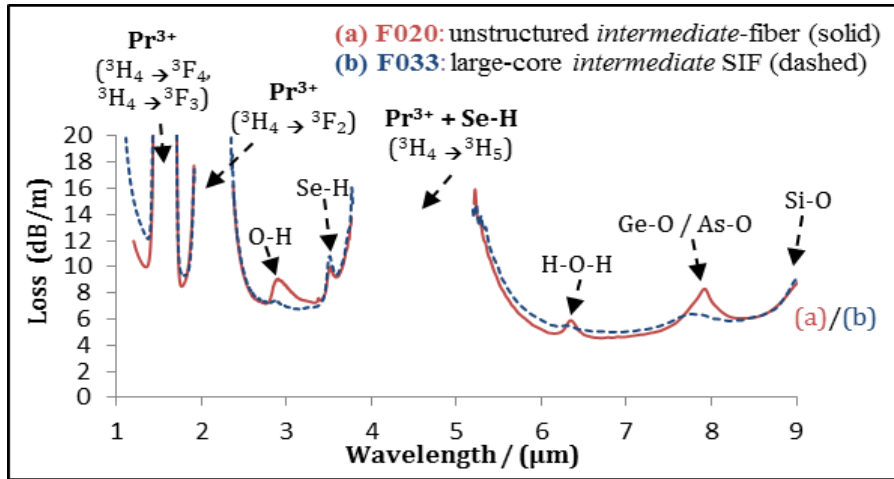


Fig. 5. Optical loss spectra across 1 - 9  $\mu\text{m}$  of the ‘unstructured  $\text{Pr}^{3+}/\text{In}$  intermediate fiber’ (section 2.2 and Fig.1(i)), which underwent glass-distillation, compared to a  $\text{Pr}^{3+}/\text{In}$  large-core *intermediate* SIF, from previous work [25] ( core /cladding = 498 ppmw  $\text{Pr}^{3+}\text{-Ge-As-In-Se}$  /  $\text{Ge-As-In-Se}$  (fiber OD =  $270 \pm 30 \mu\text{m}$ ; core diameter = 80-85% fiber OD).

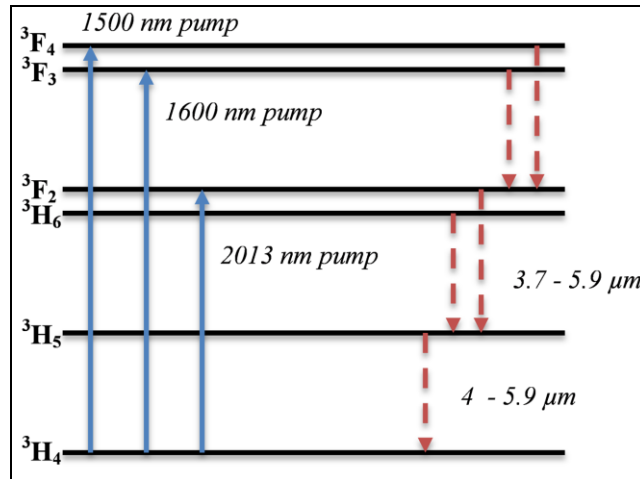


Fig. 6. Energy-level diagram of isolated  $\text{Pr}^{3+}$ -ion with absorptive transitions identified (from [31]) of interest here, together with potential longer-wavelength radiative transitions.

### 3.2 Near-field imaging and elemental mapping of the small-core SIF

Fig. 7 (a)-(e) show near-field imaging at 1.465  $\mu\text{m}$  of the ‘ $\text{Pr}^{3+}/\text{In}$  small-core SIFs’ (see section 2.2, Fig.1(ii)). Figs. 7(a), (c) demonstrate cladding modes were guided and Fig. 7(e) confirms multi-modal guiding at 1.465  $\mu\text{m}$ . Fig. 8 shows ESEM-EDX elemental mapping of In and S in the ‘ $\text{Pr}^{3+}/\text{In}$  small-core SIFs’ (see section 2.2, Fig.1(ii)). From Fig.8, the small cores were concentric, centralized and circular, and fabricated core-diameters were observed to be: (a)  $19.8 \pm 0.5 \mu\text{m}$  (*planned: 20  $\mu\text{m}$* ) and (b)  $39.4 \pm 0.6 \mu\text{m}$  (*planned: 40  $\mu\text{m}$* ).



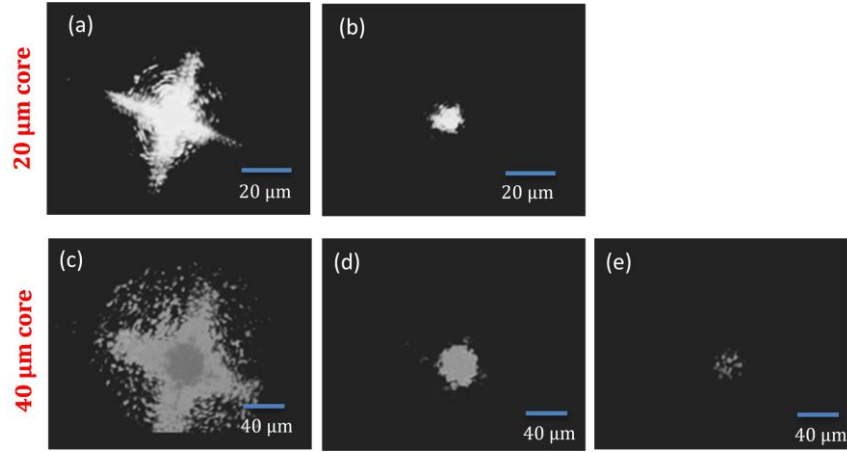


Fig. 7. Near-field images of guiding core at 1.465  $\mu\text{m}$  and different launch-angles into multimoded ‘Pr<sup>3+</sup>/In small-core SIFs’ (see section 2.2, Fig.1(ii)). (a), (b) 20  $\mu\text{m}$ -diameter small-core SIF (55 mm long). (c), (d), (e) 40  $\mu\text{m}$ -diameter small-core SIF (60 mm long)

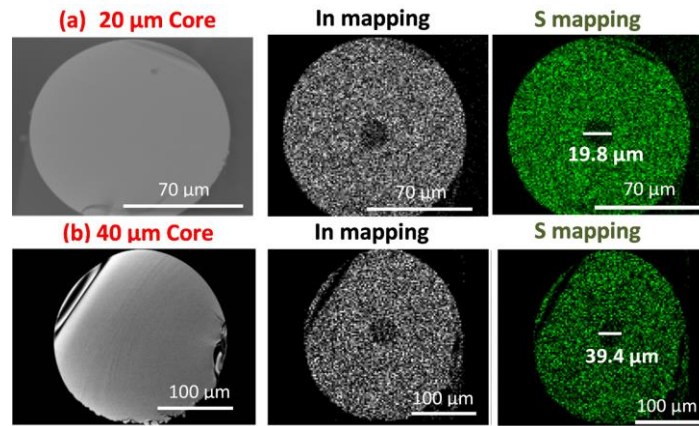


Fig. 8.(E)SEM-EDX elemental mapping of In, S in ‘Pr<sup>3+</sup>/In small-core SIFs’ (see Fig.1(ii)) showing: (a)  $19.8 \pm 0.5 \mu\text{m}$  core-diameter (designed 20  $\mu\text{m}$ ) and (b)  $39.4 \pm 0.6 \mu\text{m}$  core-diameter (designed 40  $\mu\text{m}$ ).

### 3.3 Emission in ‘Pr<sup>3+</sup>/In small-core SIFs’

Pr<sup>3+</sup> absorption bands at 1.55  $\mu\text{m}$ , 2.013  $\mu\text{m}$  were selected for pumping (see Figs.5, 6). In Fig. 9, plot (a) is prior work [25]: emission in Pr<sup>3+</sup>/In bulk glass (*melt-derived 531 ppmw Pr<sup>3+</sup>-Ge-As-In-Se*) pumped at 1.55  $\mu\text{m}$ /50.2 mW (orthogonal pump/collection across sharp corner, as in [25]), for comparison. Plot (b) is the normalized emission spectrum of the 20  $\mu\text{m}$  core-diameter ‘Pr<sup>3+</sup>/In small-core SIF’ (section 2.2, Fig.1(ii)) across 3.5-6  $\mu\text{m}$ , when end-pumped at 2.013  $\mu\text{m}$  / 25 mW, with emission collected from the same fiber-end as that pumped (as Fig. 3(c)). Under the same pump/collection conditions, the normalized emission from the 40  $\mu\text{m}$  core-diameter ‘Pr<sup>3+</sup>/In small-core SIF’ (section 2.2, Fig.1(ii)) overlaid this 3.5-6  $\mu\text{m}$  emission (*not shown*). In the absence of evidence to the contrary, the 3.5-6  $\mu\text{m}$  fluorescent emission is assigned to combined emissions: (<sup>3</sup>F<sub>2</sub>, <sup>3</sup>H<sub>6</sub> → <sup>3</sup>H<sub>5</sub>) and (<sup>3</sup>H<sub>5</sub> → <sup>3</sup>H<sub>4</sub>) (Fig.6) [37]. These transitions are radiative in selenide glass and non-radiative in fluoride glass [38] due to the high phonon energy of fluoride glasses. Fig.9 plots (c) and (d) are prior results for comparison [25], in a Pr<sup>3+</sup>/In large-core *intermediate* SIF (*core/cladding* = 498 ppmw Pr<sup>3+</sup>-Ge-As-In-Se / Ge-As-In-Se, fiber OD = 270 ± 30  $\mu\text{m}$ , core-diameter = 80-85% of fiber OD). Plot (c) is side-collection (*cf.* Fig. 3(b)) with end-pumping at 1.55  $\mu\text{m}$ /50.8 mW; plot (d)

is fiber-end collection (*cf.* Fig. 3(a)) at the opposite end from that pumped (at 1.55  $\mu\text{m}$ /24.5 mW). Plot (d) is broadened relative to plot (c) at the low energy side of the spectral emission band, indicating that photon trapping [39] occurred when emission was collected at the fibre-end opposite to that pumped (as in Fig.3(a)) and less so when collected from the side of the fiber.

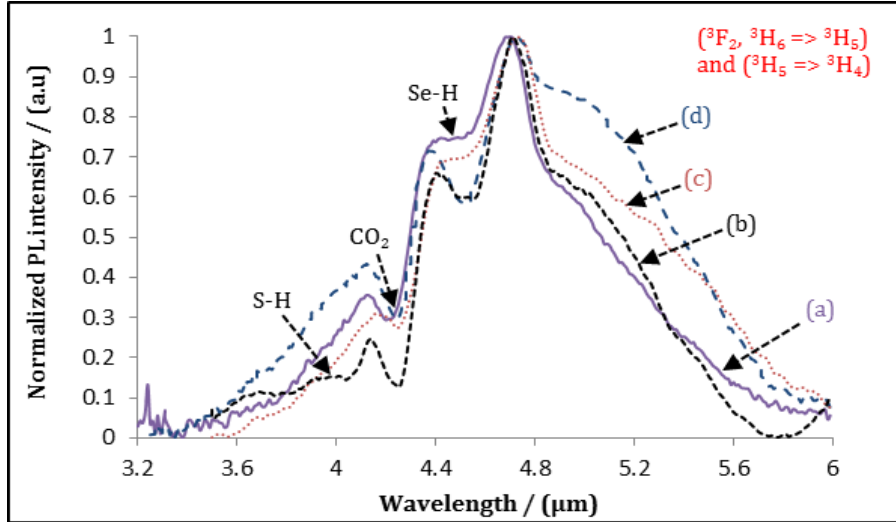


Fig. 9. Normalized fluorescence-emission at 3.5-6  $\mu\text{m}$ : (a) prior work: *bulk* glass (531 ppmw  $\text{Pr}^{3+}$ -Ge-As-In-Se) with orthogonal pump (1.55  $\mu\text{m}$ /50.2 mW), fluorescent-emission collection across a sharp corner [25]; (b) 20  $\mu\text{m}$  core-diameter ' $\text{Pr}^{3+}$ /In small-core SIF' (section 2.2, Fig.1(ii)) end-pumped at 2.013  $\mu\text{m}$  /25 mW with fluorescent-emission collection from same fiber-end as pumped (*cf.* Fig. 3(c)); (c) from prior work [25]:  $\text{Pr}^{3+}$ /In large-core *intermediate* SIF (core/cladding = 498 ppmw  $\text{Pr}^{3+}$ -Ge-As-In-Se/Ge-As-In-Se (fiber OD:  $270 \pm 30 \mu\text{m}$ , 80-85% core-diameter)) end-pumped at 1.55  $\mu\text{m}$ /50.84 mW with fluorescent-emission collection at side of fiber (Fig. 3(b) and (d) fiber same as (c) [25], but end-pumped, at 1.55  $\mu\text{m}$  / 24.5 mW with fluorescent-emission collection at *opposite end of fiber* to that pumped (Fig. 3(a)).

Fig. 10(a) shows the variation in emission intensity at 4.7  $\mu\text{m}$ , of the 20  $\mu\text{m}$  and 40  $\mu\text{m}$  core-diameter ' $\text{Pr}^{3+}$ /In small-core SIFs' (section 2.2, Fig.1(ii)), as a function of the input pump-power for end-pumping at 2.013  $\mu\text{m}$  / 0-180 mW; emission was collected from the same fiber-end as that pumped (*cf.* Fig. 3(c)). Fig. 10(b) shows the variation in pump-power exiting the opposite end of the fiber to that pumped, for pumping 2.013  $\mu\text{m}$  / 20-180 mW of 20  $\mu\text{m}$  and 40  $\mu\text{m}$  core-diameter ' $\text{Pr}^{3+}$ /In small-core SIFs' (section 2.2, Fig.1(ii)) as a function of the input pump-power. From Fig. 10, both fibers show an approximate linear increase in emission intensity at 4.7  $\mu\text{m}$  exiting the same fiber-end as that pumped, and pump-power exiting the fiber-end opposite to that pumped, with increasing input pump-power from 10-80 mW. However, from 80-170 mW pump-power, the emission intensity and pump-power exiting the pumped fiber became nonlinear: the emission intensity response became super-linear and the pump-power exiting response became sub-linear.

We suggest this behavior could be related to bleaching of the electronic ground state ( $^3\text{H}_4$ , Fig. 6), thereby limiting further pump absorption [40] and maybe evidence of approaching saturation of an excited state. Fluorescent peak intensity of RE-ions would normally start to saturate as a population-inversion is approached. We have found no reports of similar nonlinear behavior in the literature for MIR fluorescent-emission at greater than 3.9  $\mu\text{m}$  [7], nor pump-power exiting, in RE-ion doped glass fibers. It is not claimed that a population-inversion has been achieved here and gain would have been suppressed by the presence of

resonant extrinsic absorption, and possibly extrinsic non-radiative decay, [6, 41] due to [H-Se]- impurity, oxide and other impurities in the host. Our prior work, and that of others, has shown that ultra-low optical loss of undoped chalcogenide glass fibers may be achieved [16 and refs. therein]. Here, the Pr<sup>3+</sup>-dopant and In solubilizer have brought into the glass-host extra unwanted impurities and raised the optical loss in the fiber [42].

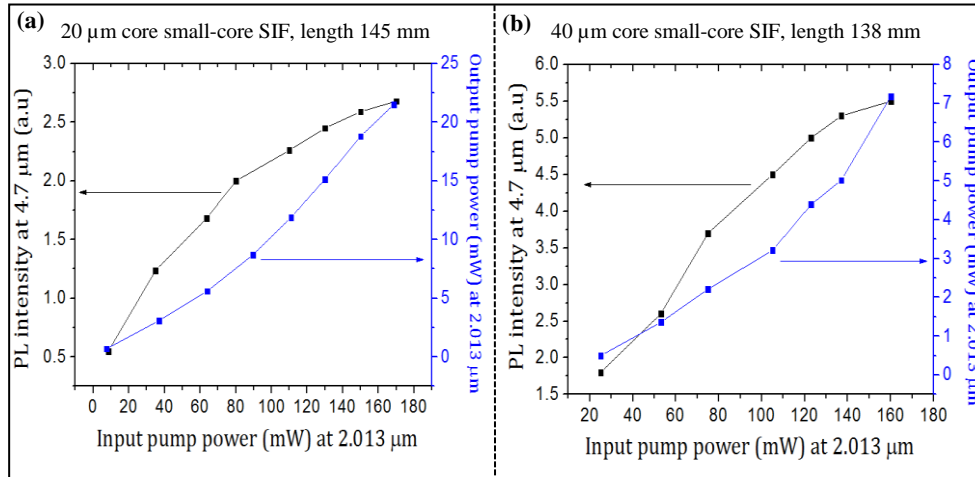


Fig. 10. For end-pumping at 2.013 μm: (a) 20 μm core-diameter ‘Pr<sup>3+</sup>/In small-core SIF’ (section 2.2, Fig.1(ii)) and (b) 40 μm core-diameter ‘Pr<sup>3+</sup>/In small-core SIF’ (section 2.2, Fig.1(ii)). Plotted for each of (a) and (b) are the output pump-power collected from the opposite fiber end to that pumped (*blue solid squares*), and fluorescent-intensity at 4.7 μm collected from the same fiber-end as that pumped (*black solid squares*) (see Fig.3(c)), plotted as a function of the incident input pump-power.

### 3.4 Fluorescent-lifetimes

Table 1 collates the fluorescent-lifetimes of the ‘Pr<sup>3+</sup>/In *intermediate* fiber and small-core SIFs’ (section 2.2, Fig.1(ii) and (ii), respectively) at 4.7 μm wavelength, when pumping at 1.55 μm or 2.013 μm compared to ‘Pr<sup>3+</sup>/In bulk-glass’ from prior work [25].

From Table 1, the Pr<sup>3+</sup> fluorescent-lifetime at 4.7 μm in the presence of an In solubilizer tends to decrease from *bulk* glass (10.1 ± 0.3 ms) [25], to *intermediate* fiber (8.0-8.2 0.5 ms) to small-core SIFs (7.1 ± 0.5 ms). This decrease in lifetime for Pr<sup>3+</sup>/In samples is in contrast to Pr<sup>3+</sup>/Ga samples, for which the Pr<sup>3+</sup> fluorescent-lifetime at 4.7 μm, when pumped at 1.55 μm, of 7.8 ± 0.3 ms for *bulk*-glass was maintained in small-core SIFs [23]).

Below, we argue that the decrease in Pr<sup>3+</sup> fluorescent-lifetime in the presence of an In solubilizer goes hand-in-hand with an increase in the cumulative-processing required to prepare each type of sample.

Processing is required to achieve shaping of preforms, cane, tubes and fiber. Each processing-step here comprised a heat-treatment where the glass was heated to some temperature between T<sub>g</sub> and the liquidus, in the supercooled-liquid regime, for shaping to take place, followed by cooling to ambient. In cases other than caning or fiberizing, an isothermal-hold took place at T<sub>g</sub> during the cooling- step in order to anneal the glass and abolish permanent stress. The supercooled-liquid is thermodynamically metastable and tends to lower free energy by undergoing nucleation then crystallisation. Accompanying this phase change (or pre-phase-change, in the case of glass nucleation) there may be a tendency for RE-ions in the melt to cluster. Any

nucleation, crystallisation or clustering ‘frozen-in’ on cooling to  $T_g$ , could ultimately lead to concentration-quenching of the excited states of the RE-ions.

Shaping of the supercooled liquid takes time, and concurrent nucleation/ crystallisation/clustering depends on the degree of undercooling during shaping, which influences not only the nucleation rate but also the supercooled liquid viscosity which is the macroscopic manifestation of the kinetic limitations to phase change and to clustering influencing the time-taken for atomic-diffusion. In a nutshell, during processing-steps, it is the cumulative time of visits to the supercooled liquid (heating, cooling or isothermal hold) which eventually leads to a non-homogeneous glass at ambient temperature.

From Table 1, labeled (I), the  $\text{Pr}^{3+}/\text{In}$  bulk glass ( $\text{Pr}^{3+}\text{-Ge-As-In-Se}$  [25]) exhibited the longest lifetime of all samples ( $10.1 \pm 0.3\text{ms}$ ); the glass was melted above the liquidus (where there is no thermodynamic tendency for devitrification or clustering) then cooled through the supercooled-liquid regime to  $T_g$ , annealed and cooled to ambient. In short, this sample received processing comprising {1x cooling} through the supercooled-liquid range.

Table 1.  $\text{Pr}^{3+}$  fluorescent-lifetime, in the presence of In, at  $4.7 \mu\text{m}$  in bulk glass, *intermediate* fiber and small-core SIF, arranged vertically downwards in order of decreasing PL lifetime (column 2), with details of samples and measurement setups. (Key: SIF: step index fiber, N/A not applicable, OD: outer diameter,  $\lambda$ : wavelength.)

<b><math>\text{Pr}^{3+}</math>-doped with In solubilizer: sample type.</b>	<b>PL lifetime /ms</b>	<b>Pumped at <math>\lambda / \mu\text{m}</math>. (Power/mW)</b>	<b>Pump launch.</b>	<b>PL collection set-up.</b>	<b>• [<math>\text{Pr}^{3+}</math>]/ ppmw. • Glass host. • (Cladding glass)</b>	<b>Sample geometry.</b>
<b>(I) Bulk glass [25]</b>	<b><math>10.1 \pm 0.3</math></b>	1.55 (57.8)	1 mm from sharp corner.	1 mm from sharp corner, orthogonal to pump.	• <b>531</b> • Ge-As-In-Se • (N/A)	Polished sharp corner.
<b>(II) Unstructured <i>intermediate</i> fiber</b>	<b><math>8.0 \pm 0.5</math></b>	1.55 (50.8)	"	Fiber side (orthogonal to pump) Fig.3(b).	• <b>500</b> • Ge-As-In-Se • (N/A)	Fiber OD $230 \pm 20 \mu\text{m}$ . (section 2.2, Fig.1(i)).
	<b><math>8.2 \pm 0.3</math></b>	"	"	Fiber end (opposite to pump) Fig.3(a).		
<b>(III) Small-core SIF</b>	<b><math>7.2 \pm 0.5</math></b>	1.55 (2.2)	"	Fiber end (opposite to pump) Fig.3(a).	• <b>500</b> • Ge-As-In- Se • (Ge-As-In- Se-S)	Core OD $20 \mu\text{m}$ . Fiber OD $130 \pm 10 \mu\text{m}$ (section 2.2, Fig.1(ii)).
	<b><math>7.0 \pm 0.1</math></b>	2.103 (12.0)	"	Fiber end (same end as pump) Fig.3(c).		

The ‘unstructured  $\text{Pr}^{3+}/\text{In}$  *intermediate* fiber’ (section 2.2, Fig.1(i)), labeled (II) in Table 1, was drawn directly from a melted, quenched, annealed  $\text{Pr}^{3+}$ -doped glass rod. In short, this sample received processing

comprising: {(1x (cooling) then 1x (heating/cooling))} in the supercooled-liquid regime. The fiber fluorescent-lifetimes measured by side- and end-collection agreed.

The 20  $\mu\text{m}$  core-diameter ‘Pr<sup>3+</sup>-doped small-core SIF’ (section 2.2, Fig.1(ii), sample IV in Table 1) had the shortest fluorescent-lifetimes ( $7.1 \pm 0.5$  ms) of all samples and underwent the most intricate fabrication processing, consisting of: (i) melting, quenching and annealing of the Pr<sup>3+</sup>-doped core and undoped cladding glasses followed by: (ii) cane-drawing Pr<sup>3+</sup>-doped core-glass; (ii) placing Pr<sup>3+</sup>-doped cane in the undoped cladding tube (made by extruding the undoped-cladding glass) and fiber-drawing ‘rod-in-tube’ (see Fig.1(ii)). The core/cladding interface was formed during the fiberizing. In short, this sample received processing comprising: {1x (cooling) then 2x (heating/cooling)} in the supercooled-liquid regime.

We speculate that the processing to form the SIF may have caused some Pr<sup>3+</sup> clustering and/or glass-nucleation/crystal growth. We showed earlier [25] that Pr<sup>3+</sup>/In selenide bulk-glasses have more tendency than their Pr<sup>3+</sup>/Ga counterparts to devitrify at  $\geq 1000$  ppmw Pr<sup>3+</sup>. An increase in light scattering loss accompanying loss of glass homogeneity could not be studied in the small-core SIF by direct optical-loss measurement due to the too large étendue of the MIR blackbody source available for measuring broadband loss.

It is important to note that the measured Pr<sup>3+</sup> lifetimes in Table 1 may be of a subset of Pr<sup>3+</sup> ions that are not quenched by, [H-Se]- or -[O]- impurities [6, 41].

Potential radiative mechanisms include that during excitation, pump upconversion could have occurred, due to potential resonance of the 2.013  $\mu\text{m}$  pump-laser and the energy absorptive process  $^3\text{H}_5 \rightarrow ^3\text{F}_3$  of Pr<sup>3+</sup>, after initial pumping to [ $^3\text{F}_2$ ,  $^3\text{H}_6$ ] followed by radiative emission:  $^3\text{H}_6 \rightarrow ^3\text{H}_5$ . The  $^3\text{H}_5$  level could also have been depopulated through this process. Further, 2.013 or 1.55  $\mu\text{m}$  pumping could incur cross-relaxation between neighbouring Pr<sup>3+</sup> ions:  $^3\text{H}_6 \rightarrow ^3\text{H}_5$  could have caused excitation upconversion of  $^3\text{H}_6 \rightarrow ^3\text{F}_3$ . In which case, emission could also have been contributing from  $^3\text{F}_3 \rightarrow ^3\text{H}_4$ , known to lase at 5.2  $\mu\text{m}$  in Pr<sup>3+</sup>-doped chloride crystals [43]. From Table 1, the measured lifetimes of the ‘Pr<sup>3+</sup>/In small-core SIFs’ were the same for 1.55  $\mu\text{m}$  or 2.013  $\mu\text{m}$  pumping. This suggests that the levels  $^3\text{H}_6$  and  $^3\text{H}_5$  are depopulated in a similar manner, whether pumping at 1.55  $\mu\text{m}$  or 2.013  $\mu\text{m}$ . Thus, 2.013  $\mu\text{m}$  pump-upconversion was probably minimized, at least for this doping concentration.

#### 4. Conclusions

20 and 40  $\mu\text{m}$  core diameter SIFs, with core / cladding =  $9.51 \times 10^{24}$  ions  $\text{m}^{-3}$  (500 ppmw) Pr<sup>3+</sup> doped Ge-As-In-Se / Ge-As-In-Se-S, have been fabricated with NA  $\sim 0.4$  (estimated from [23]) and multimode operation at 4-6  $\mu\text{m}$ , encompassing potential lasing wavelengths. Each core was circular and centralized, and of diameter close to the design. NIR light-guiding in the SIF core was demonstrated. The lowest possible SIF optical loss, across 3-9  $\mu\text{m}$  was 5 dB  $\text{m}^{-1}$  at 6.84  $\mu\text{m}$ . The SIF fluorescent-emission spanned 3.5-6  $\mu\text{m}$ , when end-pumped at 2.013  $\mu\text{m}$ . The fluorescent-decay lifetime at 4.7  $\mu\text{m}$  wavelength of the 20  $\mu\text{m}$  core fiber ( $7.1 \pm 0.5$  ms) was lower than that found in *intermediate* fiber ( $8.1 \pm 0.5$  ms) or in bulk glass ( $10.1 \pm 0.3$  ms, [25]). This trend appears to be related to the number of processing-steps to fabricate SIF causing potential decrease in glass homogeneity and some greater ordering around, or of, the doped-in Pr<sup>3+</sup>. (This result is in contrast to

analogous Pr<sup>3+</sup>/Ga selenide-glass small-core SIF which was shown, in earlier work [23], to maintain the same PL lifetime ( $7.8 \pm 0.3$  ms) in SIF as in the parent *bulk*-glass.) When end-pumping the 20 and 40  $\mu\text{m}$  core-diameter Pr<sup>3+</sup>/In SIFs at 2.013  $\mu\text{m}$  / 80-170 mW, the Pr<sup>3+</sup> fluorescent emission intensity, and residual pump-power, exiting the SIFS became super-linear, and sub-linear, respectively, as the pump-power was increased from 80-170 mW. These trends for the first time suggest MIR excited-state saturation of a Pr<sup>3+</sup>-dopant in a chalcogenide fiber.

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