

# Influence of gate bias on the avalanche ruggedness of SiC power MOSFETs

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**Abstract**—This paper investigates the effect of negative gate bias voltage ( $V_{GS}$ ) on the avalanche breakdown robustness of commercial state-of-the-art silicon carbide (SiC) power MOSFETs. The device's ability to withstand energy dissipation during avalanche regime is a connoting figure of merit for all applications requiring load dumping and/or benefiting from snubber-less converter design. The superior material properties of SiC material means that SiC MOSFETs even at 1200V exhibit significant intrinsic avalanche robustness.

**Keywords**—avalanche ruggedness; silicon carbide; unclamped inductive switching; power MOSFET; robustness

## I. INTRODUCTION

A device's ability to withstand energy dissipation in avalanche breakdown regime ( $E_{AV}$ ) is a connoting figure of merit for all applications requiring load dumping and/or benefiting from snubber-less converter design even in the presence of non-zero parasitic inductance values [1]. The avalanche ruggedness is typically achieved in silicon (Si) MOSFETs and IGBTs as a result of progressive device design and engineering optimisation by delaying parasitic BJT activation [2]. On the other hand, the specific material properties of silicon carbide (SiC) (e.g., the higher energy band gap) means that SiC MOSFETs even at 1200 V exhibit significant built-in robustness, which deserves dedicated study. However, the higher band gap in SiC makes it highly unlikely for the activation of parasitic npn BJT. Various studies have already shown that SiC power MOSFETs can dissipate  $E_{AV}$  up to approximately 1J depending on test conditions [3-5]. This paper aims to investigate the dependence of avalanche ruggedness of SiC power MOSFETs on gate bias voltage ( $V_{GS}$ ). Experimental results at different  $V_{GS}$  along with simulation results are included in this paper.

## II. EXPERIMENTAL RESULTS

Avalanche ruggedness is characterized experimentally by way of unclamped inductive switching (UIS) tests [4, 6]. The schematic of a UIS test setup is shown in Fig. 1. An auxiliary IGBT with a 3kV breakdown voltage was used in parallel to ramp up the inductor current ( $I_L$ ) to the required value, avoiding the self-heating prior to avalanche phase. A

schematic presentation of the inductor current ( $I_L$ ) and drain voltage ( $V_{DS}$ ) waveforms for a safe UIS transient are also presented in Fig. 2. The experimental results presented here are on 1.2kV 36A rated SiC power MOSFET (C2M0080120D) in a TO-247 package from CREE [7]. Different tests were carried out on device under test (DUT), referring to circuit in Fig. 1, in the first case  $V_{GS} = 0V$  was used, and in second case  $V_{GS} = -5V$  while keeping all the other test conditions unaltered. The aim of this analysis was to observe the effect of different  $V_{GS}$  on the avalanche capability of SiC power MOSFETs. All the tests presented here were carried out at input voltage ( $V_{DD}$ ) = 400V, load inductor ( $L_{LOAD}$ ) = 500 $\mu$ H and case temperature ( $T_{CASE}$ ) = 25°C while changing  $V_{GS}$ .

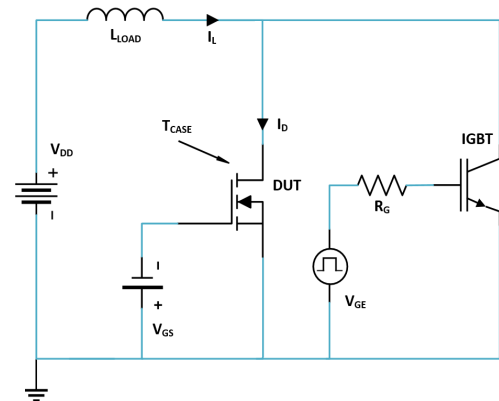


Fig. 1. Schematic of the UIS test setup

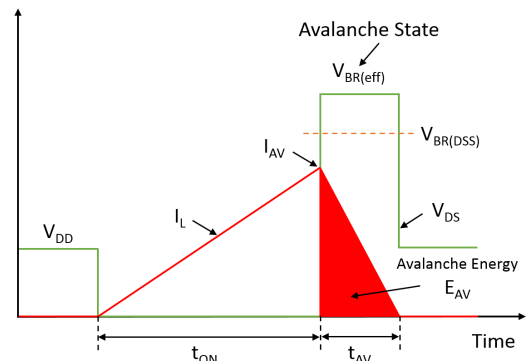


Fig. 2. Representative  $I_L$  and  $V_{DS}$  waveforms for safe UIS test

Starting with  $V_{GS} = 0V$ , Fig. 3 shows typical  $I_L$  and  $V_{DS}$  waveforms ( $I_{AV} \sim 43A$ ) for a safe UIS event. A safe UIS test is characterized by the return of  $I_L$  and  $V_{DS}$  to zero and input voltage ( $V_{DD}$ ) respectively. The on time ( $t_{ON}$ ) of the IGBT was gradually increased to increase the peak avalanche current ( $I_{AV} \sim 47A$ ) until failure was observed. Moreover, under the same test conditions, the gate bias was changed to  $V_{GS} = -5V$ . In this case, even though  $I_{AV}$  is the same, the device safely completes the avalanche phase and returns to blocking afterwards. A low output impedance power supply was used to supply  $V_{GS}$ . For comparison, Fig. 4 plots  $V_{DS}$  and  $I_L$  for both  $V_{GS} = 0V$  and  $-5V$  for the same  $I_{AV}$ . Here, it can be seen that the DUT failed for  $I_{AV} \sim 47A$  for  $V_{GS} = 0V$ . On the other hand, the DUT tested for  $V_{GS} = -5V$  sustained the avalanche breakdown regime and did not fail. Fig. 5 presents the  $I_L$  and  $V_{DS}$  waveforms ( $I_{AV} \sim 50A$ ) at failure for  $V_{GS} = -5V$ . For DUT tested at  $V_{GS} = -5V$ , higher  $I_{AV}$  i.e. higher  $E_{AV}$  was needed before failure was observed. The failure is characterized by the collapse in  $V_{DS}$  and hence, the DUT loses its ability to block voltage. At failure,  $I_L$  starts to increase again (as dictated by  $V_{DD}$  and  $L_{LOAD}$ ) since the device experiences an internal short between all the DUT terminals followed by a catastrophic failure of the device. Several papers have demonstrated characterization of SiC power MOSFETs during avalanche breakdown also showing failure [3-6, 8, 9].

The avalanche energy ( $E_{AV}$ ) can be defined using equation 1.  $E_{AV}$  was calculated for all test conditions and included in the table presented here (bold text for failure  $E_{AV}$ ). Finally, all the test conditions for the results presented here are also summarized in Table I.

$$E_{AV} = \frac{1}{2} I_0 V_{BR(eff)} t_{AV} \quad (1)$$

Table I. Summary of Experimental result

$V_{DD}$ (V)	$L_{LOAD}$ ( $\mu H$ )	$V_{GS}$ (V)	$I_{AV}$ (A)	$E_{AV}$ (J)
400	500	0	43	0.75
			47	<b>0.84</b>
		-5	47	0.88
			50	<b>0.96</b>

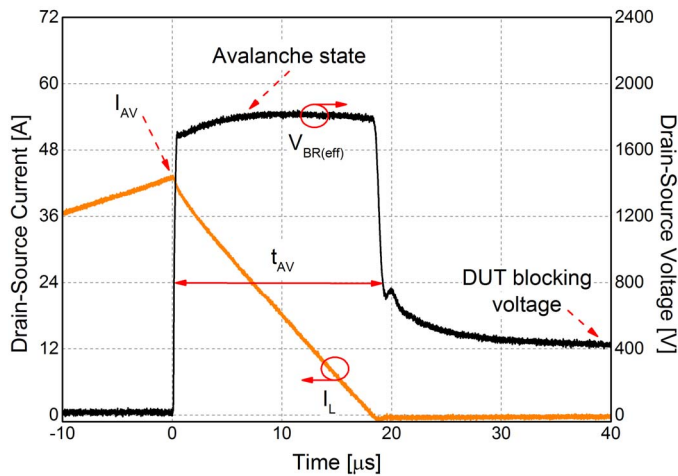


Fig. 3. Experimental  $I_L$  and  $V_{DS}$  waveforms for safe UIS test;  $V_{GS} = 0V$

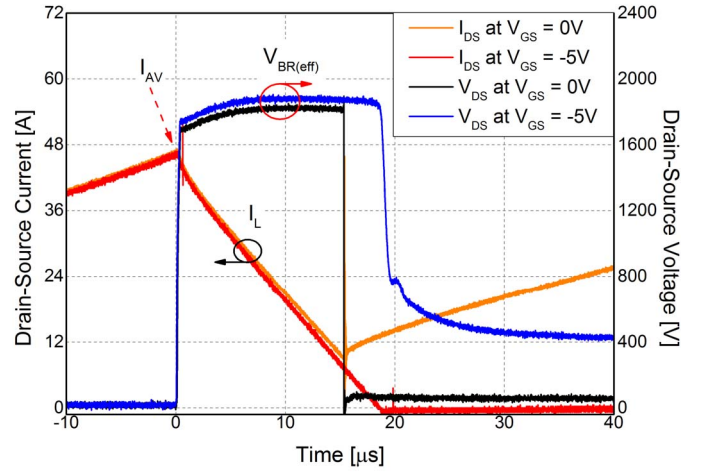


Fig. 4. Experimental  $I_L$  and  $V_{DS}$  waveforms for safe and failure UIS test for  $V_{GS} = 0V$  and  $V_{GS} = -5V$  respectively; Comparison

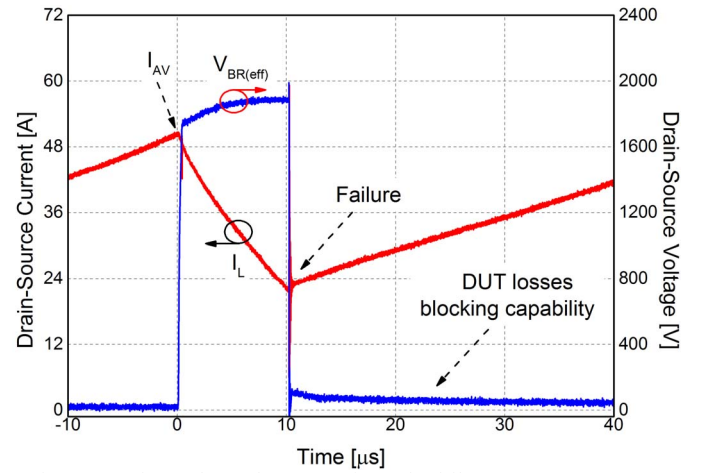


Fig. 5. Experimental  $I_L$  and  $V_{DS}$  waveforms for failure UIS test;  $V_{GS} = -5V$

### III. TCAD SIMULATION RESULTS

TCAD electro-thermal simulation results reporting the failure mechanism of SiC power MOSFETs during avalanche breakdown operation complement the study of the device behaviour under different gate source bias voltages ( $V_{GS}$ ). This work advances a previous study on the UIS failure mechanism of SiC power MOSFETs [8], providing new important insight into the dependence of avalanche robustness on gate bias voltage. Electro-thermal mixed-mode 2D simulations were used to reconstruct the observed failure and interpret the experimental results. Fig. 6 shows the simulated planar cell structure with important dimension parameters and doping concentrations.

In these simulations, two physically separated electrodes were used for the P-body and N+ implant regions of the cell structure but connected to the same electrical node. The simulations were carried out at  $V_{DD} = 400V$  and  $L_{LOAD} = 500\mu H$  while changing  $V_{GS}$ . Fig. 7 present the failure UIS simulation for  $V_{GS} = 0V$ . The current components flowing out of both P-body and N+ source contacts are also shown in Fig. 8, together with the overall device current. As can be seen in Fig.8, the electron current starts to flow into the N+ source

region as the time in avalanche ( $t_{AV}$ ) lapsed until failure occurred.

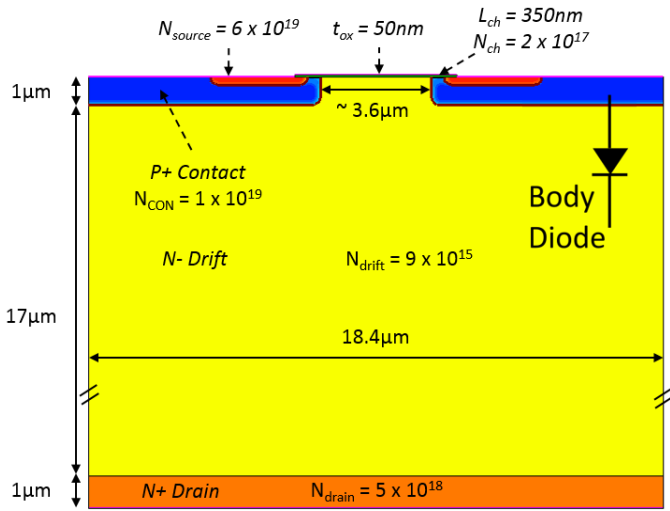


Fig. 6. Simulated planar structure (Not to scale)

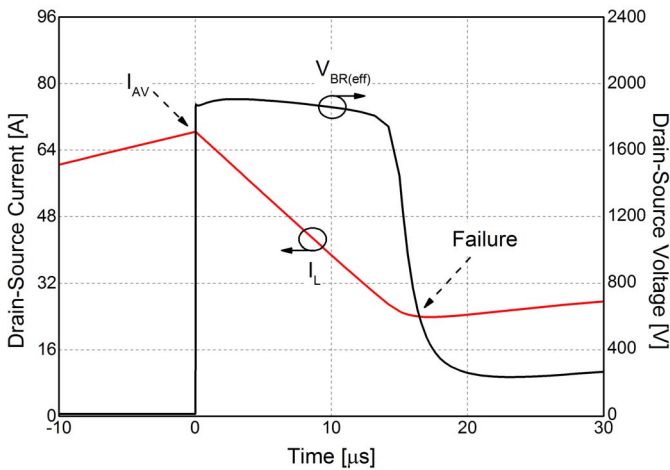


Fig. 7. Simulated  $I_L$  and  $V_{DS}$  waveforms for failure UIS event;  $V_{GS} = 0V$

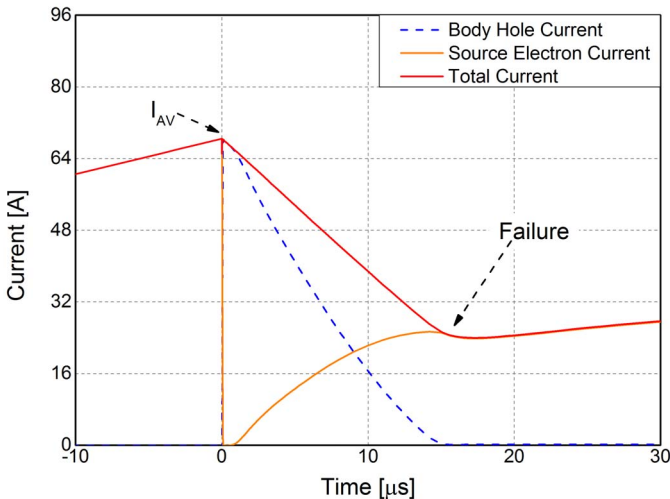


Fig. 8. Source electron and body hole current components for  $I_L$  in Fig. 7

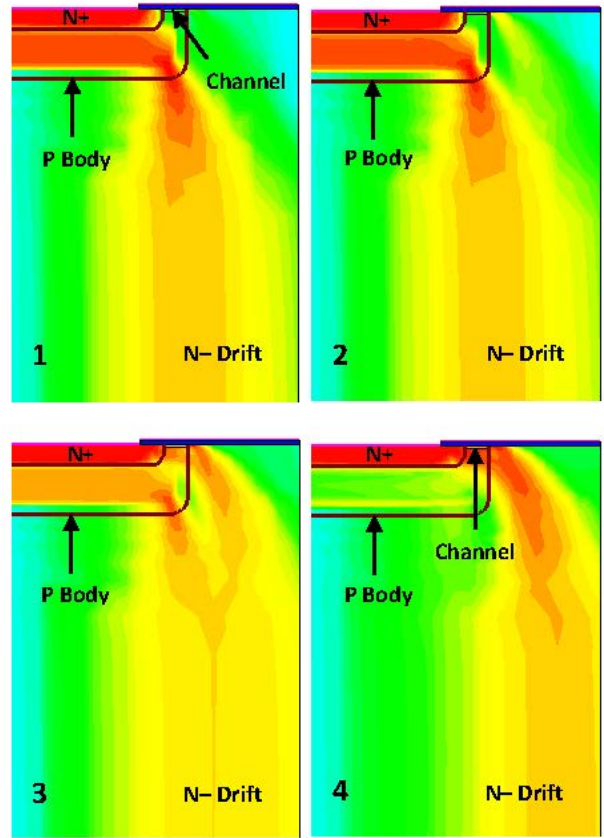


Fig. 9. Total Current Density; Zoomed in near p-body/n-drift region (1-4 with increasing  $t_{AV}$ )

During avalanche breakdown, the highest electric field occurs at the curvature of the body/drift pn junction and hence, the maximum impact ionization also happens at the same place. While the device is in avalanche, the current density usually flows over the body diode of the MOSFET going through the pn junction curvature as also presented in Fig. 9 (1). The junction temperature ( $T_J$ ) due to really high power density in such a small device (die size: 3.10 x 3.36 mm) during such short UIS transients could easily rise significantly to really high values well above 1000K. Fig. 9 presents current density at different time instances starting from time instance immediately after entering avalanche breakdown (1) until after failure (4). As a result of such high temperature increase (cell temperature simulated well above 1000K at failure), the current also starts to flow in and below the channel area aided by the reduction of threshold voltage ( $V_{th}$ ) due to consistent temperature increase near p-body as shown in Figure 9 (2 and 3). After failure, only electron current flows in and below the channel as shown in Figure 9 (4). Mixed-mode simulations results for  $V_{GS} = -5V$  are also included in Fig. 10. Here, slightly higher  $I_{AV}$  was required before failure was obtained. Simulations have shown that by using a negative  $V_{GS}$  to keep the device turned off helps to better close the channel. Therefore, it takes longer (i.e. higher temperature) before the onset of source electron current flowing in and underneath the channel. Hence, channel activation is slightly delayed with negative  $V_{GS}$  allowing the device to sustain a slightly higher  $I_{AV}$  and  $E_{AV}$  before encountering failure.

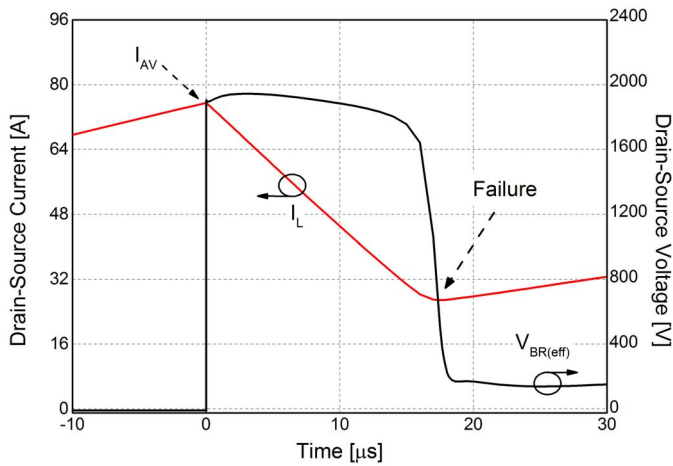


Fig. 10. Simulated  $I_L$  and  $V_{DS}$  waveforms for failure UIS event;  $V_{GS} = -5V$

Effect of temperature on  $V_{th}$  has also been investigated experimentally in [10] which also partially supports the presented interpretation of  $V_{th}$  lowering which results in current flow in and underneath channel at failure as the DUT  $T_J$  increases.

#### IV. THERMAL MAPPING

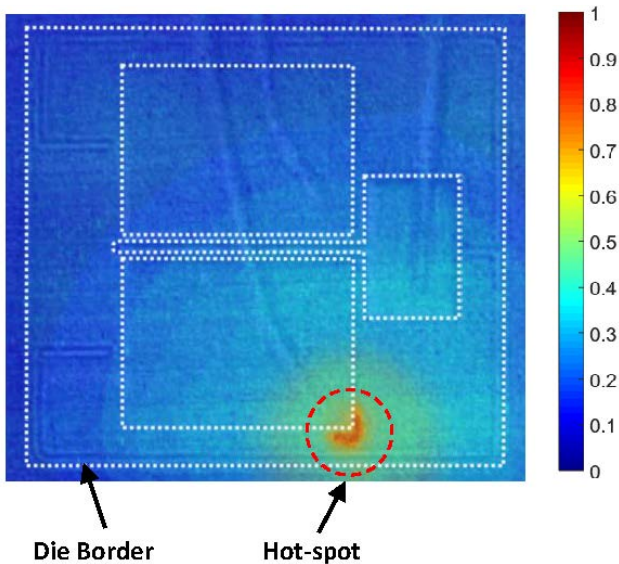


Fig. 11. Hot-spot formation during UIS

The thermal mapping on bare-die devices was also performed using infrared (IR) thermography (explained in [11]). The IR map showed the formation of a localised hot-spot due to current crowding mechanism in a small group of cells close to the source metal contact as can be seen in Fig. 11. The hot-spot formation occurs due to a positive feedback mechanism involving bipolar current flow. The formation of a hot-spot usually indicates that bipolar currents flows at failure. It is either possible that the actual device failure occurs at an earlier  $t_{AV}$  time instant when the hole current is still flowing which also explains the hot-spot formation since the simulated cell structure is not the true representation of a real device and its associated packaging features. Moreover, the simulations

involve just a single cell so the interactions from the surrounding cells is not taken into account thus making it difficult to obtain hot-spot formation with a single cell.

#### V. CONCLUSION

This paper investigated the influence of gate source bias ( $V_{GS}$ ) on the avalanche breakdown robustness of SiC power MOSFETs. Experiments showed that negative turn-off  $V_{GS} = -5V$  slightly enhances the avalanche withstand capability of the devices. Si power devices can be turned-off with  $V_{GS}$  bias up to  $-15V$  whereas SiC is limited in negative bias. The experimental results were also complemented using TCAD simulations which showed progressive shift of current from the body diode towards the channel during avalanche phase. After failure, the current flow was entirely in and underneath the channel. Lastly, thermal mapping on bare die devices are also presented. Unlike in Si, the activation of npn BJT within SiC devices seems highly unlikely due to higher band gap and weak gain of the parasitic bipolar transistor.

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