

What we think we know about the aerodynamic performance of windows

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Abstract

Natural ventilation is a low energy strategy used in many building types. Design approaches are mature but are dependent on variables with high uncertainty, such as the aerodynamic behaviour of purpose provided openings (PPOs), which need improved characterisation.

An analytical framework is used to define different types of flow through openings based on the balance of environmental forces that drive flow, and the different flow structures they create. This allows a comprehensive literature review to be made, where different studies and descriptive equations can be compared on a like-for-like basis, and from which clear gaps in knowledge, technical standards, and design data are identified. Phenomena whose understanding could be improved by analysis of existing data are identified and explored.

A Statistical Effective Area Model (SEAM) is developed from academic data to estimate the performance of butt hinged openings during the design

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stage, that accounts for the impact of aspect ratio and opening angle. Its predictions are compared against available empirical data and are found to have a standard error of 1.2%, which is substantially lower than the 15 – 25% prediction errors of *free* area models commonly used in practice.

An analytical model is made based on entrainment theory to explain the increase in flow rate that occurs through two aligned openings. This model defines characteristic design parameters and predicts a detrimental impact on the ventilation of the wider space.

Finally, an analytical model is created to explain the reduction in discharge coefficient that occurs when a large temperature difference exists across an opening. This model defines novel dimensionless parameters that characterise the flow, and predicts empirical data well, suggesting that it should be integrated into design equations.

Keywords: Purpose provided opening, ventilation, model, prediction, geometry, *free* area

Highlights

- Framework developed and used to determine aerodynamic performance
- Evaluation of existing literature
- Quantification of ambiguities implicit in existing modelling techniques
- Development of new analytical models

1. Introduction

Natural ventilation is seen by many to be a crucial part of a low energy building strategy, but is sometimes perceived to be risky and unreliable.

To encourage the widespread use of natural ventilation, reliability issues need to be addressed. A major source of prediction error stems from a poor understanding of the aerodynamic performance of window openings (PPOs) and whole building systems [1, 2, 3, 4]. This paper will show that techniques for modelling real openings are flawed, and create systematic errors in performance predictions that, when brought to light in the under-performance of the finished building [5], can damage the reputation of natural ventilation design. Improving these techniques will help design systems that are robust enough to perform under a wide range of environmental conditions, and help to restore confidence in the ability of natural ventilation to deliver efficient, functioning buildings.

Section 1 introduces the fundamental concepts behind envelope flow models and their approach to calculating air flow rates through openings. Section 2 develops a framework that systematically breaks down the assumptions made to simplify these calculations, and uses it to structure a comprehensive literature review into the behaviour of openings when these assumptions are violated. This review is used to identify key gaps in research, experimental data and technical standards. Some areas where understanding can be improved by analytical or statistical modelling are identified, and explored in Sections 3–5. Key conclusions from the work are summarised in Section 6.

24 *1.1. Designing a ventilation strategy*

25 When designing a natural ventilation strategy for a building there are
26 two key stages [6]. The first is to define the desired flow pattern of air within
27 the building, which often varies seasonally to best satisfy occupant comfort.
28 The second task is to design the envelope. This involves positioning and
29 sizing openings so that the required airflow pattern and volume flow rates
30 are achieved under the design conditions [7]. In practice, this is typically
31 achieved using envelope flow models. The main attraction of these models is
32 their simplicity: in many cases basic hand calculations suffice [3, 8, 7].

33 *1.2. Principles of envelope flow models*

34 The fundamental concepts of envelope flow models are very simple, and
35 can broadly be divided into two separable components [3, 8, 7]. The first is
36 the calculation of the pressure differentials exerted on the building envelope
37 that drive airflow through a building. These are caused by an interaction
38 between the building geometry with wind and thermal buoyancy forces [9].
39 Empirical and experimental techniques for estimating these pressure distri-
40 butions are given in [3, 8, 10], with extensive data sets for generic building
41 types given by [11]. The second component to predicting airflow through
42 envelope flow models is the characterisation of the aerodynamic performance
43 of openings in the building envelope that admit airflow. It is this second
44 component that forms the basis of this paper.

45 Openings in a building envelope can be divided into two types: adven-
46 titious openings and purpose provided openings (PPOs) [2]. Adventitious
47 openings are unintentional, and comprise cracks and gaps in the building en-
48 velope. PPOs are created intentionally as part of the ventilation scheme, and

49 often take the form of operable windows or vents. All the theory described
50 henceforth concerns PPOs, and assumes that adventitious openings account
51 for a negligible fraction of overall ventilation rates.

52 *1.2.1. Key assumptions of the treatment of PPOs in envelope flow models*

53 In a conventional envelope flow model, several assumptions are made
54 about the aerodynamic performance of PPOs to ensure their independence of
55 calculations of driving pressure, and to simplify modelling of their resistance
56 to airflow [7, 12, 13, 3, 14]. Some typical assumptions are:

- 57 • Openings in the envelope are small, so that they do not significantly
58 alter the pressure distributions on the façade.
- 59 • PPOs can be treated as an equivalent sharp-edged orifice
- 60 • Internal and external density profiles are uniform across the height of
61 the opening, and do not vary with flow rate
- 62 • Internal air motion is negligible
- 63 • Flow characteristics of openings in wind can be given by their still-air
64 characteristics
- 65 • The pressure field across the opening is approximately uniform and
66 equivalent to the pressure measured at its centre
- 67 • Ventilation is pseudo-steady - the time averaged flow characteristic is
68 unaffected by turbulence

69 To some extent these issues can be resolved by combining conventional
70 envelope flow models with zonal models, dynamic thermal models, or com-
71 putational fluid dynamics (CFD), but this comes at the cost of increased
72 computational complexity and time [15].

73 1.3. Theory of flow through openings

74 In a conventional envelope flow model, airflow through PPOs is described
75 by the orifice flow equation [3, 8, 7]. The names and definitions of the terms
76 used in this equation vary between sources, particularly those describing mea-
77 surements of area. Therefore, this paper follows the convention of Jones *et*
78 *al.* [2] to avoid ambiguity.

79 The orifice flow equation can be derived trivially by application of the
80 Bernoulli equation to a streamline passing through a constriction for the case
81 where the ambient air is quiescent on either side of the opening [7, 10, 3];
82 see Figure 1. This relates the volume flow rate to the pressure drop in the
83 constriction, and the minimum area through which the fluid passes.

$$Q = A_{min} \sqrt{\frac{2(P_1 - P_2)}{\rho}} \quad (1)$$

84 Here, Q is the volume flow rate, A_{min} is the minimum area through which
85 the fluid passes, and P_1 and P_2 are the static pressures on the streamline
86 upstream of the constriction and at the point of maximum constriction re-
87 spectively.

88 It is common in envelope flow models to treat PPOs as sharp-edged ori-
89 fices [7, 4]. When fluid flows through a sharp-edged opening, flow separation
90 occurs at the edges. This results in a characteristic flow pattern where the

91 fluid passes through a contracted area smaller than the opening, known as
 92 the *vena contracta* [4, 16]; see Figure 1. This represents the minimum area
 93 specified in Equation 1. The ratio of the area of the *vena contracta*, A_{min} , to
 94 that of the opening, A_f , is the contraction coefficient C_c [4]. The term A_f is a
 95 geometric parameter associated with the opening known as the *free* area, and
 96 is commonly defined as the minimum unobstructed area perpendicular to the
 97 flow, although this varies between sources [2, 17, 18]. The flow separation
 98 caused by the sharp edges means that the value of the contraction coefficient
 99 does not vary significantly with Reynolds number [7]. An additional factor,
 100 C_f , is included to account for frictional resistance [4]. The product of these
 101 is termed the discharge coefficient, C_d , and results in the equation

$$Q = C_d A_f \sqrt{\frac{2(P_E - P_I)}{\rho}} \quad (2)$$

102 The discharge coefficient of a two dimensional slit can be derived the-
 103 oretically, and evaluates to approximately 0.611 [19]. This is very close to
 104 experimentally derived values for a sharp edged circular orifice, which typi-
 105 cally lie between 0.6 and 0.65 [3, 10]. While the discharge coefficient would
 106 be expected to be different for different opening geometries, a discharge coef-
 107 ficient of *circa* 0.61 [8] or 0.65 [10] is commonly used to model any arbitrary
 108 PPO. Although the measurement of *free* area is trivial for a circular hole, it
 109 becomes much more complex and ambiguous for real PPO geometries; see
 110 Section 2.1.1. The product of the discharge coefficient and the *free* area is
 111 known as the *effective* area, A_{eff} , which represents the aerodynamic prop-
 112 erties of the opening in still air. Equation 2 can then be rearranged to find
 113 the *effective* area of openings required to provide a given flow rate under the

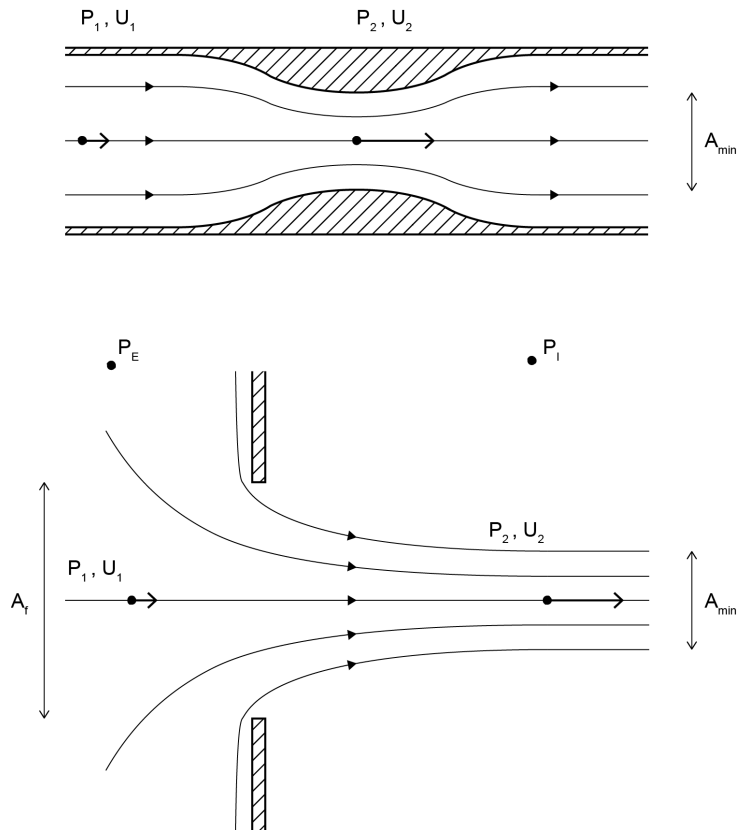


Figure 1: Comparison between the model of flow through a constriction (top), and its application to flow through a sharp-edged opening between an external (E) and internal (I) space (bottom).

114 design pressure difference. In off–design cases, the *effective* area can be used
115 to evaluate airflow rates through the building under the influence of a range
116 of weather conditions.

117 The fundamental theory is well understood, but when it is applied to real
118 buildings many of the key assumptions of the orifice flow and envelope flow
119 models are either violated (for example the still air assumption in wind driven
120 flows) or only partially fulfilled (for example the still air assumption under
121 light wind conditions) [4]. Consequently, when openings are installed in real
122 buildings their aerodynamic performance often differs from that observed
123 under laboratory conditions [6] or predicted by simple envelope flow models
124 [12]. A good deal of research has been undertaken, both analytically and
125 experimentally, to ascertain the causes of these deviations in aerodynamic
126 performance, so that they can be adequately accounted for in the design
127 process.

128 **2. Analytical framework for studying airflow through purpose pro-** 129 **vided openings**

130 One advantage of creating a structured framework for analysing flow
131 through openings, is that it enables the literature to be analysed systemati-
132 cally, gaps in the research to be identified, and the degree to which sources
133 provide useful predictive tools to be assessed.

134 To make the analysis independent of building configuration, airflow through
135 the openings is considered in isolation based on the environmental condi-
136 tions at their internal and external surfaces. Here, the problem of estimating
137 flow through window openings is broken down into assumptions affecting the

138 mechanisms that drive flow through them, and organised into a decision tree;
139 see Figure 2. The performance of an opening can then be characterised using
140 basic assumptions about the driving forces, and the impacts when each of
141 these assumptions is invalid can be systematically investigated.

142 For convenience, the framework is broken down into three sections, which
143 are shown in Figure 2. The upper section is the system definition, which
144 outlines some of the key assumptions of the model geometry, the properties of
145 the working fluid, and the flow structures present in the ambient environment.
146 Below the system definition the tree is split into two sections: one where the
147 external air is still, and the other where the external air is in motion. These
148 allow the impact of the two mechanisms that drive flow – wind and buoyancy
149 – to be evaluated both in isolation and in concert.

150 Sections 2.1–2.3 systematically describe the framework, and uses it to
151 structure a review of the literature. Section 2.4 summarises the extent of
152 knowledge identified using the framework, and identifies key gaps in the
153 research.

154 *2.1. The system definition*

155 Figure 3 shows the system definition, which details some fundamental as-
156 sumptions about the properties of the window and its environment that are
157 required before simplified modelling methods can be applied. These assump-
158 tions are applicable to both still and moving air. The first two assumptions
159 describe simplifications of model geometry; the third describes assumptions
160 about fluid properties; and the final two assumptions describe the flow struc-
161 tures on the inside and outside of the opening. Resolving the final assump-
162 tion divides the structure into two branches, describing conditions where the

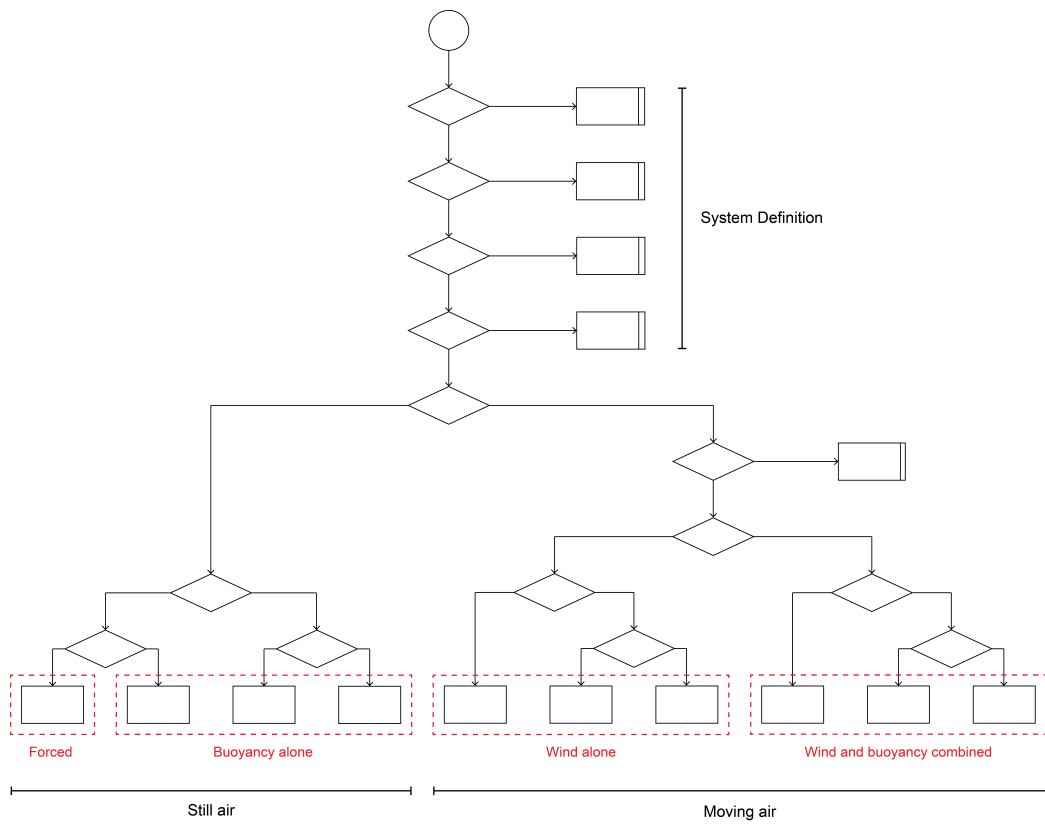


Figure 2: Overview diagram of analytical framework separating flow scenarios by driving mechanism and modelling assumptions. Branches of this framework are shown in detail in Figures 3, 7, and 8

163 external air is still and in motion, respectively.

164 *2.1.1. Two-dimensional opening assumption*

165 One of the most common simplifications of model geometry assumes that
166 any PPO can be modelled as a two-dimensional opening; see Section 1.3.
167 Much of the literature uses the two-dimensional opening assumption explic-
168 itly, both in simplified physical models [20, 21, 22, 23] and CFD analysis
169 [24, 25, 26, 27]. Many other sources study flow behaviour using three-
170 dimensional window geometries [28, 29, 30], but assume that they can be
171 represented as two-dimensional openings using unvalidated area conversions.
172 This makes it especially difficult to compare results between sources.

173 While the two-dimensional opening condition is well approximated for
174 openings where all components share a common plane with the structural
175 opening (such as sliding windows; see Figure 4), it cannot be said to be valid
176 for opening geometries that contain elements that project from the plane
177 of the structural opening. The projecting elements associated with three
178 dimensional openings can act to restrict flow, alter the shape and direction of
179 the streamlines passing through them, and change the way openings interact
180 with external flow. A few studies directly examine the impact of complex
181 opening geometry for cross ventilation [31, 32, 33], single sided ventilation
182 [34, 35], and for specialised airflow units [36, 37], but this approach is not
183 widespread.

184 When estimating airflow through an opening, it is necessary to charac-
185 terise its resistance to airflow. For design purposes, the most important
186 determinant of resistance is the *effective* area of the opening [3, 2]. While
187 it is conventionally assumed that the flow capacity of an opening is depen-

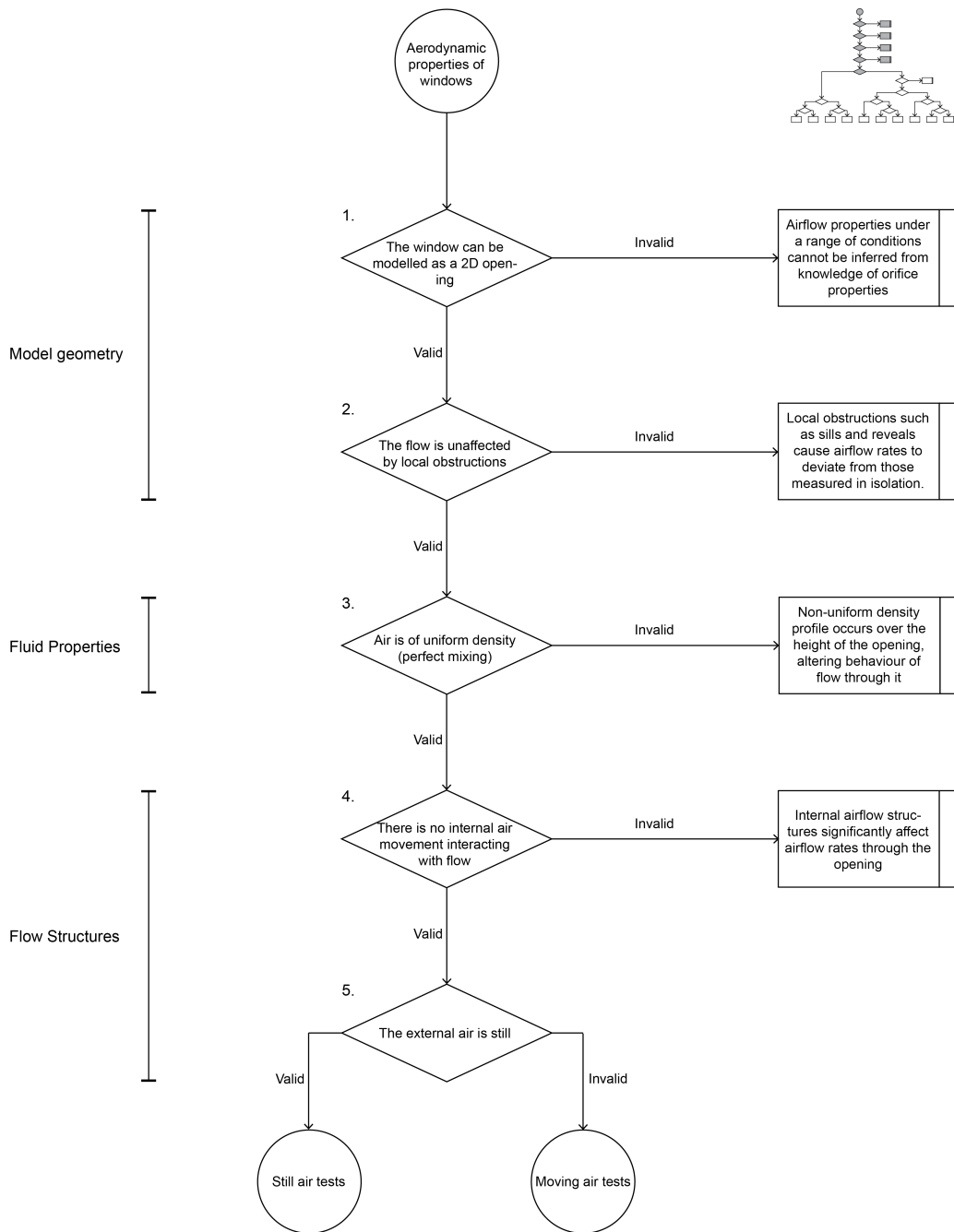


Figure 3: The system definition for the assumption tree detailing key assumptions about the nature of a ventilation opening and its environment.

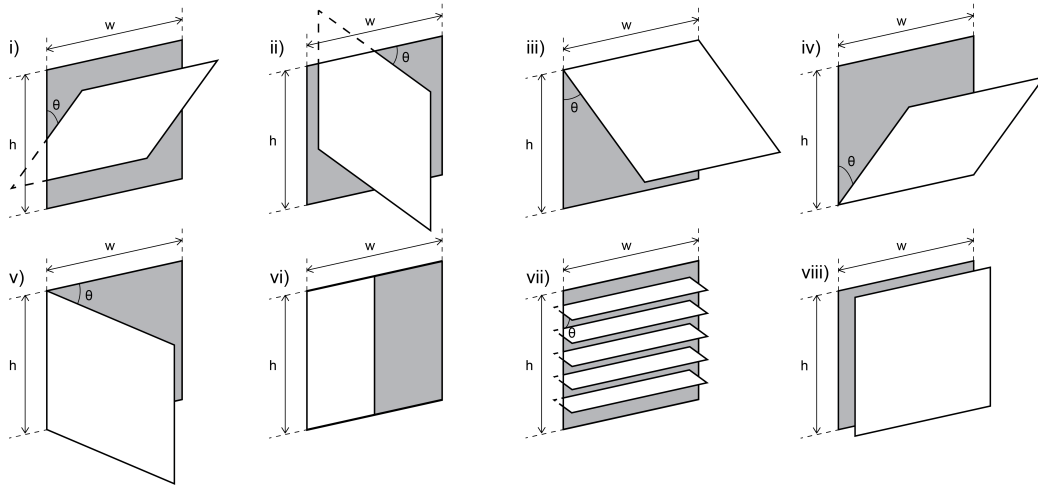


Figure 4: Common types of operable window: (i) horizontal pivot; (ii) vertical pivot; (iii) top hung; (iv) bottom hung; (v) side hung; (vi) sliding; (vii) louvre; (viii) parallel plate [39]

188 dant only on the total area [38], knowledge of the resistance distribution, and
 189 hence the area distribution, is required when the pressure across the opening
 190 is non-uniform. This effect is particularly important when the opening is
 191 very large compared to other openings in the ventilation system, or when all
 192 the openings are in similar locations in the pressure field.

193 *2.1.1.1. Operable windows.* One of the most common types of PPOs are
 194 operable windows. This paper follows the conventions given in CIBSE Guide
 195 B2 [39] for the naming of common window geometries shown in Figure 4.
 196 For the purposes of evaluating *effective* area, opening types (i-ii), and types
 197 (iii-v), can be considered identical, and are subsequently referred to as pivot
 198 and hinged openings, respectively.

199 In practice, it is common to calculate the *effective* area of an opening by
 200 assuming a constant discharge coefficient, and evaluating the *free* area based

201 on inspection of the window geometry, given by

$$A_{eff} = Cd A_f(\theta, h, w) \quad (3)$$

202 where θ is the opening angle, h is the height of the opening, w is the width
203 of the opening, and $A_f \leq hw$. Estimating the *free* area of a window is
204 often assumed to be a trivial problem. Consequently, there has not been a
205 systematic study of how this is done, or of the impacts of any errors associated
206 with its estimation on predictions of window performance. However, it is
207 clear from the literature that the definition of *free* area is ambiguous [2], and
208 that different practitioners approach it in different ways. Figure 5 illustrates
209 a range of approaches to calculating the *free* area of hinged openings, all of
210 which are based on the sum of different measured areas. Little theoretical
211 justification is given for each area model, and comparison with empirical data
212 is very rare. This ambiguity is a major source of error both in practice and
213 in academia [2].

214 Jong and Bot [31, 32] produce empirical data based on still-air pressuri-
215 sation tests for simple hinged openings, which they use to fit coefficients to
216 analytical *free* area model 'f' shown in Figure 5. This model however is
217 unnecessarily complicated by a number of analytical factors that could be
218 readily combined, and predicts *effective* areas that tend to infinity as the
219 height to width ratio becomes large.

220 An alternative approach is to define a fixed, easily measurable area for an
221 opening and to derive its discharge coefficient experimentally as a function
222 of opening angle. This defines the *effective* area and the discharge coefficient
223 as

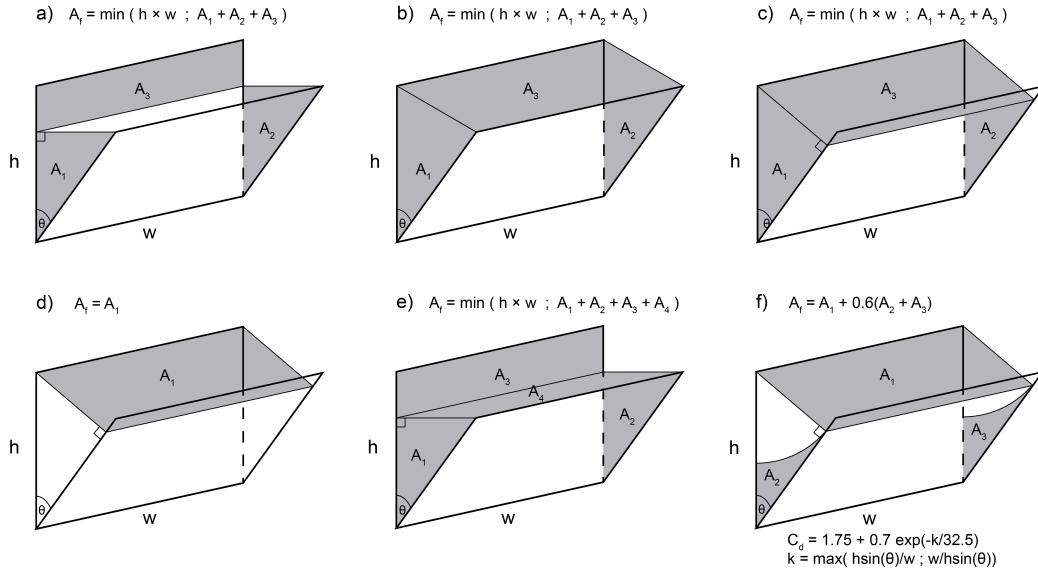


Figure 5: A range of different methods of modelling *free area* [28, 30, 34, 40, 31, 32, 18]. Model *f* is semi-empirical, and includes a co-efficient allowing for reduced efficiency of the side areas that is adjusted to fit experimental data.

$$A_{eff} = Cd(\theta, \sigma) h w \quad (4)$$

$$Cd(\theta, \sigma) = \frac{Q}{h w} \sqrt{\frac{\rho}{2 \Delta P}} \quad (5)$$

224 where σ is the aspect ratio, $h:w$. Figure 6 describes how the characteristic
 225 dimensions of a nominal window opening – its height, width, area, opening
 226 angle, and thickness – can be measured. Note that these definitions are
 227 applicable to all opening types given in Figure 4, as well as to windows that
 228 use sliding hinges where the pivot point moves in the vertical plane as θ
 229 varies.

230 The UK design guidance for the ventilation of school buildings, Building

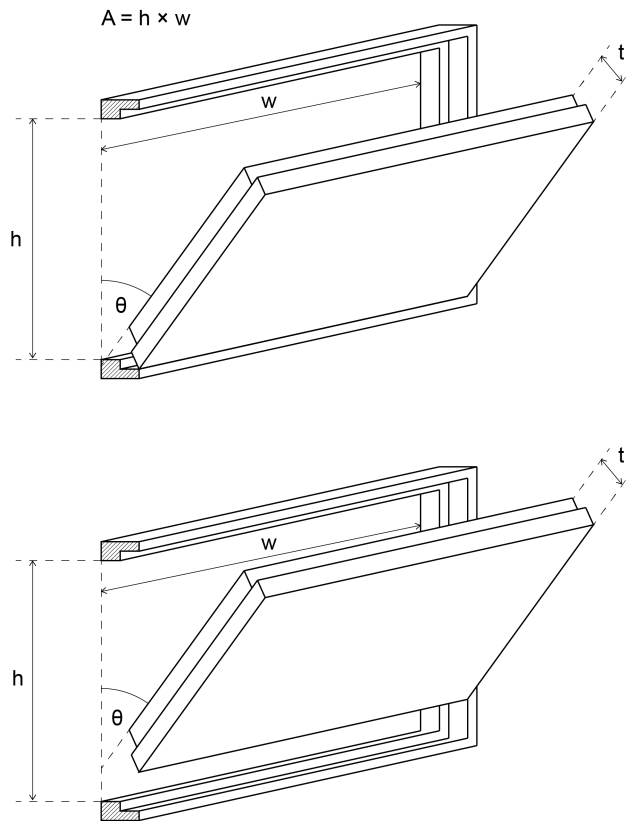


Figure 6: Simplified measurement method to define the dimensions and discharge coefficient of an opening, where A is the *free* area of the opening, h is the internal height of the fixed frame, w is the internal width of the fixed frame, t is the thickness of the opening sash, and θ is the angle between the planes of the fixed frame and opening sash known as the opening angle. A summary of its application to different opening geometries is given in Figure 4.

231 Bulletin 101 [41] (BB101), presents a simple statistical model of the variation
232 in discharge coefficient with opening angle based on this approach, fitted to
233 proprietary data¹. A new statistical model that uses academic data [31, 32,
234 33] is developed in Section 3 for use in design, and benchmarked against
235 existing models in Section 3.1.

236 *2.1.1.2. Chimneys, ducted outlets and wind catchers.* A range of ventilation
237 technologies use long ducts to transport air around a building, either to access
238 deeper spaces or provide access to favourable pressure conditions [23, 7].
239 These opening types typically cannot be modelled as a sharp-edged orifice;
240 their discharge coefficient needs to be given as a function of Reynolds number
241 [7]. The use of Reynolds-dependent discharge coefficients allows these types
242 of opening to be integrated into conventional envelope flow models.

243 In addition to their effect on discharge coefficient, these technologies af-
244 fect the driving forces available for natural ventilation. Chimneys increase
245 the stack height available for buoyancy ventilation [7, 8, 3]; solar chimneys
246 increase the air temperature within the stack, raising buoyancy pressure;
247 and wind catchers, chimney tops and roof cowls alter the wind pressure co-
248 efficients at the inlet/outlet to enhance flow [15, 42, 3, 7]. While in many
249 cases these pressures can be evaluated independently of flow rates through
250 the ducts [36, 7, 23], this is not universally true. As a result, purely empiri-
251 cal models are sometimes used to quantify the airflow performance of these
252 components under a range of conditions [43].

¹All versions available from DOI: 10.13140/RG.2.2.10748.08323

253 *2.1.1.3. Mesh screens.* To improve security and reduce the risk of ingress of
254 bugs and animals, mesh screens are often installed in openings. Flow through
255 these screens has been the subject of considerable study, and a summary of
256 experimental data and modelling methods is given by Bailey *et al.* [44]. The
257 discharge coefficient of screens are highly dependent on Reynolds's number
258 [7, 44], and this relationship can be expressed as a function of the porosity of
259 the screen and the thickness of the wires [44]. Alternatively, flow rates can
260 be characterised by a power law [45] or quadratic relationship [7] instead of
261 the conventional orifice flow equation.

262 *2.1.2. Unobstructed flow assumption*

263 To allow the properties of an opening to be determined separately from
264 the room in which it is installed, it is often assumed that airflow rates are
265 unaffected by local obstructions, such as sills and reveals. However, sills and
266 reveals can restrict the area available for air to pass through, as well as affect
267 how the window geometry interacts with external airflow.

268 In academia and in practice it is common for these local obstructions to be
269 accounted for as a reduction in *free* area [38, 18]. While the technique makes
270 analytical sense, it is subject to the same ambiguities and errors associated
271 with the geometric models discussed in Section 2.1.1.

272 The analytical technique developed by Hall [46] to improve the predic-
273 tion of single sided ventilation rates was found to reduce errors in predicted
274 performance at very small opening angles. However, the technique requires
275 empirical data to calibrate it and no justification is given for extrapolating
276 the results to higher opening angles or to cross ventilation configurations. In
277 the absence of experimental data, the resultant errors cannot be quantified.

278 In addition to the physical obstructions that occur due the installation
279 position of a PPO within a building's fabric, it is possible that external
280 obstacles - such as trees or louvres - or internal obstacles - such as people,
281 blinds, furniture, and partitions - could interact with the structure of airflow
282 through the openings. To the best of our knowledge there is no existing
283 research that quantifies the effect of these obstacles on the performance of
284 any PPO.

285 Mesh screens are typically installed within other opening types, and so
286 have the potential to interact aerodynamically. Bailey *et al.* suggest calculat-
287 ing the combined resistance to airflow caused by a mesh screen set within a
288 window frame by summing the resistance factors ($F = 1/C_d^2$) determined for
289 the two components in isolation, but provide no experimental data to sup-
290 port this. A similar approach might be applied to account for other internal
291 obstacles. Tabulated design equations for the integration of mesh screens
292 with louvres are given by Holzer and Psomas [42]. It is not clear how these
293 screens would interact with other opening geometries.

294 2.1.3. *Uniform density assumption*

295 A common simplification of envelope flow models arises from the assump-
296 tion that the air is of uniform density and perfectly mixed. This assumption
297 is known to be invalid in most cases because hot air rises from heat sources,
298 and stratifies near the ceiling [22, 47]. This may have a significant impact on
299 the pressures exerted across a window opening, resulting in substantial errors
300 in the prediction of airflow rates. In this case, the bulk of the error is in the
301 magnitude of the driving pressures and not the aerodynamic properties of
302 the opening itself. Flow through the opening only behaves differently if the

303 density profile is non-uniform across its height. A method of modelling a
304 non-uniform density profile for a room is given in CIBSE AM10 [3], but this
305 is unsubstantiated. Linden [47] develops a model based on plume physics
306 describing thermal stratification, but this cannot account for the interaction
307 between the range of heat sources and mixing mechanisms likely to be present
308 in real buildings. Given that there is no effective method of predicting the
309 density profile in a room [7], all analysis hereon assumes the density of the
310 air is uniform across the height of the opening.

311 *2.1.4. Internal air movement assumption*

312 Envelope flow models commonly assume that the internal air is static;
313 see Section 1.2.1. This assumption not only implies that the resistance to
314 airflow caused by the internal space can be neglected, but that patterns of
315 internal air movement cannot interfere with the dynamics of flow through
316 the PPOs. In reality, internal air movement can come from a number of
317 sources. Thermal plumes rising from occupants and machinery, gusting from
318 mixing fans and turbulence from movements within the space can all play a
319 role. These factors are complex to predict, and even when they are known
320 it would be hard to design an experimental procedure to account for the
321 range of possibilities. In CFD simulations, Shetabivash [25] identifies that
322 the velocity profile of an opening is insignificantly altered by its location,
323 despite the substantial variation in the internal flow pattern. This suggests
324 that a study of internal air movement is unimportant for predicting bulk
325 airflow, although it may be important in assessing local pollutant transport
326 or thermal comfort. In contrast, Hall finds that the presence of an internal
327 heater located below a bottom hung, inward opening window can reduce

328 buoyancy driven single sided ventilation rates by up to 20% [46]. This is
329 likely to be due to the fresh air supply entraining into the rising plume, which
330 leaves the space without properly mixing with the room air. Given that it is
331 common to locate emitters beneath windows to prevent cold downdraughts,
332 further research into this phenomena is warranted.

333 Internal air motion has a greater impact on bulk flow rates when the
334 inlet and outlet are in close proximity. Heiselberg and Sandberg [4] and
335 Seifert *et al.* [48] identify the formation of a stream tube between the inlet
336 and the outlet, where a flow connection causes kinetic energy to be conserved.
337 Consequently, the conventional orifice flow equation tends to underestimate
338 volume flow rates through the openings. This implies that ventilation sys-
339 tems over-perform the predictions of the orifice flow model [7] when openings
340 are closely aligned. However, airflow within a stream tube may bypass the
341 occupied portion of a room and be could be less effective at removing con-
342 taminants from there [10]. The shape of the streamlines approaching the
343 opening would also be altered, which could alter the resistance to airflow
344 provided by the opening. An analytical approach to modelling flow under
345 these conditions is developed in Section 4. This is used to create predictive
346 models for both bulk flow rates and pollutant removal rates and identify
347 characteristic parameters.

348 *2.2. Performance in still air*

349 Below the system definition shown in Figure 2 are two branches that de-
350 scribe tests in still and moving air. Still-air tests represent the most basic
351 conditions in which air can flow through an opening, and represent the con-
352 ditions upon which the conventional airflow equations are based; see Figure

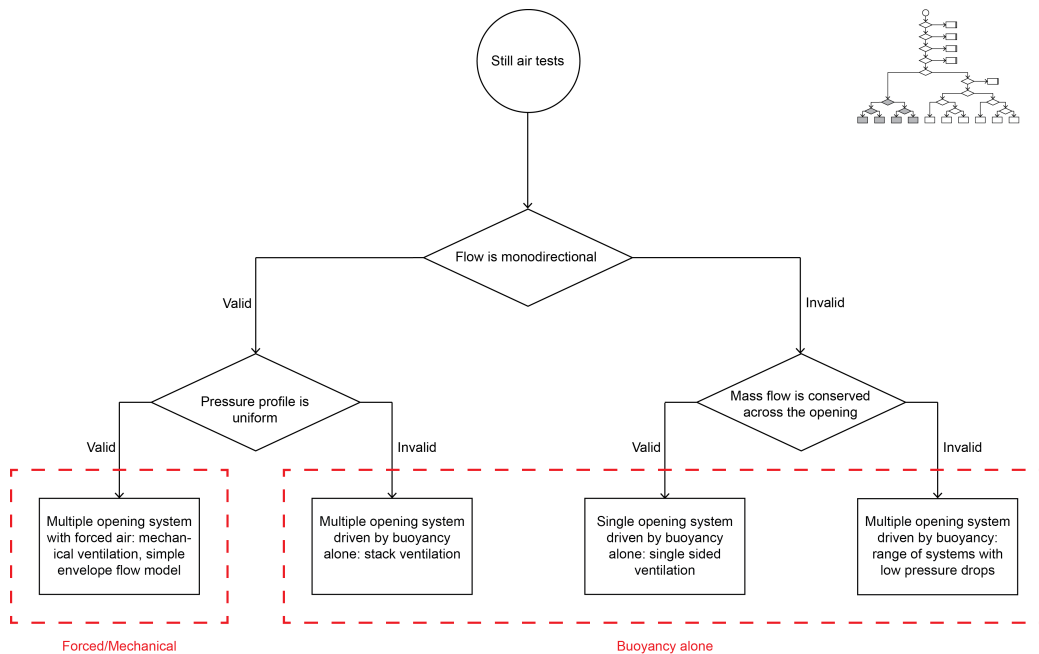


Figure 7: The assumptions that characterise still-air tests of window performance.

353 7. This assumes that the flow structure of the external environment is exclu-
 354 sively generated by airflow through the opening itself. Within this subset,
 355 flow can be conveniently divided into two types: monodirectional flow and
 356 bidirectional flow.

357 *2.2.1. Monodirectional flow*

358 Monodirectional flow is traditionally one of the simpler conditions to cal-
 359 culate. It represents stack or wind-driven ventilation where each opening
 360 acts exclusively as an inlet or an outlet.

361 *2.2.1.1. Uniform pressure profile.* The assumption that the pressure profile
 362 across an opening is uniform allows PPOs to be treated as point openings.
 363 This is the simplest set of conditions required for evaluating flow through an

364 opening, and is a key modelling assumption of the orifice flow equation [38].
365 These conditions represent forced or mechanical ventilation very well, but
366 they do not completely represent any real operating conditions of a naturally
367 ventilated building. It represents wind-driven conditions only when the wind
368 can be considered to be stationary at the building surface (well approximated
369 in the stagnation zone or in the lee of the building), and buoyancy conditions
370 when the neutral height is located an infinite distance from the opening.
371 However, it does represent some of the basic aerodynamic properties of an
372 opening upon which the effects of other factors can be analysed. For many
373 operating conditions it is likely to be a reasonable approximation of real
374 behaviour [12, 7].

375 Still-air tests that characterise airflow under these conditions are com-
376 monly used to determine the performance of components used in mechanical
377 ventilation systems, but are not common for natural ventilation openings.
378 This is partly because the larger dimensions of these openings require im-
379 practically large testing rigs, and the low pressures associated with natural
380 ventilation are hard to measure. These issues can be addressed to some ex-
381 tent using scale models [7]. Still-air tests characterising the performance of
382 real opening geometries are summarised in Section 2.1.1.

383 *2.2.1.2. Non-uniform pressure profile.* A difference in density between inter-
384 nal and external air results in a non-uniform pressure profile across its height
385 [3]. The uniformity of the pressure profile decreases as the neutral height ap-
386 proaches the window height, increasing the impact of this factor. For these
387 cases, the area distribution of the window is expected to have increased im-
388 portance.

389 Heiselberg *et al.* [38] present data for side hung windows suggesting the
390 discharge coefficient of an opening decreases when the temperature differ-
391 ence creates a non-uniform pressure profile. This is characterised by graphs
392 relating the measured discharge coefficient to a dimensionalised form of the
393 Archimedes number, given by

$$Ar' = \frac{\Delta T}{1000Q^2} \quad (6)$$

394 where ΔT is the temperature difference across the opening, and Q is the
395 volume flow rate through the opening. The presented data is specific to the
396 window geometry, opening angle, wall detail, and scale used in the experi-
397 ment, and therefore cannot be generalised to make performance predictions
398 for design.

399 Section 5 develops an analytical approach to describing this reduction
400 in discharge coefficient, and describes novel dimensionless parameters that
401 characterise this effect. The predictions of this model are compared against
402 the literature data in Section 5.1.

403 *2.2.2. Bidirectional flow*

404 Bidirectional flow is more complex than monodirectional flow. It is usu-
405 ally used to describe the ventilation of rooms with a single opening, but can
406 also occur when multiple openings are located at similar heights within a
407 façade or are substantially different in size. In still air, this represents the
408 buoyancy alone case.

409 *2.2.2.1. Mass conservation.* The most common assumption is that of mass
410 conservation across the opening – often simplified to volume conservation.

411 This flow pattern occurs where there is a single opening in a sealed room.
412 Bidirectional flow relies on a non-uniform pressure field to drive flow across
413 the opening, and so the distribution of the opening area is of great impor-
414 tance.

415 A theoretical evaluation of single opening, buoyancy driven ventilation
416 through a simple rectangular orifice can be made by integrating the orifice
417 flow equation over the height of the opening [10]. The pressure difference is
418 taken as a function of height, assuming that the neutral height occurs at the
419 centre of the opening. This results in the flow equation

$$Q = \frac{1}{3}C_dA_f\sqrt{\frac{\Delta\rho}{\rho}gh} \quad (7)$$

420 Several studies attempt to characterise buoyancy driven, single opening
421 ventilation through real windows. The experimental studies of side hung and
422 centre pivot windows of Warren and Parkins [35] present graphs of correction
423 factors to the theoretical airflow rate derived for a rectangular orifice as
424 a function of opening angle. This allows practitioners to account for the
425 geometry of these types of windows in a simple, unambiguous way. Compared
426 to analytical models, the impact of the height to width ratio is negligible.
427 Von Grabe *et al.* [34, 49] conduct similar experiments on a range of different
428 opening types, characterising the change in their performance as they open.
429 However, the performance curves are based on a potentially ambiguous *free*
430 area model, which could lead to application errors. The authors introduce the
431 idea of the thermal height of the window, providing a convincing analytical
432 explanation for the difference in the performance of different window types.
433 Their data suggests that air-flow rates through horizontal pivot windows,

434 double sliding sash windows and side hung windows increase rapidly as the
435 window is opened, suggesting they are useful for summer overheating or purge
436 ventilation [42]. In contrast, air-flow rates through top and bottom hung
437 windows increase more slowly as the window is opened, offering a greater
438 degree of control that may be more useful when ventilating for indoor air
439 quality in the winter [42].

440 Wilson and Kiel [40] identify that the ventilation rate depends on the de-
441 gree of interfacial mixing between the inflow and outflow streams. ASHRAE
442 present an equation for predicting the discharge coefficient due to this effect
443 as a function of the temperature difference across the opening ΔT [10], given
444 by

$$C_d = 0.4 + 0.0045\Delta T \quad (8)$$

445 The mixing effect is reduced at high temperature differences, and in-
446 creased by local atmospheric turbulence [40]. This suggests that experiments
447 performed in still-air could overestimate the pollutant removal rate an open-
448 ing provides when installed in a turbulent environment. It is unclear how
449 the choice of opening type affects the degree of interfacial mixing.

450 *2.2.2.2. Unbalanced flow.* Where mass flow is not conserved across an open-
451 ing, flow patterns are more complex, and the system requires one or more
452 additional airflow paths. This scenario commonly occurs where mechanical
453 extract is used in conjunction with single sided ventilation (in bathrooms or
454 kitchens), or where windows of different sizes are open simultaneously. Stud-
455 ies need to characterise both inflow and outflow rates as the neutral height is

456 varied across the opening. To the best of our knowledge, there is no research
457 into the performance of real windows under this regime, either *in-situ* or in
458 controlled conditions. Airflow network models, such as CONTAM [50], apply
459 theoretical models to describe this kind of flow through simple 2D openings.
460 To do this, the orifice equation is integrated over the height of the opening,
461 and the neutral height is varied until mass flow conservation is achieved. A
462 similar approach is used by Jones *et al.* to model infiltration in the presence
463 of mechanical extract [51].

464 2.3. Performance in moving air

465 When the wind drives flow, the air proximate to the external surface of an
466 opening can be expected to be in motion for the majority of cases. Moving
467 air tests can be used to investigate both the impact of wind alone, and of
468 wind and buoyancy combined. The analytical framework shown in Figure 8
469 breaks down the flow configurations accordingly, which are then subdivided
470 into monodirectional and bidirectional flow. For convenience, the impact of
471 non-uniform wind pressures has been given as a separate factor that can be
472 applied to flows driven by wind alone and those driven by wind and buoyancy
473 combined.

474 In many cases the aerodynamic properties measured in the presence of
475 wind may be similar to those measured in still-air tests, but significant dif-
476 ferences are also possible. The presence of external air movement can alter
477 the shape and directions of streamlines passing through the opening, and
478 projecting opening geometries can interact with external flows to alter the
479 pressure field near the surface of the opening; see Figure 9. Venturi ventila-
480 tors have been used to ensure suction pressures over outlets, with pressure

481 coefficients as low as -1 being achieved [15, 42] Here, pressure coefficients
482 derived from bluff bodies will not be suitable for use with real window ge-
483 ometries. This interaction also affects the surface pressures for some distance
484 around the opening, introducing additional uncertainties if other openings
485 are present [52]. Surface pressures and flow patterns are also sensitive to the
486 presence of buildings and trees in the immediate environment [53], meaning
487 environmental conditions can diverge from those evaluated using simple de-
488 sign techniques. In addition to this, the turbulent flow structures and other
489 unsteady behaviour associated with atmospheric wind have been proposed
490 as another mechanism for driving ventilation [35, 7, 14]. However, unsteady
491 flows are not readily compatible with envelope flow models, and are beyond
492 the scope of this paper.

493 *2.3.1. Impact of non-uniform wind pressures*

494 It is commonly assumed that wind pressures acting on an opening are
495 uniform across its surface. While this assumption is largely valid for small
496 openings, many authors state that it is likely to break down when the open-
497 ings are very large compared to the area of the façade [4, 7]. This occurs
498 because the pressure coefficients vary across a façade as a function of building
499 geometry and wind angle. It is thought that this variation could become the
500 dominant driving force behind ventilation in some configurations, but could
501 also reduce ventilation rates by the same mechanism described for thermal
502 buoyancy; see Section 2.2.1.2.

503 Non-uniform pressure profiles also occur due to interactions with complex
504 opening geometry. Iqbal *et al.* [52] find that airflow passing over a centre
505 pivot window can generate variations in static pressure across its surface

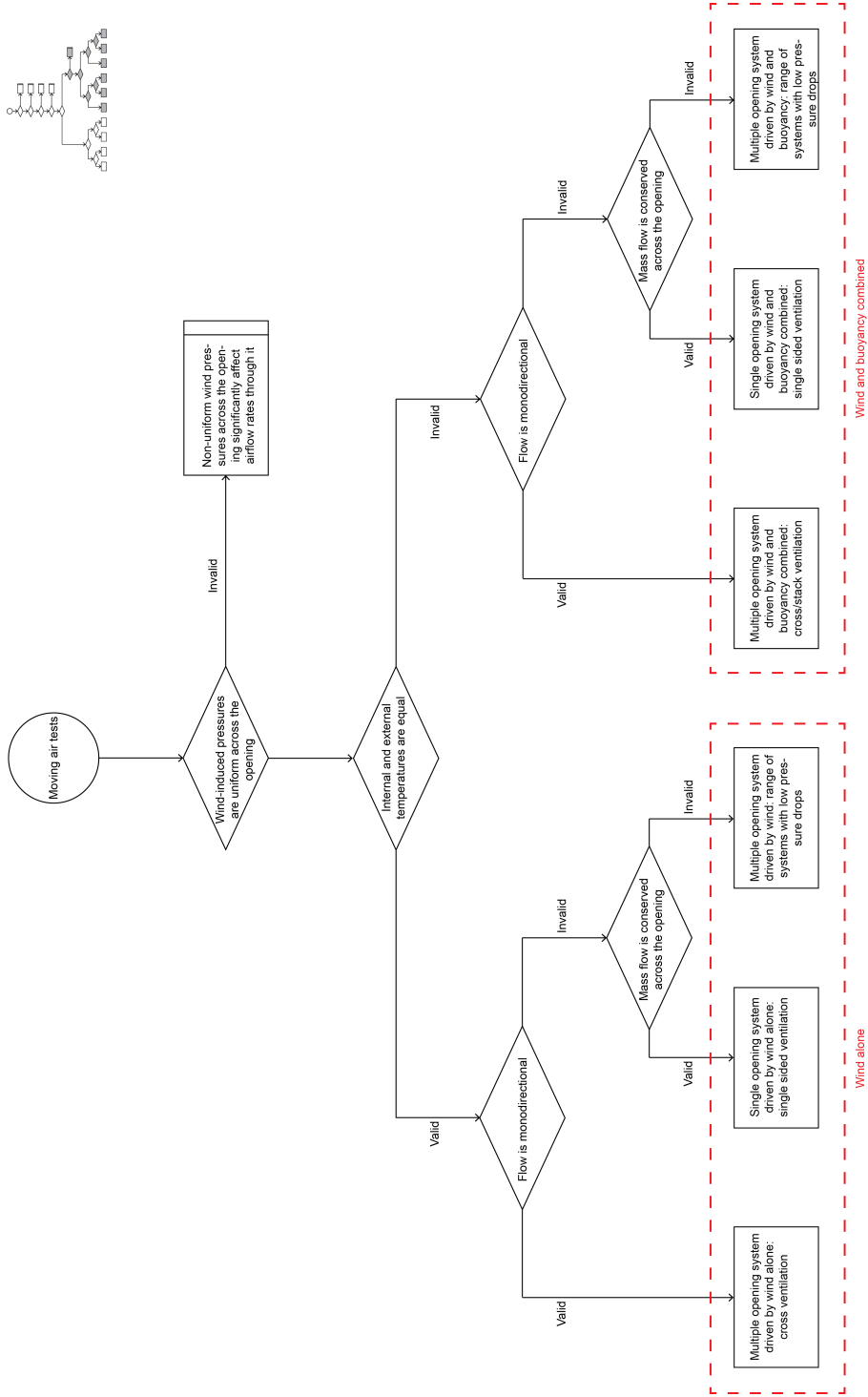


Figure 8: The further assumptions that characterise moving air tests of window performance.

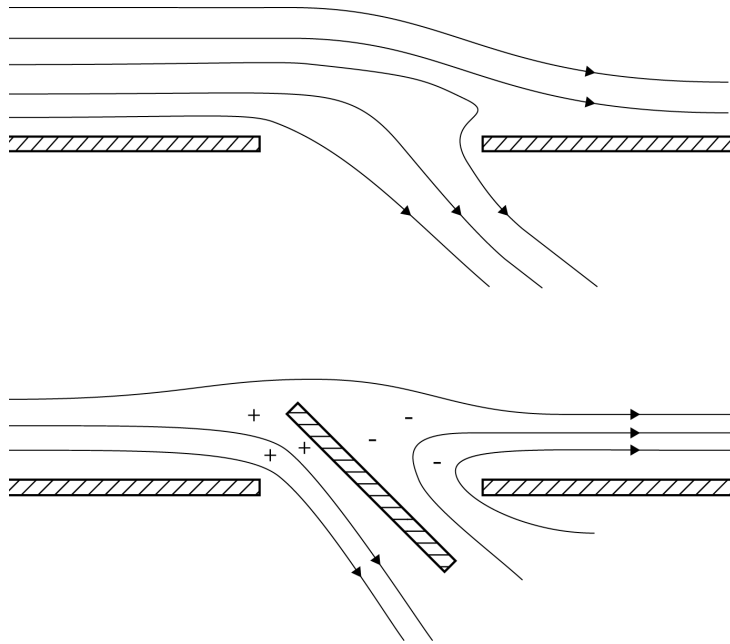


Figure 9: Diagram of the effect projecting opening geometry can have on flow patterns.

506 that are sufficient to drive steady-state bidirectional flow; see Figure 9. Air
 507 enters at the windward side of the opening, and leaves via the leeward side.

508 To the best of our knowledge there is no research studying the inter-
 509 action between non-uniform wind pressures and the non-uniform pressures
 510 generated by thermal buoyancy.

511 2.3.2. *Internal and external temperatures are equal*

512 A common simplification of wind-driven flow is that the internal and ex-
 513 ternal temperatures are equal. This represents wind alone conditions, where
 514 there is no contribution of buoyancy to ventilation rates. Although these
 515 conditions may only occur transiently in operation, this greatly simplifies
 516 the experimental treatment of the impact of external wind on PPOs. In

517 many cases, the behaviour of an opening exposed to wind alone is expected
518 to be a reasonable approximation for its behaviour in the presence of wind
519 and buoyancy forces combined.

520 *2.3.2.1. Monodirectional flow.* Envelope flow models commonly assume that
521 airflow through an opening is driven by the static pressure at the building's
522 surface, and the dynamic pressure makes no contribution to airflow rates.
523 However, Vickery and Karakatsanis [12] find that the orifice flow equation
524 systematically overestimates flow rates in the presence of external wind, and
525 the error increases as the wind angle normal to the façade increases.

526 The influence of external wind can be investigated analytically by con-
527 sidering airflow along a streamline as it enters a building; see Figure 10. It
528 is assumed that the wind induces air motion parallel to the building surface,
529 and that this air stream acts as the source of air that passes through the
530 opening [12, 35, 54]. Balancing total pressures along the streamline results
531 in an equation for airflow rate through the opening [7].

$$Q = C_d A_f \sqrt{\frac{2(\Delta P + \frac{1}{2}\rho U_L^2)}{\rho}} \quad (9)$$

532 where U_L is the local wind speed parallel to the opening. This is different
533 from the orifice flow equation, and considers the contribution of dynamic
534 pressure to airflow through the opening. One would expect the discharge
535 co-efficient defined using Equation 9 to be highly dependent on U_L , as any
536 conserved momentum in the cross flow acts to reduce the minimum area
537 through which the air passes. The influence of this on mass flow rates will,
538 to some extent, be balanced by the increased velocity of the flow owing to

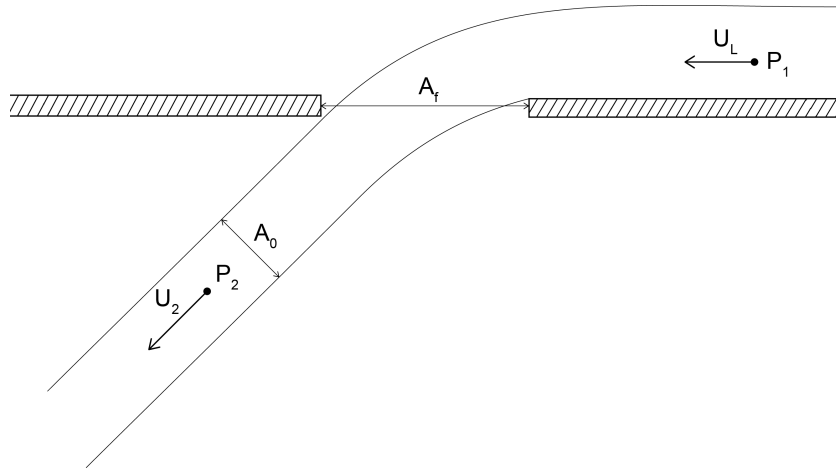


Figure 10: Diagram depicting wind induced flow through an orifice

539 the dynamic pressure term.

540 The influence of wind has been observed to cause a substantial reduction
 541 in the discharge coefficient calculated using the orifice flow equation, even
 542 for sharp edged orifices [23, 4, 55]. Kurabuchi *et al.* [54] and Obha *et al.*
 543 [55] characterise this behaviour experimentally using a dimensionless pres-
 544 sure coefficient defined as the ratio of the static pressure difference acting
 545 across the opening to the dynamic pressure in the external flow. The data is
 546 presented both graphically and through the use of an approximate curve fit,
 547 allowing the influence of wind on discharge coefficient to be integrated into
 548 simple computational models. Applying this data to envelope flow models
 549 can greatly reduce their error, but requires extensive data sets if the method
 550 is to be implemented to characterise the full range of real window geometries.

551 External air motion can interact in a range of interesting ways with real
 552 opening geometries. Kurabuchi *et al.* [54] identify some cases where discharge
 553 coefficients rise dramatically as the dynamic pressure in the wind becomes

554 large when compared to the static driving pressures. Etheridge [7] also iden-
555 tifies discharge coefficients for wind cowls rising towards infinity as the wind
556 velocity becomes very large. This may be caused by the opening interacting
557 with external wind flows to evolve static pressure at the opening surface.
558 They demonstrate how a characterisation of this effect can be significantly
559 improved using Equation 9, which integrates the dynamic pressure of the
560 external wind into the orifice equation.

561 *2.3.2.2. Bidirectional flow.* Bidirectional flow driven by wind is complex.
562 Several different mechanisms are proposed, and the volume flow rate cal-
563 culations used in practice are based on simple empirical correlations. The
564 correlations given in the CIBSE Guides and Manuals [3, 8], and the Euro-
565 pean standard EN16798-7:2017 [56], are limited to fully open windows, where
566 mass flow is conserved across the opening. They also take no account of the
567 impact of wind direction.

568 The bidirectional flow ventilation calculations presented in CIBSE Guide A
569 are based on the research of Warren and Parkins [35], who recommend calcu-
570 lating the effects of wind and buoyancy separately, and then taking the larger
571 value. For wind-driven ventilation, they model mixing across the boundary
572 of an opening due to the turbulent shear layer that forms when moving air
573 passes a region of stationary flow. A simple empirical correlation with refer-
574 ence wind speed is given to characterise a minimum flow rate to be used for
575 the sizing of openings. They also present data describing how the flow rate
576 changes for different turbulence scales and graphs of corrections for side hung
577 windows as they open and close. Ventilation rates are characterised using a
578 non-dimensional flow number F_L , defined using the velocity of the flow at the

579 building's surface. When applied to real buildings, this requires the distri-
580 bution of wind velocities on a building's surface to be known. Kurabuchi *et*
581 *al.* [54] describe simple wind tunnel techniques for measuring and presenting
582 dynamic pressure distributions, but not their direction.

583 An alternative transient method of wind-driven bidirectional proposed is
584 pulsation theory [57]. Where room volumes are large, significant volumes
585 of air can be driven into a space by fluctuating pressures at an opening,
586 without significantly pressurising the space. This leads to ventilation rates
587 that are dependent on the volume of the internal space, and the magnitude
588 and frequency of external pressure fluctuations. The use of their calculation
589 methodology is inhibited by a lack of available design data.

590 The data presented by Warren and Parkins [35] is not widely available,
591 but could be of immediate practical use to designers. More recent attempts to
592 model more complex flow mechanisms [58] and a range of opening geometries
593 [59] do not improve on the model presented by Warren and Parkins [35], as
594 they have errors of a similar magnitude that do not justify the increased
595 model complexity.

596 *2.3.2.3. Multiple openings.* Much of the literature that underpins best-practice
597 standards assumes that ventilation systems comprising multiple openings on
598 a single wall can be adequately described by treating each opening in iso-
599 lation. However, there is evidence that when multiple openings exist on a
600 façade, mass flow rates can be greater than those predicted by the single
601 opening equations [35, 21, 14]. This is primarily due to differing local pres-
602 sure coefficients between any two openings driving flow. Here, mass flow
603 rates cannot be said to be conserved through each window, and surface av-

604 eraged pressure coefficients cannot justifiably be used. This might simplify
605 the flow through the opening to the monodirectional flow case, but it is also
606 possible that bidirectional mechanisms occur where differential pressures are
607 sufficiently small [8].

608 *2.3.3. Wind and buoyancy combined*

609 Sections 2.2.1.2 and 2.2.2.1 show that the uneven pressure profile asso-
610 ciated with a temperature difference across an opening can both decrease
611 monodirectional flow rates and increase bidirectional flow rates. To the best
612 of our knowledge there is no research studying wind-driven monodirectional
613 flow in the presence of buoyancy forces, or how this may impact predictions
614 of volume flow rates. However, the internal flows generated by wind forces
615 exceed those due to thermal buoyancy, even in light winds [12]. This suggests
616 the influence of buoyancy on the aerodynamic properties of an opening in
617 the presence of wind may be small.

618 *2.3.3.1. Bidirectional flow.* The European standard used to predict ventila-
619 tion rates through single openings [56] is based on the correlations of De Gids
620 and Phaff [60], which seek to account for the effect of wind and buoyancy
621 combined. The simplified equations they produced are used as the basis for
622 further research by Larsen and Heiselberg [20], who account for the wind
623 direction to reduce the error in the model from 29% to 23%. However, the
624 use of this equation requires information about the variation in wind pressure
625 across the surface of the opening - which will vary with building geometry,
626 opening location and wind direction - making it of less practical use in the
627 early design of a naturally ventilated building.

Environmental conditions / flow configuration	Predictive equations	Source
Still-air, monodirectional, uniform pressure profile [Forced air, multiple opening]	$Q = C_d A \sqrt{\frac{2\Delta P}{\rho}}$ $C_d = f(Re)$ $Q = C\Delta P^n$	ASHRAE Fundamentals [10], CIBSE Guides A and AM10 [8, 3], CONTAM [50] Etheridge [7], Bailey <i>et al.</i> [44] CONTAM [50], Sherman [45]

Table 1: Predictive equations under conditions for forced ventilation - often used to describe monodirectional flow under any driving force.

628 While the predictive equations used in practice assume the forces of buoyancy
629 and wind act constructively, Caciolo *et al.* [28] identify cases where
630 the interaction of wind reduces the ventilation rate expected from buoyancy
631 alone. This reinforces the measurements of Kiel and Wilson [40], who show
632 that interfacial mixing by the wind can reduce ventilation efficiency.

633 2.4. Extent of knowledge

634 The framework set out in Sections 2.1 – 2.3 breaks down a range of
635 characteristic environmental conditions that drive flow through PPOs and
636 the range of flow structures that can occur within these openings. As the
637 mechanisms vary, so do the equations that describe flow through them. A
638 range of equations given in the literature to describe these flow scenarios are
639 given in Tables 1–4.

640 In principle, each flow scenario needs to be characterised with its own
641 testing regime, and the degree to which it can be described using still-
642 air discharge coefficients assessed. In practice this is very rarely achieved.
643 EN 13141-1 standardises still-air pressurisation tests of PPOs [17], but this
644 does not require parametisation that would enable modeling under a range of
645 ventilation pressures. Similarly, European technical standards specify test-
646 ing regimes for roof outlets in the presence of wind [62, 43], but these do

Environmental conditions / flow configuration	Predictive equations	Source
Still-air, monodirectional, non-uniform pressure profile [buoyancy alone, multiple opening]	$Q = C_q C_d A \sqrt{\frac{2\Delta P}{\rho}}$ where $C_q = f\left(\frac{2\Delta P}{\Delta\rho gh}\right)$	Proposed in this paper; see Section 5
Still-air, bidirectional, flow-conservation [buoyancy alone, single opening]	$Q = \frac{1}{3} C_d A \sqrt{\frac{\Delta\rho}{\rho} gh}$ where $\frac{\Delta\rho}{\rho} = \frac{\Delta T}{T}$ and $C_d = 0.4 + 0.0045\Delta T$	Warren and Parkins [35], ASHRAE Fundamentals [10], CIBSE guides A and AM10 [8, 3] ASHRAE Fundamentals [10], Wilson and Kiel [40]
Still-air, bidirectional, unbalanced flow [buoyancy alone, multiple opening]	$Q_{out} = w C_d \int_{z_n}^{h_T} \sqrt{2\Delta P(z)} dz$ $Q_{in} = w C_d \int_{h_B}^{z_n} \sqrt{\Delta P(z)} dz$	CONTAM [50]

Table 2: Predictive equations for buoyancy only ventilation

Environmental conditions / flow configuration	Predictive equations	Source
Moving-air, uniform temperature, monodirectional [Wind alone, multiple opening]	$Q = C_d A \sqrt{\frac{2\Delta P}{\rho}}$ where $C_d = f\left(\frac{2\Delta P}{\rho U_L^2}\right)$ $Q = C_d A \sqrt{\frac{2(\Delta P + \frac{1}{2}\rho U_L^2)}{\rho}}$ where $C_d = f\left(\frac{U_L A}{Q}\right)$	Kurabuchi <i>et al.</i> Etheridge [7], Chiu and Etheridge [23]
Moving-air, uniform temperature, bidirectional, flow conservation [wind alone, single opening]	$Q = 0.025 A U_R$ $Q = F_L A U_L$ $Q = C_d A \sqrt{U_R^2 - \frac{2\gamma P_a}{\rho V} \omega}$	CIBSE guides A and AM10 [8, 3], Warren and Parkins [35] Warren and Parkins [35] Cockroft and Robertson [57]
Moving-air, uniform temperature, bidirectional, unbalanced flow [Wind alone, multiple opening]	$Q = A^* U_R \sqrt{0.32\Delta C_p + 0.09\sigma_{\Delta C_p}}$	Daish <i>et al.</i> [14]

Table 3: Predictive equations for wind alone ventilation

Environmental conditions / flow configuration	Predictive equations	Source
Moving air, temperature difference, mono-directional flow [Wind and buoyancy combined, multiple openings]	–	–
Moving air, temperature difference, bidirectional flow, flow conservation [Wind and Buoyancy combined, single opening]	$Q = \frac{1}{2} C_d A \sqrt{C_1 U_{10}^2 + C_2 h \Delta T + C_3}$ $Q = \frac{1}{2} C_d A \sqrt{\max(C_1 U_{10}^2; C_2 h \Delta T)}$ $Q = A \sqrt{a \Delta T + b U_R^2}$	De Gids and Phaff [60], BS EN 15242:2007 [61] BS EN 16798-7:2017 [56] ASHRAE Fundamentals [10]
Moving air, temperature difference, bidirectional flow, flow conservation, non-uniform wind pressure [Wind and Buoyancy combined, single opening]	$Q = A \sqrt{C_U + C_T + C_{\Delta P}}$ $C_U = C_1 C_P U_R^2$ $C_T = C_2 h \Delta T$ $C_{\Delta P} = C_3 \Delta C_{P(\text{opening})} \frac{\Delta T}{U_R^2}$	Larsen and Heiselberg [20]
Moving air, temperature difference, bidirectional flow, unbalanced flow [Wind and Buoyancy combined, multiple openings]	–	–

Table 4: Predictive equations for wind and buoyancy combined ventilation

647 not yield parameters suitable for modelling. No similar standard is found for
 648 inflow openings, or for openings in walls.

649 Standardised test methods to evaluate the aerodynamic properties of
 650 PPOs are largely absent for a range of driving mechanisms. As a result, data
 651 sets provided by manufacturers cannot confidently be applied for a range of
 652 design conditions. However, there is scope within the existing literature to
 653 derive such tests. The experimental procedures of Warren and Parkins [35]
 654 could be used as the basis for standardised tests for buoyancy driven and
 655 wind driven bidirectional flows through PPOs in the single opening configu-
 656 ration. Similarly, the procedures developed by kurabuchi *et al.* could be used
 657 to standardise performance tests of monodirectional flow in the presence of
 658 wind. For many categories of experimental conditions identified within the
 659 framework, academic research does not yet provide adequate procedures to

660 evaluate the aerodynamic performance of specific PPOs; see Sections 2.1.3,
661 2.1.4, 2.2.2.2, 2.3.1, 2.3.2.3, and 2.3.3.

662 The use of sharp-edged, flush orifices are ubiquitous in investigations of
663 ventilation phenomena, but academic data characterising the aerodynamic
664 performance of specific PPOs, or types of PPO, is scarce. Data is available
665 for still-air discharge coefficients of hinged openings [31, 32, 33]; buoyancy
666 alone, single opening discharge coefficients for side hung and horizontal pivot
667 openings [35]; and wind alone, single opening discharge coefficients for hinged
668 openings [35]. Predictive equations describing still-air performance for airflow
669 through insect mesh [44], and combinations of mesh and louvres [42] are
670 available in the literature. This study identifies no sufficiently comprehensive
671 data sets for other opening types or flow configurations.

672 3. Statistical Effective Area Model

673 To address the failings of *free* area models discussed in Section 2.1.1,
674 a Statistical Effective Area Model (SEAM) has been created. This model
675 is based on that proposed in BB101 [41], and fit using academic data for
676 hinged openings [31, 32, 33]. The discharge coefficient is defined according
677 to Equation 5, and described for a constant aspect ratio by

$$C_d(\theta) = B (1 - e^{-M\theta}) \quad (10)$$

678 where B and M are coefficients that can be fit to experimental data. The
679 fitted coefficients B and M are plotted as a function of aspect ratio, and
680 described by empirical correlations given by

$$B = 0.18e^{-0.78(\sigma)} + 0.61 \quad (11)$$

$$M = 0.016(\sigma + 1) \quad (12)$$

681 Although this model is defined for top or bottom hung openings, rotational
 682 symmetry enables the same model to be used for side hung openings when
 683 the inverse aspect ratio $w:h$ is substituted.

684 3.1. Comparing the performance of different area models

685 To quantify the errors associated with each modelling technique, we have
 686 compared the discharge coefficients predicted by the *free* area models, the
 687 BB101 online calculator, and *SEAM* with experimental data from the liter-
 688 ature [31, 32, 33]; see Figure 11. This shows that that the ambiguity of *free*
 689 area models can lead to significant variations in predicted performance that
 690 can either under or over-estimate airflow rates. The predicted values of the
 691 discharge coefficient are calculated using Equation 4 for the range of height
 692 to width ratios and opening angles present in the literature, substituting the
 693 *effective* areas predicted using each modelling technique. The *effective* area
 694 predicted by the *free* area models are calculated using a discharge coefficient
 695 of 0.61, although 0.65 is also common [10].

696 The difference between model predictions and experimental data comes
 697 from systematic error caused by a poorly fitting model and random error in
 698 the experimental data. These errors can be combined to assess the confidence
 699 in each model when used to predict the aerodynamic performance of an
 700 opening.

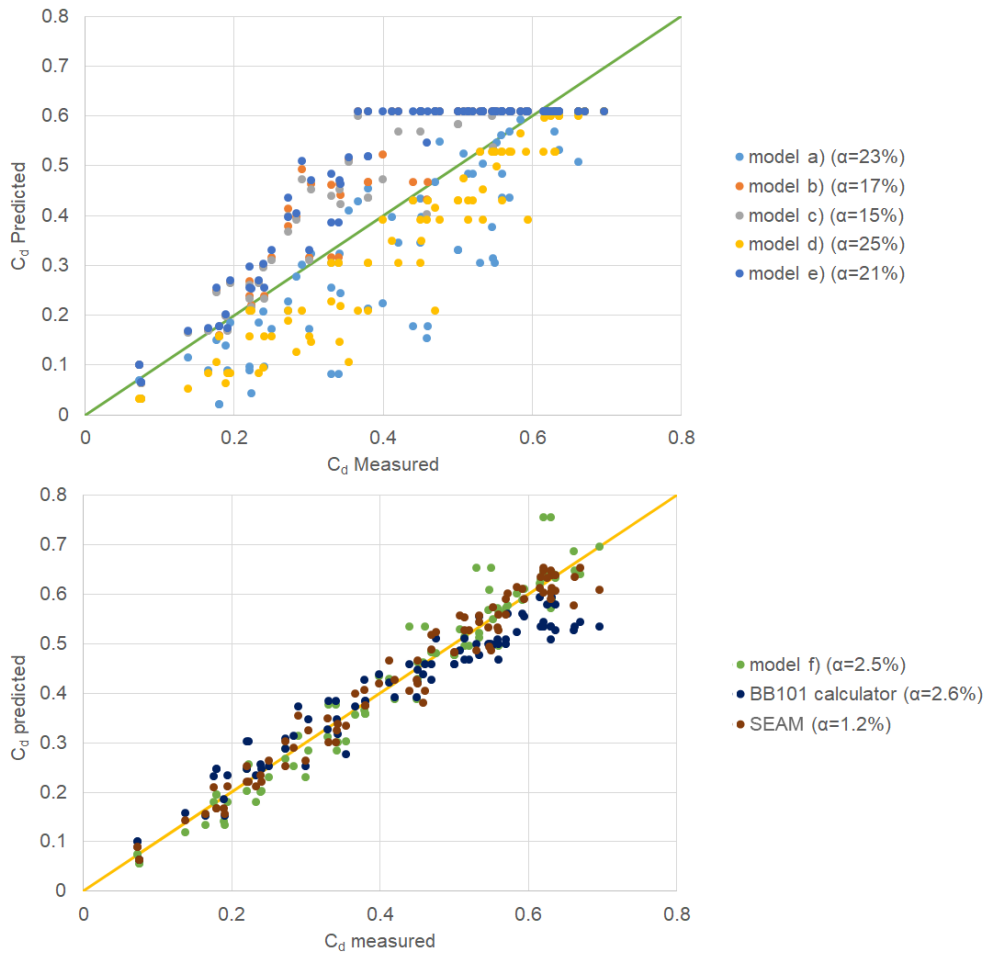


Figure 11: Relationship between predicted and measured values of the discharge coefficient (as defined in Figure 6) for a range of predictive models and their standard error, α .

Top, purely analytical *free* area models; Bottom, semi-empirical models (see Figure 5 for descriptions).

701 If the deviation from the model is entirely random there is an equal chance
702 of deviations occurring above or below the model predictions, and therefore
703 the mean deviation would be zero. The *mean* systematic error can therefore
704 be estimated from the mean deviation. It must be stressed that this is a
705 *mean* systematic error, and a model that overestimates in some regions and
706 underestimates in others may have an artificially low *mean* systematic error.
707 This is to some extent compensated for by an increase in random error.
708 Some height to width ratios and opening angles result in greater systematic
709 errors than others, which can be as large as 80%. In addition, while the
710 percentage deviation between model and data is normally distributed for the
711 quasi-empirical models, this is not true for the purely analytical models.

712 The model proposed in this paper - *SEAM* - fits the data the best, and can
713 predict opening performance with a standard error of 1.2%. Free area model *d*
714 is used in the safety-critical application of smoke ventilation, and is the only
715 model that systematically underestimates aerodynamic performance. This
716 will result in the specification of openings that outperform design predictions,
717 and therefore the model does not need updating urgently. *SEAM* will be
718 included in an updated BB101 calculator. The model is based on data from
719 large openings where $t/h \ll 1$ (see Figure 6), so an analytical model based
720 on geometric similarity of the *free* area has been included in the calculator for
721 smaller or thicker openings where the opening thickness cannot be neglected.

722 This analysis shows that purely analytical *free* area models cannot be
723 applied with confidence to predict the aerodynamic performance of PPOs.
724 Predictive models created to support system design must be calibrated with
725 empirical data for the range of geometric parameters within which it will be

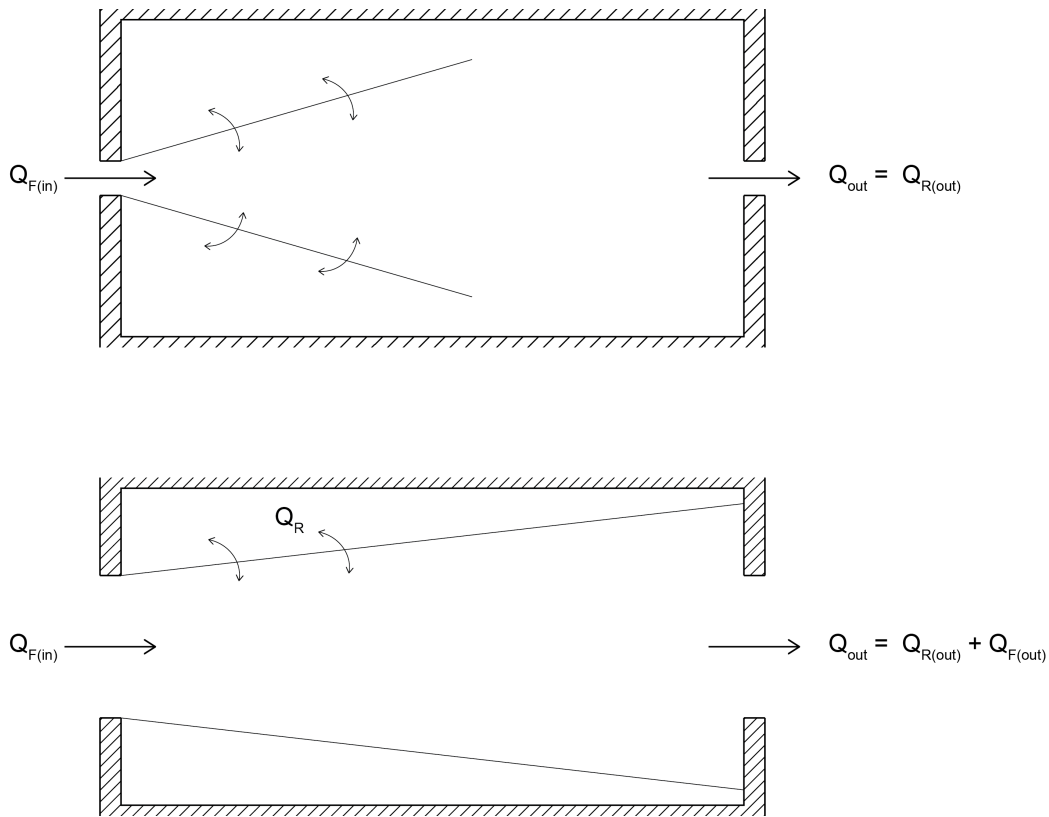


Figure 12: An illustration of the difference between an expanding jet from a small opening (top) and a large opening (bottom).

726 applied.

727 4. Modelling the impact of flow connection

728 Section 2.1.4 identified evidence of increases in bulk airflow rate that
 729 occur when inflow and outflow openings are closely aligned. This effect may
 730 be caused by conservation of kinetic energy between the inflow jet, or it may
 731 be associated with a change in streamline shape in the approach to the outlet.
 732 Moreover, it is unclear whether this effect is beneficial, as it is possible that

733 this flow connection inhibits the removal of pollutants from the wider space.

734 We have developed a simplified analytical approach to this question by
735 looking at the stream tube formed between two aligned openings in a large
736 space; see Figure 12. In this model, we treat the inflow air as a free jet,
737 which expands as it entrains air on its journey towards an outlet located on
738 the opposite wall. When the inlet is small, the inflow jet mixes thoroughly
739 with the room air before being extracted and so the extracted air can be
740 considered to be made up entirely of room air. The kinetic energy is also
741 completely dissipated, resulting in still-air conditions at the surface of the
742 outflow opening. Conversely, when the inlet is large, much of the fresh air
743 leaves via the outlet without mixing with the room air, and the jet reaches
744 the outlet with a significant velocity. Room air can be removed from the
745 space only by entrainment into the jet.

746 The effect of flow connection on a natural ventilation strategy can be
747 broken down into two key phenomena; the increase in bulk airflow rate due
748 to conservation of kinetic energy and the reduction in ventilation effectiveness
749 [8], E_v , caused by short circuiting of fresh air.

750 A simple model for estimating bulk airflow rates can be made using a
751 modified envelope flow model, which allows a proportion of the dynamic
752 pressure in the inflow jet to be conserved to drive air through the outflow
753 opening. The dynamic pressure in the jet available to drive airflow can be
754 evaluated using the entrainment equations for ideal free jets [19]. Assum-
755 ing the discharge coefficients of the openings are unchanged by the altered
756 streamlines, a dimensionless volume flow rate can be evaluated

$$\frac{Q}{Q_E} \sqrt{\frac{1 + A^{*2}}{1 + A^{*2} - 16 \frac{A_{eff(1)}}{x^2}}} \quad (13)$$

757 where Q_E is the volume flow rate predicted by conventional envelope flow
 758 models, A^* is the ratio of the *effective* area of the inlet $A_{eff(1)}$ to the *effective*
 759 area of the outlet $A_{eff(2)}$, and x is the distance between the two openings.

760 The ventilation effectiveness [8] can be defined as the proportion of room
 761 air in the jet at the outlet. Similarly, this can be evaluated using the entrain-
 762 ment equations for free jets [19], giving

$$E_v = \frac{Q_{R(out)}}{Q_{out}} = 1 - 4 \frac{\sqrt{A_{eff(1)}}}{x} \quad (14)$$

763 where $Q_{R(out)}/Q_{out}$ is the proportion of room air extracted from the space.
 764 The effective ventilation rate of room air can be calculated as the product
 765 of the ventilation effectiveness and the volume flow rate. Equations 13 and
 766 14 suggest that the relevant dimensionless parameter is the ratio $\sqrt{A_{eff}}/x$,
 767 rather than the commonly favoured opening porosity [4].

768 4.1. Comparison with literature data

769 The predictions of these equations can be compared against the data
 770 presented by Seifert *et al.* [48]. They present a CFD study of a 6m cube,
 771 where the area of the inlet and outlet are gradually increased. By applying
 772 the model to this data, the mass flow rate of room and fresh air can be
 773 plotted as opening area is increased; see Figure 13. Once flow connection
 774 has been formed, the rate at which room air is removed drops and is not
 775 sufficiently offset by increasing flow rate of fresh air. This contrasts with
 776 conventional wisdom that larger airflow rates imply higher pollutant dilution

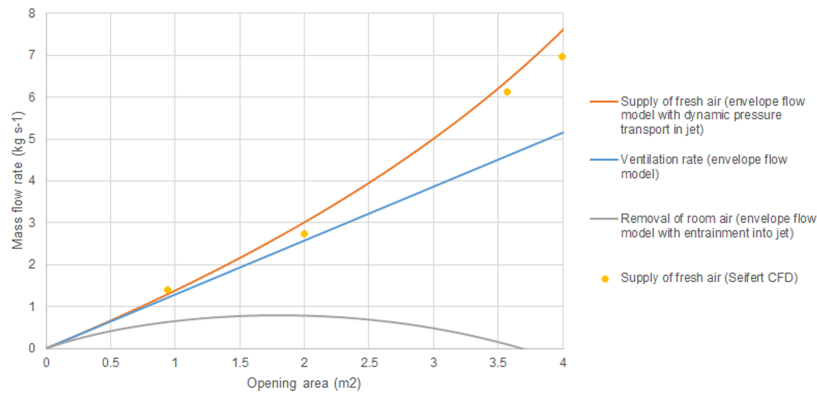


Figure 13: The predicted ventilation rate of room air has been plotted alongside CFD data presented by Seifert [48].

777 rates and provide more effective ventilative cooling. It also demonstrates the
 778 continuous predictions made by combining free jet and envelope flow models
 779 agree well with the discrete data points produced by the CFD. This suggests
 780 that the increase in volume flow rate is caused primarily by transmission of
 781 dynamic pressure in the jet, and not by a reduction in discharge coefficient.
 782 The agreement is surprising given the scale of the difference in complexity
 783 and computation time between the two models.

784 Counter-intuitively, this model suggests that, under certain circumstances,
 785 increasing the open area can reduce pollutant removal from a space. Ven-
 786 tilation strategies should be designed to prevent flow contact between the
 787 inflow jet and the outflow opening. This can either be achieved by interfer-
 788 ing with the transmission of the jet through the space, or by manipulating
 789 the openings to adjust the size, velocity and direction of the inflow jet. These
 790 parameters represent a set of aerodynamic properties that need to be charac-
 791 terised for different opening types, beyond merely their resistance to airflow.

792 Knowledge of these properties would be useful when designing for thermal
793 comfort, as well as ensuring contact with thermal mass.

794 In real buildings, the reduction in pollutant removal rates is likely to be
795 less severe than predicted by the model, as three-dimensional opening geome-
796 try, buoyancy, and internal obstacles will interfere with the clean propagation
797 of the jet and encourage mixing. However, formation of a wall jet caused by
798 locating openings near ceilings would reduce the entrainment coefficient [19],
799 which could cause pollutant removal rates to be lower than predicted.

800 5. Modelling the impact of buoyancy-induced non-uniform pres- 801 sure profiles

802 A reduction in the discharge coefficient that occurs in buoyancy driven
803 ventilation is identified in Section 2.2.1.2. This might occur because the
804 non-uniform pressure profile associated with a high temperature difference
805 invalidates the point area assumption used in the orifice equation. If this is
806 the case, it should be possible to evaluate flow rates analytically using an
807 area profile for the opening.

808 Side hung windows do not have a uniform area profile, and so the paths of
809 least resistance are at the top and bottom of the window. In order to analyse
810 a worst case scenario, it is assumed that the measured *effective* area can be
811 represented by two equal, point openings at the window's extremities; see
812 Figure 14. Evaluating the flow through these two openings yields a correction
813 factor, C_q , that can be applied to the conventional orifice flow equation, where

$$C_q = \frac{1}{2} \left(\sqrt{1 + \frac{1}{P^*}} + \sqrt{1 - \frac{1}{P^*}} \right) \quad (15)$$

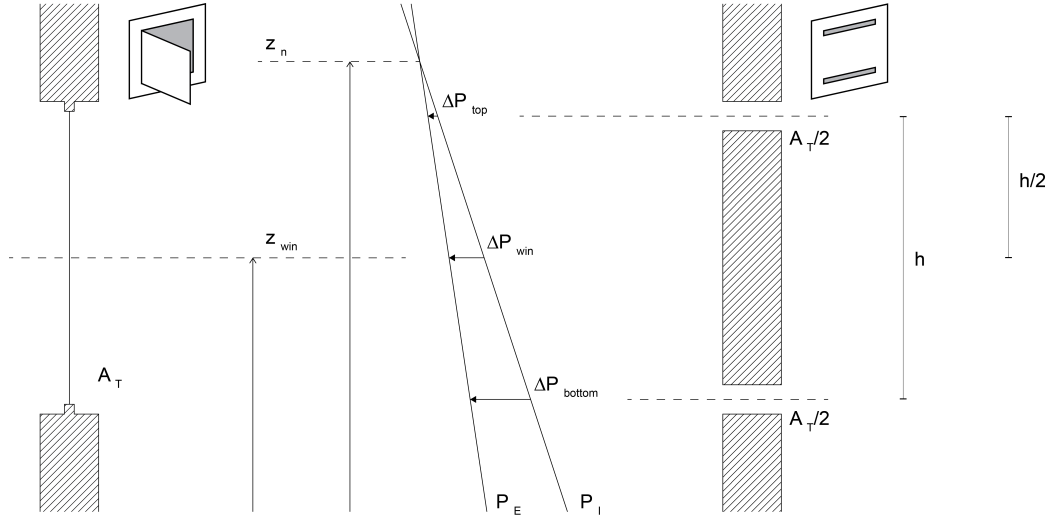


Figure 14: Diagram of a side hung window (left) against the model proposed to evaluate flow through it (right)

814 This derivation defines a dimensionless pressure ratio P^* that characterises
 815 the flow regime through the opening and is independent of building geometry,
 816 given by

$$P^* = \frac{2 \Delta P_{win}}{\Delta \rho g h} \quad (16)$$

817 where ΔP_{win} is the measured pressure difference across the centre of the
 818 window, $\Delta \rho$ is the density difference between indoor and outdoor air and h is
 819 the height of the opening; see Figure 14. Flow is monodirectional when $P^* \geq$
 820 1 or ≤ -1 , and bidirectional where $-1 \leq P^* \leq 1$. Still-air pressurisation
 821 tests describe behaviour where $P^* \rightarrow \pm \infty$.

822 For buoyancy only ventilation, the non-dimensional pressure can be shown
 823 to be equivalent to a non-dimensional height, h^* , which allows the correction

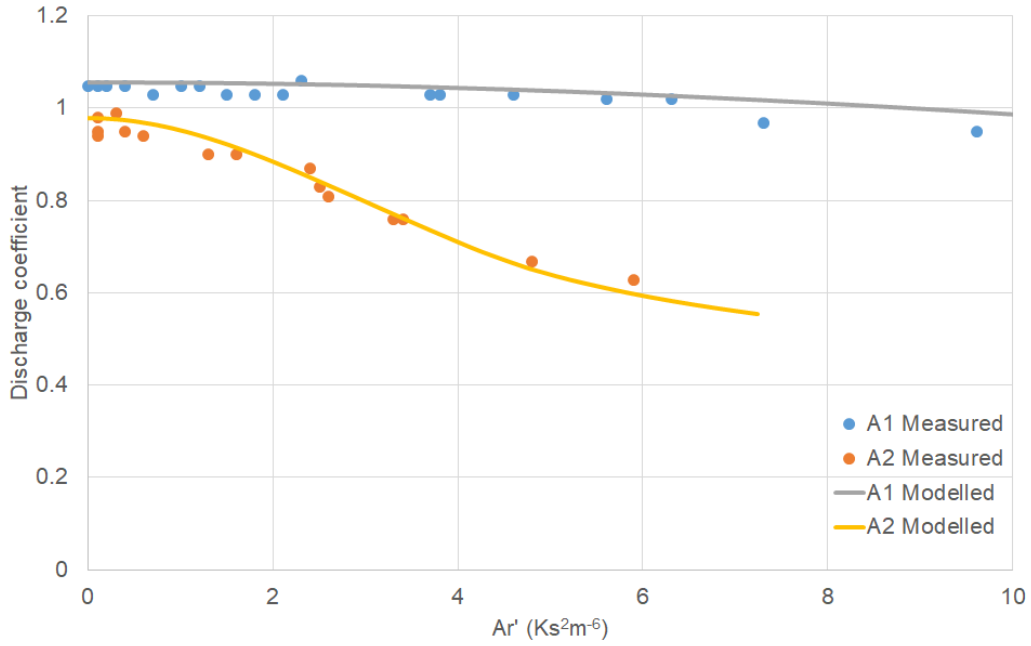


Figure 15: The analytical correction factor compared against the data measured by Heiselberg *et al.* [38].

824 factor to be described independently of the pressure across the window and
 825 the density difference;

$$h^* = \frac{2(z_n - z_{win})}{h} \quad (17)$$

826 where z_n is the neutral pressure height, defined as the height on the facade
 827 where the internal and external pressures are equal, and z_{win} is the height
 828 of the window at its centre. Thus, the correction factor can be found in the
 829 design case by considering the building geometry alone, isolated from the
 830 environmental conditions.

831 *5.1. Comparison with literature data*

832 To enable the analytical model to be compared with the data presented by
 833 Heiselberg *et al.* [38], Equations 6, 15, and 16 have been combined to describe
 834 the relationship between C_q and Ar' that can be solved by iteration.

$$C_q(Ar') = \frac{1}{2} \left(\sqrt{1 + \frac{1000C_q^2 A_{eff}^2 Ar' gh}{T}} + \sqrt{1 - \frac{1000C_q^2 A_{eff}^2 Ar' gh}{T}} \right) \quad (18)$$

835 The height of the window and an estimate of the mean temperature can
 836 be found directly from the reference, and an estimate of the window's *effec-*
 837 *tive* area can be made using the product of the discharge coefficient measured
 838 at $Ar' = 0$ and the stated *free* area used to calculate this discharge coeffi-
 839 cient. The predicted relationship between Ar' and the discharge coefficient
 840 calculated in the reference can then be given by

$$C_d(Ar') = C_{dAr'(0)} C_q(Ar') \quad (19)$$

841 The agreement between the analytical model and the data of Heiselberg *et*
 842 *al.* [38] is good, suggesting it can be used to predict the reduction in the
 843 discharge coefficient; see Figure 15. The model is expected to become in-
 844 creasingly inaccurate as the opening angle increases.

845 **6. Conclusions**

846 The analytical framework is an effective tool for defining different types
 847 of flow through openings. The use of this tool enables a clear and compre-

848 hensive literature review to be made, where different studies and descriptive
849 equations can be compared on a like-for-like basis.

850 The framework identifies a lack of standardised testing regimes for char-
851 acterising opening performance under a range of environmental conditions.
852 In some areas, academic research does not yet provide adequate procedures
853 to evaluate the performance of specific PPOs.

854 While the use of sharp-edged rectangular orifices are ubiquitous in the
855 investigation of ventilation phenomena, aerodynamic performance data for
856 other types of types of PPO are scarce. Use of ambiguous *free* area models
857 are commonplace, and are potentially a source of large variations in measured
858 aerodynamic properties between papers.

859 The framework is used to identify three key areas where understanding
860 can be improved by analysis of existing data: the prediction still-air per-
861 formance characteristics for butt hinged openings; modelling the increase in
862 airflow rate that occurs through two aligned openings; and modelling the
863 reduction in discharge coefficient that occurs when a large temperature dif-
864 ference exists across the opening.

865 A Statistical Effective Area Model (SEAM) is developed from academic
866 data to estimate the still-air performance of hinged openings in the design
867 stage, accounting for the impact of aspect ratio and opening angle. This
868 model predicts literature data with a standard error of 1.2%, compared to
869 a 15 – 25% error offered by *free* area models commonly used in industry.
870 Most analytical *free* area models, such as that given in CIBSE AM10 [3],
871 overestimate airflow through openings and require urgent revision. However,
872 one model used in safety-critical smoke ventilation applications [18] system-

873 atically underestimates flow rate and does not require urgent revision.

874 An analytical model is made based on entrainment theory to explain the
875 increase in flow rate that occurs through two aligned openings. This model
876 identifies the dimensionless ratio \sqrt{A}/x as the characteristic parameter de-
877 scribing the phenomena, rather than the opening porosity commonly cited
878 in the literature. The predictions of the analytical model match CFD predic-
879 tions of airflow rate given in the literature well, and predicts a detrimental
880 impact on pollutant removal from the wider space. The latter phenomena is
881 not identified in the literature, and represents opportunity for further study.

882 Finally, an analytical model is created to explain the reduction in dis-
883 charge coefficient that occurs when a large temperature difference exists
884 across an opening. This model defines a novel dimensionless parameter that
885 characterises the flow based on the ratio of the pressure drop across the cen-
886 tre of the opening to the variation in pressure due to buoyancy across its
887 height. This can be determined in isolation from building geometry, and de-
888 termines whether mono-directional or bidirectional flow is occurring through
889 the opening. The model predicts literature data well, suggesting it can be
890 directly integrated into design equations.

891 The results here suggest a range of avenues where further work may be
892 required, and new predictive tools have been created that can be directly used
893 to reduce design errors in the engineering of a naturally ventilated building.

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